

WKB APPROXIMATION AND THE POWER LAW POTENTIAL

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References: Griffiths, David J. (2005), Introduction to Quantum Mechanics, 2nd Edition; Pearson Education - Problem 8.11.

We've seen that if we apply the WKB approximation to a potential well and require that the WKB wave functions match up in the region between the turning points, we get the condition

$$(0.1) \quad \int_{x_1}^{x_2} p(x) dx = \left(n - \frac{1}{2}\right) \pi \hbar$$

where x_1 is the left turning point, x_2 is the right turning point and $n = 1, 2, 3, \dots$. We've applied this result to the case of the harmonic oscillator, where the potential is $V(x) = \frac{1}{2}m\omega^2x^2$, but it's interesting to work out the more general case where the potential is any power law:

$$(0.2) \quad V(x) = \alpha |x|^v$$

where α and v are positive constants. The turning points for a particle with energy E are found by solving $E = V(x)$ so we get

$$(0.3) \quad x_{1,2} = \pm \left(\frac{E}{\alpha}\right)^{1/v}$$

Since the potential is an even function, we can write 0.1 as

$$(0.4) \quad \int_{x_1}^{x_2} p(x) dx = 2 \int_0^{(E/\alpha)^{1/v}} \sqrt{2m(E - \alpha x^v)} dx$$

Maple is unable to handle the integral as it stands, so we can help it by using the substitution

$$(0.5) \quad u = x^\nu$$

$$(0.6) \quad du = \nu x^{\nu-1} dx$$

$$(0.7) \quad = \nu u^{1-1/\nu} dx$$

$$(0.8) \quad dx = \frac{1}{\nu} u^{1/\nu-1} du$$

The turning point in the u coordinates is E/α , so the integral becomes

$$(0.9) \quad 2 \int_0^{(E/\alpha)^{1/\nu}} \sqrt{2m(E - \alpha x^\nu)} dx = \frac{2\sqrt{2m}}{\nu} \int_0^{E/\alpha} \sqrt{E - \alpha u} u^{1/\nu-1} du$$

Maple can do this integral, with the result

$$(0.10) \quad \frac{2\sqrt{2m}}{\nu} \int_0^{E/\alpha} \sqrt{E - \alpha u} u^{1/\nu-1} du = \frac{\sqrt{2\pi m}}{\nu} \frac{1}{\alpha^{1/\nu}} \frac{\Gamma(\frac{1}{\nu})}{\Gamma(\frac{1}{\nu} + \frac{3}{2})} E^{\frac{1}{\nu} + \frac{1}{2}}$$

$$(0.11) \quad = \left(n - \frac{1}{2}\right) \pi \hbar$$

Using the gamma function identity

$$(0.12) \quad \Gamma(z+1) = z\Gamma(z)$$

with $z = 1/\nu$, we can solve for E and simplify the expression slightly:

$$(0.13) \quad E_n = \left[\left(n - \frac{1}{2}\right) \pi \hbar \frac{\Gamma(\frac{1}{\nu} + \frac{3}{2})}{\Gamma(\frac{1}{\nu})} \frac{\nu \alpha^{1/\nu}}{\sqrt{2\pi m}} \right]^{\frac{2\nu}{\nu+2}}$$

$$(0.14) \quad = \left[\left(n - \frac{1}{2}\right) \pi \hbar \frac{\Gamma(\frac{1}{\nu} + \frac{3}{2})}{\Gamma(\frac{1}{\nu} + 1)} \frac{\alpha^{1/\nu}}{\sqrt{2\pi m}} \right]^{\frac{2\nu}{\nu+2}}$$

We also have

$$(0.15) \quad \left(\alpha^{1/\nu}\right)^{\frac{2\nu}{\nu+2}} = \alpha^{\frac{2}{\nu+2}} = \alpha^{\frac{\nu+2-\nu}{\nu+2}} = \alpha \cdot \alpha^{-\frac{\nu}{\nu+2}} = \alpha \cdot \left[(\alpha)^{-1/2}\right]^{\frac{2\nu}{\nu+2}}$$

so

$$(0.16) \quad E_n = \alpha \left[\left(n - \frac{1}{2} \right) \hbar \frac{\Gamma\left(\frac{1}{\nu} + \frac{3}{2}\right)}{\Gamma\left(\frac{1}{\nu} + 1\right)} \frac{\sqrt{\pi}}{\sqrt{2m\alpha}} \right]^{\frac{2\nu}{\nu+2}}$$

In the special case of the harmonic oscillator, $\nu = 2$ and $\alpha = \frac{1}{2}m\omega^2$. The values of the gamma function required here are

$$(0.17) \quad \Gamma(2) = 1$$

$$(0.18) \quad \Gamma\left(\frac{3}{2}\right) = \frac{1}{2}\sqrt{\pi}$$

so we get

$$(0.19) \quad E_n = \left(n - \frac{1}{2} \right) \hbar\omega$$