

## HARMONIC OSCILLATOR IN AN ELECTRIC FIELD

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References: Griffiths, David J. (2005), Introduction to Quantum Mechanics, 2nd Edition; Pearson Education - Problem 6.5.

As another example of second order perturbation theory we start with the harmonic oscillator potential, except this time we place a charged particle (charge  $q$ ) into the potential and turn on a small electric field  $E$ , so that the perturbation in the potential is

$$V = -qEx \quad (1)$$

We'll begin by looking at the first order correction, for which we have

$$E_{n1} = \langle n0 | V | n0 \rangle \quad (2)$$

From our earlier calculations of the matrix elements in the harmonic oscillator potential, we have

$$\langle n0 | x | j0 \rangle = \sqrt{\frac{\hbar}{2m\omega}} \left( \sqrt{j+1} \delta_{n,j+1} + \sqrt{j} \delta_{n,j-1} \right) \quad (3)$$

Since all diagonal elements are zero,  $\langle n0 | V | n0 \rangle = 0$  and there is no first order correction to the energy.

There are only two non-zero terms in the second order correction series

$$E_{n2} = \sum_{j \neq n} \frac{|\langle j0 | V | n0 \rangle|^2}{E_{n0} - E_{j0}} \quad (4)$$

That is, the terms for  $j = n \pm 1$ . Since  $E_{n0} = \hbar\omega \left( n + \frac{1}{2} \right)$ , we get

$$E_{n2} = \frac{\hbar q^2 E^2}{2m\omega} \left( \frac{-(n+1) + n}{\hbar\omega} \right) \quad (5)$$

$$= -\frac{q^2 E^2}{2m\omega^2} \quad (6)$$

$$= -\frac{q^2 E^2}{2k} \quad (7)$$

In this case, we can actually find an exact solution for the perturbed potential. The Schrödinger equation for the perturbed potential is

$$-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} + \left( \frac{1}{2}kx^2 - qEx \right) \psi = E_n \psi \quad (8)$$

Multiplying through by  $2/k$  we get

$$-\frac{2}{k} \frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} + x \left( x - \frac{qE}{k} - \frac{qE}{k} \right) \psi = \frac{2}{k} E_n \psi \quad (9)$$

We can now make the variable transformation

$$\xi \equiv x - \frac{qE}{k} \quad (10)$$

Since  $d\xi = dx$ , this transforms the Schrödinger equation to

$$-\frac{2}{k} \frac{\hbar^2}{2m} \frac{d^2\psi}{d\xi^2} + \left( \xi + \frac{qE}{k} \right) \left( \xi - \frac{qE}{k} \right) \psi = \frac{2}{k} E_n \psi \quad (11)$$

$$-\frac{2}{k} \frac{\hbar^2}{2m} \frac{d^2\psi}{d\xi^2} + \left( \xi^2 - \left( \frac{qE}{k} \right)^2 \right) \psi = \frac{2}{k} E_n \psi \quad (12)$$

$$-\frac{\hbar^2}{2m} \frac{d^2\psi}{d\xi^2} + \frac{1}{2}k\xi^2\psi = \left( E_n + \frac{q^2E^2}{2k} \right) \psi \quad (13)$$

The last line has the same form as the Schrödinger equation for the unperturbed potential except that the energies are given by  $E_n + \frac{q^2E^2}{2k}$ , which is a constant (independent of  $\xi$ ), so the solutions must be the same as for the unperturbed oscillator. That is

$$E_n + \frac{q^2E^2}{2k} = \hbar\omega \left( n + \frac{1}{2} \right) \quad (14)$$

$$E_n = \hbar\omega \left( n + \frac{1}{2} \right) - \frac{q^2E^2}{2k} \quad (15)$$

The correction to the energy is the same as the second order term found above, so in this case, second order perturbation theory gives the *exact* correction to the energy.