

MAGNETIC DIPOLE FIELD OF A FINITE SOLENOID

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Reference: Griffiths, David J. (2007) Introduction to Electrodynamics, 3rd Edition; Prentice Hall - Problem 5.61.

For our last post on magnetostatics, we'll consider a finite solenoid of radius R and length L , with a surface charge density σ rotating at angular speed ω . We know that the field outside an infinite solenoid is zero, but what about a finite solenoid? For points far from the axis, we can use a dipole approximation.

We align the axis along the z axis, and consider it to be a stack of individual current loops, each with its own dipole moment. The moment of a current loop is

$$(1) \quad \mathbf{m} = \pi IR^2 \hat{\mathbf{z}}$$

In terms of the parameters of the problem, each loop has a thickness of dz and thus carries a current of $I = \sigma R \omega dz$. The contribution of the loop at coordinate z is therefore

$$(2) \quad d\mathbf{m} = \pi \omega \sigma R^3 dz \hat{\mathbf{z}}$$

The field of the dipole from this current loop is

$$(3) \quad d\mathbf{B}_{dip} = \frac{\mu_0}{4\pi r^3} [3(d\mathbf{m} \cdot \hat{\mathbf{r}}) \hat{\mathbf{r}} - d\mathbf{m}]$$

We now need to consider carefully what is meant by $\hat{\mathbf{r}}$. We'll take the observation point to be a distance s along a line perpendicular to the axis and intersecting the axis at its midpoint. We'll define s to be on the $\hat{\mathbf{x}}$ axis. Then for a given current loop at coordinate z , the vector \mathbf{r} points from the centre of this loop to s . Therefore

$$\begin{aligned}
(4) \quad r &= \sqrt{s^2 + z^2} \\
(5) \quad \cos \theta &= -\frac{z}{r} \\
(6) \quad \sin \theta &= \frac{s}{r} \\
(7) \quad d\mathbf{m} \cdot \hat{\mathbf{r}} &= \pi\omega\sigma R^3 \cos \theta dz \\
(8) \quad &= -\pi\omega\sigma R^3 \frac{z}{r} dz
\end{aligned}$$

where θ is, as usual, the angle between $\hat{\mathbf{r}}$ and $\hat{\mathbf{z}}$. Therefore

$$\begin{aligned}
(9) \quad \hat{\mathbf{r}} &= \cos \theta \hat{\mathbf{z}} + \sin \theta \hat{\mathbf{x}} \\
(10) \quad &= -\frac{z}{r} \hat{\mathbf{z}} + \frac{s}{r} \hat{\mathbf{x}}
\end{aligned}$$

The total dipole field of the solenoid is therefore

$$(11) \quad \mathbf{B}_{dip} = \frac{\mu_0}{4\pi} \pi\omega\sigma R^3 \int_{-L/2}^{L/2} \left[\left(\frac{3z^2}{(s^2 + z^2)^{5/2}} - \frac{1}{(s^2 + z^2)^{3/2}} \right) \hat{\mathbf{z}} + \frac{sz}{(s^2 + z^2)^{5/2}} \hat{\mathbf{x}} \right] dz$$

$$(12) \quad = -\frac{\mu_0\omega\sigma R^3 L}{4 \left(s^2 + \left(\frac{L}{2} \right)^2 \right)^{3/2}} \hat{\mathbf{z}} + 0 \hat{\mathbf{x}}$$

$$(13) \quad = -\frac{\mu_0\omega\sigma R^3 L}{4 \left(s^2 + \left(\frac{L}{2} \right)^2 \right)^{3/2}} \hat{\mathbf{z}}$$

Note that as $L \rightarrow \infty$, the field does tend to zero as $\frac{1}{L^2}$ which is the correct value for an infinite solenoid.