

HYDROGEN ATOM - MEAN RADIUS OF ELECTRON POSITION

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References: Griffiths, David J. (2005), Introduction to Quantum Mechanics, 2nd Edition; Pearson Education - Problem 4.13.

We can work out the mean and mean square distances for an electron in a hydrogen atom using the standard method of interpreting the square modulus of the wave function as a probability density.

In the ground state, the quantum numbers are $n = 1$, $l = 0$, $m = 0$ and the wave function can be worked out from the formulas we have used earlier.

$$\begin{aligned} (1) \quad \psi_{100} &= R_{10}Y_0^0 \\ (2) \quad &= c_0 \frac{1}{a} e^{-r/a} Y_0^0 \end{aligned}$$

The constant c_0 can be found from normalization as usual, and we have $c_0 = 2/\sqrt{a}$ so, with the normalized spherical harmonic $Y_0^0 = 1/\sqrt{4\pi}$ we have

$$(3) \quad \psi_{100} = \frac{1}{\sqrt{\pi}a^{3/2}} e^{-r/a}$$

Since the ground state does not depend on either angle, we have

$$(4) \quad \langle r \rangle = \int_0^{2\pi} \int_0^\pi \int_0^\infty r |R_{10}|^2 r^2 |Y_0^0|^2 \sin \theta dr d\theta d\phi$$

$$(5) \quad \int_0^\infty r |R_{10}|^2 r^2 dr$$

$$(6) \quad = 4a^{-3} \int_0^\infty r^3 e^{-2r/a} dr$$

$$(7) \quad = \frac{3}{2}a$$

For the mean square:

$$\begin{aligned}
 (8) \quad \langle r^2 \rangle &= \int_0^\infty r^2 |R_{10}|^2 r^2 dr \\
 (9) &= 4a^{-3} \int_0^\infty r^4 e^{-2r/a} dr \\
 (10) &= 3a^2
 \end{aligned}$$

Since the ground state is symmetric, we can work out the means for the rectangular coordinates separately without doing any more integrals:

$$\begin{aligned}
 (11) \quad \langle x \rangle &= 0 \\
 (12) \quad \langle x^2 \rangle &= \frac{1}{3} \langle r^2 \rangle \\
 (13) &= a^2
 \end{aligned}$$

The same results would be obtained for y and z .

For higher states, we don't have the spherical symmetry of the ground state. As an example, we'll look at the state $n = 2, l = 1, m = 1$. We need $R_{21}(r) = \frac{1}{\sqrt{24}} a^{-3/2} (r/a) e^{-r/2a}$ from our previous example and the spherical harmonic $Y_1^1(\theta, \phi) = -(\sqrt{3/8\pi}) \sin \theta e^{i\phi}$. Using $x = r \sin \theta \cos \phi$ we get

$$\begin{aligned}
 (14) \quad \langle x^2 \rangle &= \int_0^{2\pi} \int_0^\pi \int_0^\infty (r \sin \theta \cos \phi)^2 \left(\frac{1}{\sqrt{24}} a^{-3/2} (r/a) e^{-r/2a} \right)^2 (\sqrt{3/8\pi} \sin \theta)^2 r^2 \sin \theta dr d\theta d\phi \\
 (15) &= 12a^2
 \end{aligned}$$

using Maple to do the integrals.

We can use the same technique to find means for the other coordinates. Because ψ_{211} is independent of ϕ , $\langle x \rangle = \langle y \rangle = 0$ and $\langle y^2 \rangle = \langle x^2 \rangle = 12a^2$. In the z direction, the wave function is even about the angle $\theta = \pi/2$, so we would expect $\langle z \rangle = 0$. We can work out $\langle z^2 \rangle$ by doing the integral using $z = r \cos \theta$ and we get $\langle z^2 \rangle = 6a^2$.

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