

HYDROGEN ATOM: POWERS OF THE MOMENTUM OPERATOR

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References: Griffiths, David J. (2005), Introduction to Quantum Mechanics, 2nd Edition; Pearson Education - Problem 6.15.

In this post, we'll derive results concerning powers of the momentum operator p when applied to the $l = 0$ states of hydrogen. The general form of the hydrogen wave function is

$$(0.1) \quad \psi_{nlm}(r, \theta, \phi) = R_{nl}(r) Y_l^m(\theta, \phi)$$

where R is the radial function and Y_l^m is a spherical harmonic. If $l = 0$, then the only possible value of m is $m = 0$ and $Y_0^0 = 1/\sqrt{4\pi}$, which is independent of θ and ϕ . In spherical coordinates, the square of the momentum operator is then

$$(0.2) \quad p^2 = -\hbar^2 \nabla^2 = -\frac{\hbar^2}{r^2} \frac{d}{dr} \left(r^2 \frac{d}{dr} \right)$$

We'd like to show that this operator is hermitian, that is, for two functions $f(r)$ and $g(r)$ that

$$(0.3) \quad \langle f | p^2 g \rangle = \langle p^2 f | g \rangle$$

We start with

$$(0.4) \quad p^2 g = -\frac{\hbar^2}{r^2} \frac{d}{dr} \left(r^2 \frac{dg}{dr} \right)$$

$$(0.5) \quad = -\frac{\hbar^2}{r^2} (2rg' + r^2 g'')$$

$$(0.6) \quad = -\hbar^2 \left(2\frac{g'}{r} + g'' \right)$$

We then get

(0.7)

$$\langle f | p^2 g \rangle = -4\pi\hbar^2 \int_0^\infty \frac{f}{r^2} (2rg' + r^2g'') r^2 dr$$

$$(0.8) \quad = -4\pi\hbar^2 \int_0^\infty f (2rg' + r^2g'') dr$$

$$(0.9) \quad = -4\pi\hbar^2 \left[r^2 fg' \Big|_0^\infty - \int_0^\infty r^2 (f'g' + fg'') dr + \int_0^\infty r^2 fg'' dr \right]$$

$$(0.10) \quad = -4\pi\hbar^2 \left[r^2 fg' \Big|_0^\infty - \int_0^\infty r^2 f'g' dr \right]$$

$$(0.11) \quad = -4\pi\hbar^2 \left[r^2 fg' \Big|_0^\infty - r^2 f'g \Big|_0^\infty + \int_0^\infty (2rf' + r^2f) g dr \right]$$

$$(0.12) \quad = -4\pi\hbar^2 \int_0^\infty (2rf' + r^2f) g dr$$

$$(0.13) \quad = \langle p^2 f | g \rangle$$

where in the second-to-last line we used the fact that all radial functions in the hydrogen atom have an $e^{-r/na}$ term multiplied by a polynomial in r . The exponential ensures the integrated terms are zero at infinity, and the r^2 factor ensures they are zero at $r = 0$. Thus p^2 is hermitian.

For p^4 , we start from 0.6 and apply 0.2:

$$(0.14) \quad p^4 = -\hbar^4 \nabla^2 \left(2\frac{g'}{r} + g'' \right)$$

For the first term, we have

$$\begin{aligned}
 (0.15) \quad \nabla^2 \frac{g'}{r} &= \nabla \cdot \left(\nabla \frac{g'}{r} \right) \\
 (0.16) \quad &= \nabla \cdot \left[g' \nabla \frac{1}{r} + \frac{1}{r} \nabla g' \right] \\
 (0.17) \quad &= g' \nabla^2 \frac{1}{r} + 2 \left(\nabla \frac{1}{r} \right) \cdot (\nabla g') + \frac{1}{r} \nabla^2 g' \\
 (0.18) \quad &= -4\pi \delta(\mathbf{r}) g' - \frac{2}{r^2} \hat{\mathbf{r}} \cdot (g'' \hat{\mathbf{r}}) + \frac{1}{r} \nabla^2 g' \\
 (0.19) \quad &= -4\pi \delta(\mathbf{r}) g' - 2 \frac{g''}{r^2} + \frac{1}{r^3} \frac{d}{dr} (r^2 g'') \\
 (0.20) \quad &= -4\pi \delta(\mathbf{r}) g' - 2 \frac{g''}{r^2} + 2 \frac{g''}{r^2} + \frac{g^{(3)}}{r} \\
 (0.21) \quad &= -4\pi \delta(\mathbf{r}) g' + \frac{g^{(3)}}{r}
 \end{aligned}$$

where the notation $g^{(i)}$ denotes the i th derivative and we've used a couple of earlier results to get the fourth line:

$$(0.22) \quad \nabla^2 \frac{1}{r} = -4\pi \delta(\mathbf{r})$$

$$(0.23) \quad \nabla \frac{1}{r} = -\frac{\hat{\mathbf{r}}}{r^2}$$

For the second term in 0.14 we have

$$(0.24) \quad \nabla^2 g'' = \frac{1}{r^2} \frac{d}{dr} (r^2 g^{(3)})$$

$$(0.25) \quad = 2 \frac{g^{(3)}}{r} + g^{(4)}$$

Inserting 0.21 and 0.25 into 0.14 we get

$$(0.26) \quad p^4 g = \hbar^4 \left(\frac{4}{r} g^{(3)} + g^{(4)} - 8\pi \delta(\mathbf{r}) g' \right)$$

Now we want to calculate $\langle f | p^4 g \rangle$ and compare it with $\langle g | p^4 f \rangle$, so we have

$$(0.27) \quad \frac{1}{\hbar^4} \langle f | p^4 g \rangle = 4\pi \int_0^\infty \left(4rf g^{(3)} + r^2 f g^{(4)} \right) - 8\pi \int \delta(\mathbf{r}) f g' d^3 \mathbf{r}$$

$$(0.28) \quad \frac{1}{\hbar^4} \langle g | p^4 f \rangle = 4\pi \int_0^\infty \left(4rg f^{(3)} + r^2 g f^{(4)} \right) - 8\pi \int \delta(\mathbf{r}) g f' d^3 \mathbf{r}$$

The aim is to integrate by parts enough times to eliminate the derivatives of g under the integral. Again, this is tedious, but we can plow onwards, or else just use some software to ease the task. Using Maple's IntegrationTools[Parts] operation, we find (after eliminating all terms evaluated at $r = \infty$ because they contain an $e^{-r/na}$ factor, and those terms containing a factor of r or r^2 evaluated at $r = 0$):

$$(0.29) \quad \int_0^\infty 4rf g^{(3)} dr = 4f(0)g'(0) - 8f'(0)g(0) - \int_0^\infty g \left(12f'' + 4rf^{(3)} \right) dr$$

$$(0.30) \quad \int_0^\infty r^2 f g^{(4)} dr = -2f(0)g'(0) + 6f'(0)g(0) + \int_0^\infty g \left(12f'' + 8rf^{(3)} + r^2 f^{(4)} \right) dr$$

Adding these together and adding on the delta function term in 0.27, we get, by comparing the result with 0.28

$$(0.31) \quad \frac{1}{4\pi\hbar^4} \langle f | p^4 g \rangle = 2f(0)g'(0) - 2f'(0)g(0) + \int_0^\infty g \left(r^2 f^{(4)} + 4rf^{(3)} \right) dr - 8\pi f(0)g'(0)$$

$$(0.32) \quad \langle f | p^4 g \rangle = 8\pi\hbar^4 \left(f(0)g'(0) - f'(0)g(0) \right) + \langle g | p^4 f \rangle + 8\pi\hbar^4 \left(g(0)f'(0) - f(0)g'(0) \right)$$

$$(0.33) \quad = \langle g | p^4 f \rangle$$

Thus p^4 is also hermitian.

[Note that this is the opposite result to that specified in Griffiths's problem 6.15, where he asks us to prove that p^4 is *not* hermitian. However, Griffiths corrects this result in his errata. Thanks to Jack Whaley-Baldwin (see comments below) for pointing this out.]

PINGBACKS

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