

POLARIZATION VECTORS

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The 4d wave equation (Maxwell's equations in Lorenz gauge) has the general solution

$$A^\mu(x) = \sum_{r,\mathbf{k}} \frac{1}{\sqrt{2V\omega_{\mathbf{k}}}} \varepsilon_r^\mu A_r(\mathbf{k}) e^{-ikx} + \sum_{r,\mathbf{k}} \frac{1}{\sqrt{2V\omega_{\mathbf{k}}}} \varepsilon_r^\mu A_r^\dagger(\mathbf{k}) e^{ikx} \quad (1)$$

Here, $A_r(\mathbf{k})$ and $B_r^\dagger(\mathbf{k})$ are (possibly complex) numerical coefficients. The ε_r^μ are the polarization vectors, which serve as a basis for the 4d space. We can choose any 4 independent vectors to serve as ε_r .

One choice is to align ε_3 with \mathbf{k} , the direction of travel of the photon. Although there is a sum over \mathbf{k} in 1, we can consider situations where the field is travelling in a single direction, so all the \mathbf{k} vectors are parallel. The vectors can have different frequencies, which we would derive from the boundary conditions on the volume V in which the field is contained.

With this choice, we have

$$\varepsilon_3 = \frac{\mathbf{k}}{|\mathbf{k}|} \quad (2)$$

If we choose the other polarization vectors to be orthogonal, then

$$\varepsilon_1 \cdot \mathbf{k} = \varepsilon_2 \cdot \mathbf{k} = 0 \quad (3)$$

Since electromagnetic waves are transverse waves, the electric and magnetic fields \mathbf{E} and \mathbf{B} are perpendicular to \mathbf{k} , so the vectors ε_1 and ε_2 can be chosen to be parallel to \mathbf{E} and \mathbf{B} .

To see how this works, suppose we also align the coordinate system so that ε_i points along the x^i direction. Now consider the case where A^μ is polarized along the x^1 direction, and consider the spatial part of 1. Also taking A_1 to be real, so that $A_1^\dagger = A_1$, we have, considering only a single vector \mathbf{k} :

$$A^i(x) = \frac{1}{\sqrt{2V\omega_{\mathbf{k}}}} \varepsilon_1^i A_1(\mathbf{k}) [e^{-ikx} + e^{ikx}] \quad (4)$$

$$= \sqrt{\frac{2}{V\omega_{\mathbf{k}}}} \varepsilon_1^i A_1(\mathbf{k}) \cos kx \quad (5)$$

Since the wave is propagating in the $x^3 = z$ direction, we have

$$kx = \omega_{\mathbf{k}}t - pz \quad (6)$$

where $p = |\mathbf{p}|$ is the magnitude of the momentum. Also

$$\varepsilon_1 = \begin{pmatrix} 0 & 1 & 0 & 0 \end{pmatrix} \quad (7)$$

We can now find the magnetic field \mathbf{B} from

$$\mathbf{B} = \nabla \times \mathbf{A} \quad (8)$$

$$= \sqrt{\frac{2}{V\omega_{\mathbf{k}}}} A_1(\mathbf{k}) \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \partial_x & \partial_y & \partial_z \\ \cos kx & 0 & 0 \end{vmatrix} \quad (9)$$

$$= \sqrt{\frac{2}{V\omega_{\mathbf{k}}}} A_1(\mathbf{k}) [\hat{y}(\partial_z \cos kx) + \hat{z}(\partial_y \cos kx)] \quad (10)$$

$$= \sqrt{\frac{2}{V\omega_{\mathbf{k}}}} A_1(\mathbf{k}) \hat{y} (p \sin kx) + 0 \quad (11)$$

$$= \sqrt{\frac{2}{V\omega_{\mathbf{k}}}} A_1(\mathbf{k}) p \sin kx \hat{y} \quad (12)$$

where the $+0$ in the third line is because kx does not depend on y (see 6). Thus the magnetic field is parallel to \hat{y} which is also the x^2 direction, and is perpendicular to \mathbf{A} .

The electric field can be found from

$$\nabla \times \mathbf{E} = \nabla \times \left(-\nabla\Phi - \frac{\partial \mathbf{A}}{\partial t} \right) \quad (13)$$

$$= -\frac{\partial}{\partial t} (\nabla \times \mathbf{A}) \quad (14)$$

$$= -\frac{\partial}{\partial t} \mathbf{B} \quad (15)$$

$$= -\sqrt{\frac{2\omega_{\mathbf{k}}}{V}} A_1(\mathbf{k}) p \cos kx \hat{y} \quad (16)$$

since curl of grad is always zero. Thus the curl of \mathbf{E} is parallel to the curl of \mathbf{A} which is true if \mathbf{E} is parallel to \mathbf{A} , so it lies in the x^1 direction. Both \mathbf{E} and \mathbf{B} are perpendicular to \mathbf{k} , so the wave is transverse, as required.

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