

## ELECTROMAGNETIC FIELD TENSOR: A COUPLE OF EXAMPLES

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Reference: Griffiths, David J. (2007), Introduction to Electrodynamics, 3rd Edition; Pearson Education - Chapter 12, Problems 12.52-53.

Maxwell's equations can be expressed quite simply in terms of the electromagnetic field tensor

$$(1) \quad F^{ij} = \begin{bmatrix} 0 & E_x & E_y & E_z \\ -E_x & 0 & B_z & -B_y \\ -E_y & -B_z & 0 & B_x \\ -E_z & B_y & -B_x & 0 \end{bmatrix}$$

and its dual tensor:

$$(2) \quad G^{ij} = \begin{bmatrix} 0 & B_x & B_y & B_z \\ -B_x & 0 & -E_z & E_y \\ -B_y & E_z & 0 & -E_x \\ -B_z & -E_y & E_x & 0 \end{bmatrix}$$

In his section 12.3.4, Griffiths shows that all four of Maxwell's equations can be expressed as

$$(3) \quad \partial_j F^{ij} = \mu_0 J^i$$

$$(4) \quad \partial_j G^{ij} = 0$$

where  $J^i$  is a four-vector known as the four-current with components

$$(5) \quad J^i = [c\rho, J_x, J_y, J_z]$$

$$(6) \quad = \frac{\rho_0}{\sqrt{1-u^2/c^2}} [c, u_x, u_y, u_z]$$

In this formula, we're looking at a small element of charge that is moving with velocity  $\mathbf{u}$  and has charge density  $\rho_0$  in its rest frame. The four-current is just  $\rho_0$  times the four-velocity:

$$(7) \quad J^i = \rho_0 \eta^i$$

If we take the divergence (in four dimensions) of 3 we get, since  $F^{ij} = -F^{ji}$

$$(8) \quad \partial_i \partial_j F^{ij} = -\partial_i \partial_j F^{ji} = 0$$

which leads to

$$(9) \quad \partial_i J^i = 0$$

This is a compact way of expressing the continuity condition, since from 5

$$(10) \quad \partial_i J^i = \frac{\partial (c\rho)}{\partial x^0} + \nabla \cdot \mathbf{J}$$

$$(11) \quad = \frac{\partial \rho}{\partial t} + \nabla \cdot \mathbf{J} = 0$$

The equation 4 can be written in terms of  $F_{ij}$  as follows. First, we can write  $G^{ij}$  in terms of the components of  $F_{ij}$  by comparing the two matrices 2 and  $F_{ij}$ :

$$(12) \quad F_{ij} = \begin{bmatrix} 0 & -E_x & -E_y & -E_z \\ E_x & 0 & B_z & -B_y \\ E_y & -B_z & 0 & B_x \\ E_z & B_y & -B_x & 0 \end{bmatrix}$$

We get

$$(13) \quad G^{ij} = \begin{bmatrix} 0 & F_{23} & F_{31} & F_{12} \\ F_{32} & 0 & F_{03} & F_{20} \\ F_{13} & F_{30} & 0 & F_{01} \\ F_{21} & F_{02} & F_{10} & 0 \end{bmatrix}$$

Now we can read off the four equations contained in 4 by scanning the rows of this matrix:

$$(14) \quad \partial_j G^{0j} = \partial_1 F_{23} + \partial_2 F_{31} + \partial_3 F_{12}$$

$$(15) \quad \partial_j G^{1j} = \partial_0 F_{32} + \partial_2 F_{03} + \partial_3 F_{20}$$

$$(16) \quad \partial_j G^{2j} = \partial_0 F_{13} + \partial_1 F_{30} + \partial_3 F_{01}$$

$$(17) \quad \partial_j G^{3j} = \partial_0 F_{21} + \partial_1 F_{02} + \partial_2 F_{10}$$

All four of these equations have the form

$$(18) \quad \partial_j G^{ij} = \partial_a F_{bc} + \partial_b F_{ca} + \partial_c F_{ab}$$

where  $i, a, b$  and  $c$  are all different. Note that the indexes  $a, b$  and  $c$  are cyclically permuted in each term on the RHS. Therefore we can write 4 as

$$(19) \quad \partial_a F_{bc} + \partial_b F_{ca} + \partial_c F_{ab} = 0$$

where  $a, b$  and  $c$  are a subset of  $\{0, 1, 2, 3\}$ .

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