

ANGULAR MOMENTUM IN ELECTROMAGNETIC FIELDS

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References: Griffiths, David J. (2007), Introduction to Electrodynamics, 3rd Edition; Pearson Education - Problem 8.7.

The momentum density of an electromagnetic field is given by

$$(1) \quad \mathbf{p}_{em} = \epsilon_0 \mu_0 \mathbf{S} = \epsilon_0 \mathbf{E} \times \mathbf{B}$$

If we have linear momentum, then we automatically have *angular momentum* with respect to some origin by using the classical definition of angular momentum $\mathbf{L} = \mathbf{r} \times \mathbf{p}$. We can define the angular momentum density of an electromagnetic field by

$$(2) \quad \mathfrak{L}_{em} \equiv \mathbf{r} \times \mathbf{p}_{em} = \epsilon_0 \mathbf{r} \times (\mathbf{E} \times \mathbf{B})$$

Just as with linear momentum, even static fields can have angular momentum. As an example, suppose we have a long solenoid with n turns per unit length carrying current I_0 and a radius R , with its axis along the z axis. The magnetic field inside the solenoid is

$$(3) \quad \mathbf{B}_0 = \mu_0 n I_0 \hat{\mathbf{z}}$$

The field is zero outside the solenoid.

Now suppose we add two other cylinders (not solenoids), both coaxial with the solenoid. One cylinder has radius $a < R$ (so it lies inside the solenoid) and carries surface charge $+Q$; the other cylinder has radius $b > R$ (outside the solenoid) and carries charge $-Q$. Both cylinders have length ℓ . From Gauss's law, the electric field between these two cylinders is, for $a < r < b$

$$(4) \quad \mathbf{E}_0 = \frac{Q}{2\pi\epsilon_0\ell r} \hat{\mathbf{r}}$$

That is, the field points radially outward from the axis. The electric field is zero for $r < a$ and $r > b$. (We're neglecting end effects, so we're assuming that $\ell \gg b > a$.)

The linear momentum density is non-zero in the region $a < r < R$ (where both fields are non-zero) and we have

$$(5) \quad \mathbf{p}_{em} = -\frac{\mu_0 n I Q}{2\pi\ell} \frac{\hat{\phi}}{r}$$

so the angular momentum density is

$$(6) \quad \mathfrak{L}_{em} = \mathbf{r} \times \mathbf{p}_{em}$$

$$(7) \quad = -\frac{\mu_0 n I Q}{2\pi\ell} \hat{\mathbf{z}}$$

Conveniently, this is constant so the total angular momentum is just the density times the volume of the cylindrical tube in the region $a < r < R$

$$(8) \quad \mathbf{L}_{em} = -\ell\pi(R^2 - a^2) \frac{\mu_0 n I Q}{2\pi\ell} \hat{\mathbf{z}}$$

$$(9) \quad = -(R^2 - a^2) \frac{\mu_0 n I Q}{2} \hat{\mathbf{z}}$$

Now suppose we (quasistatically) discharge the two cylinders by connecting a resistor \mathcal{R} between them. We'd like to show that the angular momentum gets transferred from the fields to the physical devices in the problem. The two cylinders are effectively a capacitor with some capacitance C , so we know that the current in the resistor will decay exponentially

$$(10) \quad I(t) = \frac{V_0}{\mathcal{R}} e^{-t/\mathcal{R}C}$$

where V_0 is the potential difference between the cylinders at $t = 0$. The force $d\mathbf{F}$ on a segment of the resistor of length dr is

$$(11) \quad d\mathbf{F} = I(t) dr \hat{\mathbf{r}} \times \mathbf{B}_0$$

$$(12) \quad = -I(t) dr B_0 \hat{\phi}$$

so the torque on this segment is

$$(13) \quad d\mathbf{N} = \mathbf{r} \times d\mathbf{F}$$

$$(14) \quad = -I(t) B_0 r dr \hat{\mathbf{z}}$$

The total torque on the resistor at time t is

$$(15) \quad \mathbf{N}(t) = -I(t)B_0\hat{\mathbf{z}} \int_a^R r dr$$

$$(16) \quad = -\frac{1}{2}I(t)B_0(R^2 - a^2)\hat{\mathbf{z}}$$

The angular impulse is the integral of torque over time, so we get

$$(17) \quad \mathbf{I} = \int_0^\infty \mathbf{N}(t) dt$$

$$(18) \quad = -\frac{1}{2}B_0(R^2 - a^2)\hat{\mathbf{z}} \int_0^\infty I(t) dt$$

$$(19) \quad = -\frac{1}{2}B_0(R^2 - a^2)\hat{\mathbf{z}} \int_0^\infty \frac{V_0}{\mathcal{R}} e^{-t/\mathcal{R}C} dt$$

$$(20) \quad = -\frac{1}{2}B_0(R^2 - a^2)CV_0\hat{\mathbf{z}}$$

$$(21) \quad = -\frac{1}{2}\mu_0 nI(R^2 - a^2)Q\hat{\mathbf{z}}$$

where we used the relation between capacitance, charge and voltage $Q = CV$. We see that this agrees with 9, so all the angular momentum in the fields is transferred to the resistor as the electric field is reduced to zero.

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