

## RADIATION REACTION WITH A DELTA-FUNCTION EXTERNAL FORCE

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Here's another example of applying an external force to a charge feeling the radiation reaction force. In general, a charge's acceleration obeys the differential equation

$$a = \tau \dot{a} + \frac{F}{m} \quad (1)$$

where  $F$  is the external force and

$$\tau \equiv \frac{\mu_0 q^2}{6\pi m c} \quad (2)$$

Suppose now that the force is a delta function:

$$F = k\delta(t) \quad (3)$$

for some constant  $k$ . In the earlier post, we showed that if  $F$  is finite everywhere, then  $a$  must be continuous everywhere. However, here  $F$  is not finite at  $t = 0$ . As before we start by integrating 1 over a small time interval around  $t = 0$ :

$$\int_{-\epsilon}^{\epsilon} a \, dt = \tau [a(\epsilon) - a(-\epsilon)] + \frac{1}{m} \int_{-\epsilon}^{\epsilon} F \, dt \quad (4)$$

Provided that  $a$  is finite everywhere, the integral on the LHS goes to zero as  $\epsilon \rightarrow 0$  so we're left with

$$\tau \Delta a = -\frac{1}{m} \int_{-\epsilon}^{\epsilon} F \, dt \quad (5)$$

$$= -\frac{k}{m} \int_{-\epsilon}^{\epsilon} \delta(t) \, dt \quad (6)$$

$$= -\frac{k}{m} \quad (7)$$

$$\Delta a = -\frac{k}{m\tau} \quad (8)$$

We can repeat the calculations we did earlier to check that energy is conserved here. Since  $F = 0$  everywhere except  $t = 0$ , the general solution of 1 is

$$a(t) = \begin{cases} a_0 e^{t/\tau} & t < 0 \\ a_1 e^{t/\tau} & t > 0 \end{cases} \quad (9)$$

If we eliminate the runaway acceleration for  $t > 0$  by requiring  $a_1 = 0$  then the condition 8 requires  $a_0 = k/m\tau$ , so

$$a(t) = \begin{cases} \frac{k}{m\tau} e^{t/\tau} & t < 0 \\ 0 & t > 0 \end{cases} \quad (10)$$

By requiring  $v = 0$  at  $t = -\infty$  and that  $v$  is continuous at  $t = 0$  we get

$$v(t) = \begin{cases} \frac{k}{m} e^{t/\tau} & t < 0 \\ \frac{k}{m} & t > 0 \end{cases} \quad (11)$$

The work done by the force is

$$W = \int_{-\infty}^{\infty} Fv dt \quad (12)$$

$$= k \int_{-\infty}^{\infty} \delta(t) v dt \quad (13)$$

$$= kv(0) \quad (14)$$

$$= \frac{k^2}{m} \quad (15)$$

The energy radiated  $R$  is given by integrating the Larmor formula

$$P = \frac{\mu_0 q^2 a^2}{6\pi c} = m\tau a^2 \quad (16)$$

so we get

$$R = m\tau \int_{-\infty}^0 \left( \frac{k}{m\tau} \right)^2 e^{2t/\tau} dt \quad (17)$$

$$= \frac{k^2}{2m} \quad (18)$$

The final kinetic energy is

$$K = \frac{1}{2} m \frac{k^2}{m^2} = \frac{k^2}{2m} \quad (19)$$

Thus

$$W = R + K \quad (20)$$

and energy is conserved.