

## RADIATION REACTION WITH A DELTA-FUNCTION EXTERNAL FORCE

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References: Griffiths, David J. (2007), Introduction to Electrodynamics, 3rd Edition; Pearson Education - Chapter 11, Post 28.

Here's another example of applying an external force to a charge feeling the radiation reaction force. In general, a charge's acceleration obeys the differential equation

$$(1) \quad a = \tau \dot{a} + \frac{F}{m}$$

where  $F$  is the external force and

$$(2) \quad \tau \equiv \frac{\mu_0 q^2}{6\pi m c}$$

Suppose now that the force is a delta function:

$$(3) \quad F = k\delta(t)$$

for some constant  $k$ . In the earlier post, we showed that if  $F$  is finite everywhere, then  $a$  must be continuous everywhere. However, here  $F$  is not finite at  $t = 0$ . As before we start by integrating 1 over a small time interval around  $t = 0$ :

$$(4) \quad \int_{-\varepsilon}^{\varepsilon} a \, dt = \tau [a(\varepsilon) - a(-\varepsilon)] + \frac{1}{m} \int_{-\varepsilon}^{\varepsilon} F \, dt$$

Provided that  $a$  is finite everywhere, the integral on the LHS goes to zero as  $\varepsilon \rightarrow 0$  so we're left with

$$\begin{aligned}
(5) \quad \tau \Delta a &= -\frac{1}{m} \int_{-\varepsilon}^{\varepsilon} F dt \\
(6) &= -\frac{k}{m} \int_{-\varepsilon}^{\varepsilon} \delta(t) dt \\
(7) &= -\frac{k}{m} \\
(8) \quad \Delta a &= -\frac{k}{m\tau}
\end{aligned}$$

We can repeat the calculations we did earlier to check that energy is conserved here. Since  $F = 0$  everywhere except  $t = 0$ , the general solution of 1 is

$$(9) \quad a(t) = \begin{cases} a_0 e^{t/\tau} & t < 0 \\ a_1 e^{t/\tau} & t > 0 \end{cases}$$

If we eliminate the runaway acceleration for  $t > 0$  by requiring  $a_1 = 0$  then the condition 8 requires  $a_0 = k/m\tau$ , so

$$(10) \quad a(t) = \begin{cases} \frac{k}{m\tau} e^{t/\tau} & t < 0 \\ 0 & t > 0 \end{cases}$$

By requiring  $v = 0$  at  $t = -\infty$  and that  $v$  is continuous at  $t = 0$  we get

$$(11) \quad v(t) = \begin{cases} \frac{k}{m} e^{t/\tau} & t < 0 \\ \frac{k}{m} & t > 0 \end{cases}$$

The work done by the force is

$$\begin{aligned}
(12) \quad W &= \int_{-\infty}^{\infty} F v dt \\
(13) &= k \int_{-\infty}^{\infty} \delta(t) v dt \\
(14) &= kv(0) \\
(15) &= \frac{k^2}{m}
\end{aligned}$$

The energy radiated  $R$  is given by integrating the Larmor formula

$$(16) \quad P = \frac{\mu_0 q^2 a^2}{6\pi c} = m\tau a^2$$

so we get

$$(17) \quad R = m\tau \int_{-\infty}^0 \left( \frac{k}{m\tau} \right)^2 e^{2t/\tau} dt$$

$$(18) \quad = \frac{k^2}{2m}$$

The final kinetic energy is

$$(19) \quad K = \frac{1}{2} m \frac{k^2}{m^2} = \frac{k^2}{2m}$$

Thus

$$(20) \quad W = R + K$$

and energy is conserved.

#### PINGBACKS

Pingback: Tunnelling through a potential barrier with the radiation reaction force