

WKB APPROXIMATION OF DOUBLE-WELL POTENTIAL: WAVE FUNCTIONS

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References: Griffiths, David J. (2005), Introduction to Quantum Mechanics, 2nd Edition; Pearson Education - Problem 8.15b-f.

Continuing our application of the WKB approximation to the problem of a double potential well, we can now look at determining the allowed energies for bound states. Since the potential $V(x)$ is even, the wave function is a linear combination of even and odd functions. We worked out the WKB wave functions for $x > 0$ and they are

$$\psi(x) \approx \begin{cases} \frac{D}{\sqrt{|p(x)|}} \left[2 \cos \theta e^{\int_{x_1}^{x_1} |p(x')| dx' / \hbar} + \sin \theta e^{-\int_{x_1}^{x_1} |p(x')| dx' / \hbar} \right] & 0 \leq x < x_1 \\ \frac{2D}{\sqrt{|p(x)|}} \sin \left[\int_x^{x_2} p(x') dx' / \hbar + \frac{\pi}{4} \right] & x_1 < x < x_2 \\ \frac{D}{\sqrt{|p(x)|}} \exp \left[-\int_{x_2}^x |p(x')| dx' / \hbar \right] & x > x_2 \end{cases} \quad (1)$$

where

$$\theta \equiv \frac{1}{\hbar} \int_{x_1}^{x_2} p(x') dx' \quad (2)$$

The odd extension of 1 is thus

$$\psi(-x) = -\psi(x) \quad (3)$$

for $x \geq 0$. For this case, we must have $\psi(0) = 0$, so we have

$$\psi(0) = \frac{D}{\sqrt{|p(0)|}} \left[2 \cos \theta e^{\int_0^{x_1} |p(x')| dx' / \hbar} + \sin \theta e^{-\int_0^{x_1} |p(x')| dx' / \hbar} \right] \quad (4)$$

$$= \frac{D}{\sqrt{|p(0)|}} e^{-\int_0^{x_1} |p(x')| dx' / \hbar} \left[2 \cos \theta e^{\phi} + \sin \theta \right] \quad (5)$$

$$\phi \equiv \frac{2}{\hbar} \int_0^{x_1} |p(x')| dx' = \frac{1}{\hbar} \int_{-x_1}^{x_1} |p(x')| dx' \quad (6)$$

where the last line follows from the fact that if $V(x)$ is even, then $p(x) = \sqrt{2m(E - V(x))}$ is also even. Setting 5 to zero gives the condition

$$\tan \theta = -2e^\phi \quad (7)$$

Since both ϕ and θ depend on $p(x)$, they both contain the energy E , so this condition imposes constraints on E .

For the even extension of the WKB function, we have

$$\psi(-x) = \psi(x) \quad (8)$$

$$\psi'(0) = 0 \quad (9)$$

The latter condition gives us

$$\psi'(x) = -\frac{1}{2} \frac{D}{|p(x)|} p'(x) \psi(x) + \frac{D}{\hbar} \sqrt{|p(x)|} \left[-2 \cos \theta e^{\int_x^{x_1} |p(x')| dx' / \hbar} + \sin \theta e^{-\int_x^{x_1} |p(x')| dx' / \hbar} \right] \quad (10)$$

Because $p(x)$ is even, $p'(0) = 0$ so

$$\psi'(0) = \frac{D}{\hbar} \sqrt{|p(0)|} \left[-2 \cos \theta e^{\int_0^{x_1} |p(x')| dx' / \hbar} + \sin \theta e^{-\int_0^{x_1} |p(x')| dx' / \hbar} \right] \quad (11)$$

$$= \frac{D}{\hbar} \sqrt{|p(0)|} e^{-\int_0^{x_1} |p(x')| dx' / \hbar} \left[-2 \cos \theta e^\phi + \sin \theta \right] \quad (12)$$

Setting this to zero gives the other energy quantization condition

$$\tan \theta = 2e^\phi \quad (13)$$

so the combined conditions are

$$\boxed{\tan \theta = \pm 2e^\phi} \quad (14)$$

If the central part of the potential (between $x = -x_1$ and $x = +x_1$) is high and/or broad then from 6, ϕ will become large, since $p(x)$ depends on $V(x)$. In that case we see from 14 that $\tan \theta \rightarrow \pm \infty$ so $\theta \rightarrow (n + \frac{1}{2}) \frac{\pi}{2}$ for some integer n . Rewriting 14 we get

$$\theta = \operatorname{arccot} \left(\pm \frac{1}{2} e^{-\phi} \right) \quad (15)$$

$$= \left(n + \frac{1}{2} \right) \pi \mp \frac{1}{2} e^{-\phi} + \mathcal{O} \left(e^{-3\phi} \right) \quad (16)$$

Now suppose we give the potential a specific formula:

$$V(x) = \begin{cases} \frac{1}{2}m\omega^2(x+a)^2 & x < 0 \\ \frac{1}{2}m\omega^2(x-a)^2 & x > 0 \end{cases} \quad (17)$$

That is, we're dealing with two linked harmonic oscillator potentials. The turning points are found from the condition $p(x) = 0$ and are, for $x > 0$:

$$x_1 = -\sqrt{\frac{2E}{m\omega^2}} + a \quad (18)$$

$$x_2 = \sqrt{\frac{2E}{m\omega^2}} + a \quad (19)$$

We can now work out θ from 2 by first working out the integral using Maple

$$\frac{1}{\hbar} \int p(x) dx = -\frac{m^{3/2}\omega^2}{2\hbar} (a-x) \left(-x^2 + 2ax + \frac{2E}{m\omega^2} - a^2 \right)^{1/2} - \frac{E}{\hbar\omega} \sin^{-1} \left(\sqrt{\frac{m\omega^2}{2E}} (a-x) \right) \quad (20)$$

The first term is zero at both x_1 and x_2 , and the argument of the arcsine is -1 at x_2 and $+1$ at x_1 so

$$\theta = \frac{1}{\hbar} \int_{x_1}^{x_2} p(x) dx \quad (21)$$

$$= \frac{\pi E}{\hbar\omega} \quad (22)$$

Using the approximation 16 we get

$$E_n^\pm \approx \left(n + \frac{1}{2} \right) \hbar\omega \mp \frac{\hbar\omega}{2\pi} e^{-\phi} \quad (23)$$

The + energy is for the even wave function, and corresponds to the minus sign on the RHS, so the energies of particles in the even state ψ_n^+ are slightly lower than those in the odd state ψ_n^- .

We can get a full time-dependent (approximate) wave function for a particle that starts out in some linear combination of ψ_n^+ and ψ_n^- by using the usual technique. Suppose we start the particle out in the following state:

$$\Psi(x, 0) = \frac{1}{\sqrt{2}} (\psi_n^+ + \psi_n^-) \quad (24)$$

This particle is entirely within the well in the region $x > 0$, since for $x < 0$, $\psi_n^+(x) = -\psi_n^-(x)$ and for $x > 0$, $\psi_n^+(x) = \psi_n^-(x)$ so the total wave

function is zero for $x < 0$. Then the full time-dependent wave function is, using 23

$$\Psi(x, t) = \frac{1}{\sqrt{2}} \left(\psi_n^+ e^{-iE_n^+ t/\hbar} + \psi_n^- e^{-iE_n^- t/\hbar} \right) \quad (25)$$

$$= \frac{1}{\sqrt{2}} e^{-i(n+\frac{1}{2})\omega t} \left(\psi_n^+ e^{i\omega t e^{-\phi}/2\pi} + \psi_n^- e^{-i\omega t e^{-\phi}/2\pi} \right) \quad (26)$$

The probability density is $|\Psi(x, t)|^2$, which can be calculated fairly easily since ψ_n^+ and ψ_n^- are both real functions, as we can see from 1. We get

$$|\Psi(x, t)|^2 = \frac{1}{2} \left[|\psi_n^+|^2 + |\psi_n^-|^2 + \psi_n^+ \psi_n^- \left(e^{i\omega t e^{-\phi}/\pi} + e^{-i\omega t e^{-\phi}/\pi} \right) \right] \quad (27)$$

$$= \frac{1}{2} \left(|\psi_n^+|^2 + |\psi_n^-|^2 \right) + \psi_n^+ \psi_n^- \cos \frac{\omega t e^{-\phi}}{\pi} \quad (28)$$

If $x > 0$, $\psi_n^+ \psi_n^- = |\psi_n^+|^2$ while if $x < 0$, $\psi_n^+ \psi_n^- = -|\psi_n^+|^2$ so as time progresses, the probability that the particle is one well or the other oscillates between the two wells, with a period given by

$$\frac{\omega e^{-\phi}}{\pi} = \frac{2\pi}{\tau} \quad (29)$$

$$\tau = \frac{2\pi^2 e^{\phi}}{\omega} \quad (30)$$

Finally, we can calculate ϕ for the double harmonic oscillator potential, using 6. We get

$$\phi = \frac{2}{\hbar} \int_0^{x_1} \sqrt{2m \left(\frac{1}{2} m \omega^2 (x-a)^2 - E \right)} dx \quad (31)$$

$$= \frac{2\sqrt{2mE}}{\hbar} \int_0^{x_1} \sqrt{\left(\frac{m\omega^2}{2E} (x-a)^2 - 1 \right)} dx \quad (32)$$

To transform this to a form that Maple can handle, we use the substitution

$$u = \sqrt{\frac{m\omega^2}{2E}} (x-a) \quad (33)$$

$$du = \sqrt{\frac{m\omega^2}{2E}} dx \quad (34)$$

The limits transform as

$$x = 0 \rightarrow u = -a\sqrt{\frac{m\omega^2}{2E}} \quad (35)$$

$$x = x_1 = -\sqrt{\frac{2E}{m\omega^2}} + a \rightarrow u = -1 \quad (36)$$

So

$$\phi = \frac{2\sqrt{2mE}}{\hbar} \sqrt{\frac{2E}{m\omega^2}} \int_{-a\sqrt{\frac{m\omega^2}{2E}}}^{-1} \sqrt{u^2 - 1} du \quad (37)$$

$$= \frac{4E}{\hbar\omega} \int_1^{a\sqrt{\frac{m\omega^2}{2E}}} \sqrt{u^2 - 1} du \quad (38)$$

where the last step uses the fact that the integrand is even. Since $V(0) = \frac{1}{2}m\omega^2 a^2 \equiv V_0$ we can write this as

$$\phi = \frac{4E}{\hbar\omega} \int_1^{a\sqrt{V_0/E}} \sqrt{u^2 - 1} du \quad (39)$$

$$= \frac{2E}{\hbar\omega} \left(u\sqrt{u^2 - 1} - \ln(u + \sqrt{u^2 - 1}) \right) \Big|_1^{a\sqrt{V_0/E}} \quad (40)$$

$$= \frac{2E}{\hbar\omega} \left[\sqrt{\frac{V_0}{E}} \sqrt{\frac{V_0}{E} - 1} - \ln \left(\sqrt{\frac{V_0}{E}} + \sqrt{\frac{V_0}{E} - 1} \right) \right] \quad (41)$$

For a high central barrier, $V_0 \gg E$ and we can approximate this by

$$\phi \approx \frac{2E}{\hbar\omega} \left[\frac{V_0}{E} - \ln \left(2\sqrt{\frac{V_0}{E}} \right) \right] \quad (42)$$

$$\approx \frac{2V_0}{\hbar\omega} = \frac{m\omega a^2}{\hbar} \quad (43)$$

where we've dropped the logarithm term as it's much smaller than V_0/E .