

## SCHRÖDINGER EQUATION FOR 2 PARTICLES - SEPARATION OF VARIABLES

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Reference: Griffiths, David J. (2005), Introduction to Quantum Mechanics, 2nd Edition; Pearson Education - Problem 5.1.

The Schrödinger equation that we've looked at so far involves the wave function for a single particle moving in a potential. To extend this to multi-particle systems, we need to make the wave function and the potential functions of the positions of all the particles and the time. Thus the Schrödinger equation for a system of  $n$  particles becomes

$$(1) \quad -\frac{\hbar^2}{2} \sum_{i=1}^n \frac{1}{m_i} \nabla_i^2 \Psi(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_n, t) +$$

$$(2) \quad V(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_n, t) \Psi(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_n, t) = i\hbar \frac{\partial}{\partial t} \Psi(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_n, t)$$

Needless to say, finding solutions of this equation for even as few as 2 particles is extremely difficult. In one case, however, we can make some progress. In a 2-particle system, if the potential  $V$  is not time-dependent and depends only on the separation  $\mathbf{r} \equiv \mathbf{r}_1 - \mathbf{r}_2$  of the two particles, we can fiddle with it a bit and produce a simpler form.

First, we define the centre of mass

$$(3) \quad \mathbf{R} \equiv \frac{m_1 \mathbf{r}_1 + m_2 \mathbf{r}_2}{m_1 + m_2}$$

If we also introduce the reduced mass

$$(4) \quad \mu \equiv \frac{m_1 m_2}{m_1 + m_2}$$

then we get

$$(5) \quad \mathbf{r}_1 = \mathbf{R} + \frac{\mu}{m_1} \mathbf{r}$$

$$(6) \quad \mathbf{r}_2 = \mathbf{R} - \frac{\mu}{m_2} \mathbf{r}$$

In the coordinates  $\mathbf{r}$  and  $\mathbf{R}$ , we can find the gradient operators. We use a dummy function  $f$  to give the gradient something to operate on. We'll consider the  $x$  component and use the chain rule (remember that  $\mathbf{r}_1$  and  $\mathbf{r}_2$  are independent vectors, each with 3 components, so there is a total of 6 independent position variables):

$$\begin{aligned}
 (7) \quad \frac{\partial f}{\partial r_{1x}} &= \frac{\partial f}{\partial r_x} \frac{\partial r_x}{\partial r_{1x}} + \frac{\partial f}{\partial R_x} \frac{\partial R_x}{\partial r_{1x}} \\
 (8) \quad &= \frac{\partial f}{\partial r_x} (1) + \frac{\partial f}{\partial R_x} \frac{m_1}{m_1 + m_2} \\
 (9) \quad &= \frac{\partial f}{\partial r_x} + \frac{\partial f}{\partial R_x} \frac{\mu}{m_2} \\
 (10) \quad \frac{\partial f}{\partial r_{2x}} &= \frac{\partial f}{\partial r_x} (-1) + \frac{\partial f}{\partial R_x} \frac{m_2}{m_1 + m_2} \\
 (11) \quad &= -\frac{\partial f}{\partial r_x} + \frac{\partial f}{\partial R_x} \frac{\mu}{m_1}
 \end{aligned}$$

The relations for the other two components are similar, so dropping the test function  $f$ , we get for the gradients:

$$\begin{aligned}
 (12) \quad \nabla_1 &= \nabla_r + \frac{\mu}{m_2} \nabla_R \\
 (13) \quad \nabla_2 &= -\nabla_r + \frac{\mu}{m_1} \nabla_R
 \end{aligned}$$

To get the Laplacian operators, we differentiate the  $x$  component expressions above.

$$\begin{aligned}
 (14) \quad \frac{\partial^2 f}{\partial r_{1x}^2} &= \left( \frac{\partial}{\partial r_x} + \frac{\mu}{m_2} \frac{\partial}{\partial R_x} \right) \left( \frac{\partial f}{\partial r_x} + \frac{\partial f}{\partial R_x} \frac{\mu}{m_2} \right) \\
 (15) \quad &= \frac{\partial^2 f}{\partial r_x^2} + \frac{\partial f}{\partial r_x \partial R_x} \frac{2m_1}{m_1 + m_2} + \frac{\partial^2 f}{\partial R_x^2} \left( \frac{\mu}{m_2} \right)^2 \\
 (16) \quad \frac{\partial^2 f}{\partial r_{2x}^2} &= \frac{\partial^2 f}{\partial r_x^2} - \frac{\partial f}{\partial r_x \partial R_x} \frac{2m_2}{m_1 + m_2} + \frac{\partial^2 f}{\partial R_x^2} \left( \frac{\mu}{m_1} \right)^2
 \end{aligned}$$

The combination of these two expressions that appears in the Schrödinger equation is, after cancelling terms and putting the remaining terms over common denominators:

$$(17) \quad -\frac{\hbar^2}{2m_1} \frac{\partial^2 f}{\partial r_{1x}^2} - \frac{\hbar^2}{2m_2} \frac{\partial^2 f}{\partial r_{2x}^2} = -\frac{\hbar^2}{2(m_1 + m_2)} \frac{\partial^2 f}{\partial R_x^2} - \frac{\hbar^2}{2\mu} \frac{\partial^2 f}{\partial r_x^2}$$

The calculations for the other two components are similar, so the final result is:

$$(18) \quad -\frac{\hbar^2}{2(m_1 + m_2)} \nabla_R^2 \psi - \frac{\hbar^2}{2\mu} \nabla_r^2 \psi + V(\mathbf{r}) \psi = E \psi$$

We can now try the usual technique of separation of variables, so we try

$$(19) \quad \psi = A(r)B(R)$$

We get, after substituting and dividing through by  $AB$ :

$$(20) \quad -\frac{\hbar^2}{2(m_1 + m_2)B} \nabla_R^2 B - \frac{\hbar^2}{2\mu A} \nabla_r^2 A + V(\mathbf{r}) = E$$

As usual, the terms involving each of the variables  $r$  and  $R$  must separately be equal to constants, so we get

$$(21) \quad -\frac{\hbar^2}{2(m_1 + m_2)} \nabla_R^2 B = E_R B$$

$$(22) \quad -\frac{\hbar^2}{2\mu} \nabla_r^2 A + AV(\mathbf{r}) = E_r A$$

The first equation is that of a free particle with mass  $m_1 + m_2$ , while the second is that of a particle with mass  $\mu$  moving in a potential  $V$ . Thus the system separates into one equation for a free particle with the total mass and a position at the centre of mass and another for a single particle with the reduced mass  $\mu$ . The total energy is  $E = E_R + E_r$ .

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