

## COAXIAL CABLE WITH VARYING CURRENT

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References: Griffiths, David J. (2007), Introduction to Electrodynamics, 3rd Edition; Pearson Education - Problem 7.16.

Here's another (mis)application of Faraday's law, using the quasistatic approximation to calculate the fields involved. This time we have a coaxial cable in which the current flows in the  $+z$  direction on the inner wire and back in the  $-z$  direction on the coaxial cylinder. The current varies with time according to

$$I(t) = I_0 \cos \omega t \quad (1)$$

The magnetic field due to the inner wire is given by

$$\mathbf{B}_w(r, t) = \hat{\phi} \frac{\mu_0 I_0}{2\pi r} \cos \omega t \quad (2)$$

The magnetic field due to the coaxial cylinder is just the negative of this outside the cylinder, and zero inside:

$$\mathbf{B}_c(r, t) = \begin{cases} -\hat{\phi} \frac{\mu_0 I_0}{2\pi r} \cos \omega t & r > a \\ 0 & r < a \end{cases} \quad (3)$$

Therefore the net magnetic field outside the cylinder is always zero, and the field inside is given by  $\mathbf{B}_w$  alone. The induced electric field is parallel to the  $z$  axis, as can be seen by a similar argument to that given for the magnetic field of a solenoid. First,  $\mathbf{E}$  cannot have a radial component, since if we reversed the directions of the currents, then  $\mathbf{B}$  and hence  $\mathbf{E}$  reverse their directions. But reversing the currents is equivalent to turning the wire through 180 degrees, and that shouldn't affect the radial component of the field, so  $E_r = 0$ .

$\mathbf{E}$  also cannot have a  $\phi$  component. Since  $\mathbf{B}$  is entirely in the  $\phi$  direction, so is  $-\partial\mathbf{B}/\partial t$ , so if we choose a path of integration that is a circle at constant  $r$ , and a corresponding surface of integration that is perpendicular to the  $z$  direction (that is, to the cable), then  $d\mathbf{a}$  is always perpendicular to  $-\partial\mathbf{B}/\partial t$  so the integral on the right is zero:

$$\oint \mathbf{E} \cdot d\mathbf{l} = - \int \frac{\partial\mathbf{B}}{\partial t} \cdot d\mathbf{a} \quad (4)$$

By symmetry,  $E_\phi$  must be constant around the circle so we get  $\oint \mathbf{E} \cdot d\ell = 2\pi r E_\phi = 0$ , so  $E_\phi = 0$ . Therefore,  $\mathbf{E}$  must have a  $z$  component only.

To figure out the actual value of  $\mathbf{E}$ , we can use the same argument as in calculating the solenoid's magnetic field. We choose a square loop with side length  $\ell$  that lies in the  $yz$  plane. First, we take the loop to lie entirely outside the coaxial cable, where  $\mathbf{B} = 0$  always. This means that (taking the near edge to be a distance  $s$  from the cable's axis):

$$\oint \mathbf{E} \cdot d\ell = \ell (E(s+\ell) - E(s)) = 0 \quad (5)$$

If we assume that  $\mathbf{E} \rightarrow 0$  at infinity, then  $\mathbf{E} = 0$  everywhere outside the cable.

To find  $\mathbf{E}$  inside the cable, we can take the loop of integration to have its inner edge at distance  $s$  from the axis, where  $0 < s < a$ , and the outer edge outside the cylinder, where  $\mathbf{E} = 0$ . Then

$$\oint \mathbf{E} \cdot d\ell = \ell E(s) \quad (6)$$

$$= - \int \frac{\partial \mathbf{B}}{\partial t} \cdot d\mathbf{a} \quad (7)$$

$$= \frac{\ell \omega \mu_0 I_0}{2\pi} \sin \omega t \int_s^a \frac{dr}{r} \quad (8)$$

$$= \frac{\ell \omega \mu_0 I_0}{2\pi} \sin \omega t \ln \frac{a}{s} \quad (9)$$

To get the direction of  $\mathbf{E}$  we use Lenz's law and the right hand rule. If the current flows in the  $+z$  direction and is increasing, then the magnetic field is in the  $+\phi$  direction and increasing, so the flux through the loop is increasing. The induced electric field will oppose this increase, which means that  $\mathbf{E}$  inside the cylinder must point in the  $-z$  direction. Conversely, if  $I$  is decreasing,  $\mathbf{E}$  points in the  $+z$  direction. Since  $I$  starts off at its maximum value at  $t = 0$ , the current is decreasing initially, so the initial value of  $\mathbf{E}$  must point in the  $+z$  direction, thus:

$$\mathbf{E}(s, t) = \begin{cases} \hat{\mathbf{z}} \frac{\omega \mu_0 I_0}{2\pi} \sin \omega t \ln \frac{a}{s} & 0 < s < a \\ 0 & s > a \end{cases} \quad (10)$$

As always, we need to issue a caution that this analysis is deeply flawed, since we're applying the quasistatic approximation to a cable of infinite length, and we're assuming that the current can change everywhere along the length of the cable at the same rate, which doesn't happen in real life.

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