CYCLOID MOTION IN CROSSED ELECTRIC AND MAGNETIC FIELDS

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Reference: Griffiths, David J. (2007) Introduction to Electrodynamics, 3rd Edition; Prentice Hall - Problem 5.2.

One of the classic examples of how mixing electric and magnetic fields can produce unusual results is the case of cycloid motion. The idea is that we have a constant magnetic field pointing in the *x* direction, and a constant electric field perpendicular to the magnetic field. We might as well align the electric field with the *z* axis.

If we place a particle with charge q at the origin and either give it an initial velocity perpendicular to \mathbf{B} or start it at rest, then neither the electric nor the magnetic field will exert a force in the x direction, so the motion is restricted entirely to the yz plane.

Combining the Lorentz force law $\mathbf{F} = q\mathbf{v} \times \mathbf{B}$ and electrostatics with Newton's law $\mathbf{F} = m\mathbf{a}$ we get, using a dot to denote a derivative w.r.t. time:

$$(0.1) m\ddot{y} = qB\dot{z}$$

$$m\ddot{z} = qE - qB\dot{y}$$

At this point, we can note that it's quite easy to generalize this system so that \mathbf{E} has an arbitrary direction perpendicular to \mathbf{B} . However, there's not much point in doing so, since this amounts merely to a rotation of the coordinates about the x axis. Giving \mathbf{E} a component along x, however, does introduce significant complexity into the system, since we then have a force component along x, so we need another differential equation in the above set.

We can solve the system of 2 ODEs above by using software (or by hand, if you're familiar with methods of solving systems like this), and we get

(0.3)
$$y(t) = -\frac{c_3}{\omega}\cos\omega t + \frac{c_2}{\omega}\sin\omega t + \frac{E}{B}t + c_4$$

(0.4)
$$z(t) = \frac{c_3}{\omega} \sin \omega t + \frac{c_2}{\omega} \cos \omega t + \frac{E}{\omega B} + c_1$$

where $\omega \equiv \frac{qB}{m}$ and the c_i s are constants of integration. Note that ω has units of inverse time, as we can see from the Lorentz force law above. qB

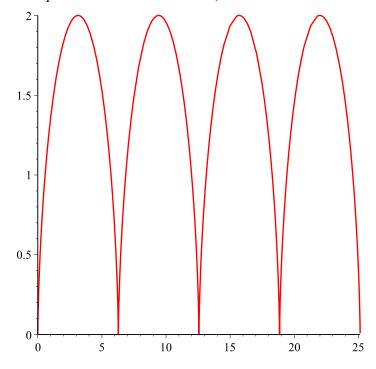
has units of force/velocity, which has units of mass \times time⁻¹. (For those reading Griffiths's book, the constants in our solution can be redefined to produce Griffiths's equation 5.6, since E, B, q, m and ω are all constants here.)

The four constants can be determined by specifying the initial position and velocity of the particle. If it starts at rest at the origin (that is, x(0) = $y(0) = \dot{x}(0) = \dot{y}(0) = 0$), we can apply these conditions to the solution (again, using software to ease the mathematics), and get

$$(0.5) y(t) = -\frac{E}{\omega B}\sin \omega t + \frac{E}{B}t$$

$$(0.6) z(t) = -\frac{E}{\omega B}\cos \omega t + \frac{E}{\omega B}$$

These two equations give the trajectory in parametric form, so we can vary t to get a plot of z versus y (we've taken all the constants to be 1 E = B = q = m = 1 for convenience):



Taking an initial velocity of $\frac{E}{B}$ along the y axis $(\dot{y}(0) = \frac{E}{B})$, we get

$$y(t) = \frac{E}{B}t$$

$$(0.8) z(t) = 0$$

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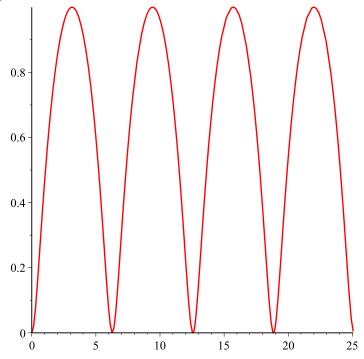
In this case the particle moves at constant speed along the y axis. Why? The electric force is qE in the +z direction (assuming q is positive), and the initial magnetic force is $qvB = q\frac{E}{B}B = qE$ in the -z direction (use the right hand rule in the cross product). Thus the two forces cancel each other, so there is no net acceleration and the initial velocity remains unchanged.

If
$$\dot{y}(0) = \frac{E}{2B}$$
, we get

$$(0.9) y(t) = -\frac{E}{2\omega B}\sin \omega t + \frac{E}{B}t$$

$$(0.10) z(t) = -\frac{E}{2\omega B}\cos \omega t + \frac{E}{2\omega B}$$

This gives a similar trajectory to the first example, but with a reduced amplitude:

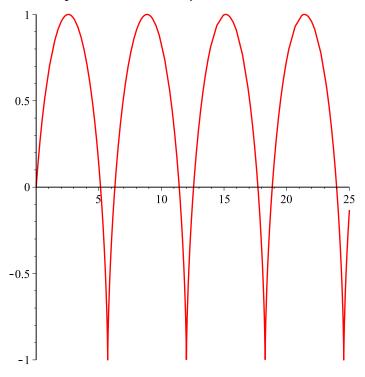


Finally, if we start the particle off travelling at an angle of $\pi/4$ w.r.t. the y axis, so that $\dot{y}(0) = \dot{z}(0) = \frac{E}{B}$, we get

(0.11)
$$y(t) = -\frac{E}{\omega B}\cos \omega t + \frac{E}{B}t + \frac{E}{\omega B}$$

$$(0.12) z(t) = \frac{E}{\omega B} \sin \omega t$$

This gives a trajectory identical in shape to the first case, but shifted down so that its midpoint is now on the *y* axis:



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