

DELTA FUNCTION IN TIME PERTURBATION

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References: Griffiths, David J. (2005), Introduction to Quantum Mechanics, 2nd Edition; Pearson Education - Problem 9.3.

We've seen that we can solve the Schrödinger equation with a time-dependent potential in a two-state system if we split the hamiltonian into a time-independent part H^0 and a time-dependent part H' , so that the complete hamiltonian is

$$(1) \quad H = H^0 + H'$$

The solution is

$$(2) \quad \Psi(x, t) = c_a(t) \psi_a(x) e^{-iE_a t/\hbar} + c_b(t) \psi_b(x) e^{-iE_b t/\hbar}$$

As an application of the situation where a perturbation to the hamiltonian in a two-state system is switched on at $t = 0$ and then off again at some later time $t = \tau$ we can look at a perturbation that is on only momentarily, which we can model by using a delta function:

$$(3) \quad H' = U(x) \delta(t)$$

We can think of this as the limit as $\tau \rightarrow 0$ of the earlier problem if we start with

$$(4) \quad H' = \begin{cases} \frac{U}{\tau} & 0 \leq t \leq \tau \\ 0 & \text{otherwise} \end{cases}$$

This works because the integral

$$(5) \quad \int_0^\tau \frac{dt}{\tau} = 1$$

so as $\tau \rightarrow 0$ 4 becomes 3.

If we assume that the matrix elements of H' in the two-state system satisfy

$$(6) \quad U_{aa} = U_{bb} = 0$$

$$(7) \quad U_{ab} = U_{ba}^* \equiv \alpha$$

then in the interval $0 \leq t \leq \tau$ we have

$$(8) \quad H'_{aa} = H'_{bb} = 0$$

$$(9) \quad H'_{ab} = (H'_{ba})^* = \frac{\alpha}{\tau}$$

If $c_a(-\infty) = 1$ and $c_b(-\infty) = 0$ then, since the perturbation is switched off until $t = 0$ these coefficients remain constant until that time, so we must have $c_a(0) = 1$ and $c_b(0) = 0$. This means that this system is the same as that we treated in the last post, where a perturbation is switched on at $t = 0$ and off at a later time, so the coefficients are

$$(10) \quad c_a(t) = e^{-i\omega_0 t/2} \left[\cos(Qt) + \frac{i\omega_0}{2Q} \sin(Qt) \right]$$

$$(11) \quad c_b(t) = -\frac{i|H'_{ab}|}{\hbar Q} e^{i\omega_0 t/2} \sin(Qt)$$

$$(12) \quad = -\frac{i\alpha}{\hbar Q \tau} e^{i\omega_0 t/2} \sin(Qt)$$

where

$$(13) \quad Q \equiv \frac{1}{2} \sqrt{\omega_0^2 + \frac{4|H'_{ab}|^2}{\hbar^2}}$$

$$(14) \quad = \frac{1}{2} \sqrt{\omega_0^2 + \frac{4\alpha^2}{\hbar^2 \tau^2}}$$

$$(15) \quad \omega_0 \equiv \frac{E_b - E_a}{\hbar}$$

Since the perturbation is switched off at $t = \tau$, c_a and c_b with remain constant from $t = \tau$ onwards.

Since this is a special case of the earlier problem, we have $|c_a|^2 + |c_b|^2 = 1$ here as well.

To get a delta function perturbation, we let $\tau \rightarrow 0$. In this limit, we get

$$(16) \quad Q \sim \frac{|\alpha|}{\hbar\tau}$$

and the coefficients become

$$(17) \quad c_a(\tau) = e^{-i\omega_0\tau/2} \left[\cos(Q\tau) + \frac{i\omega_0}{2Q} \sin(Q\tau) \right]$$

$$(18) \quad \sim \cos\left(\frac{|\alpha|}{\hbar}\right)$$

$$(19) \quad c_b(\tau) = -\frac{i\alpha}{\hbar Q\tau} e^{i\omega_0\tau/2} \sin(Q\tau)$$

$$(20) \quad \sim \pm i \sin\left(\frac{|\alpha|}{\hbar}\right)$$

where the \pm in the last line arises because of the quotient $\alpha/|\alpha|$. Since the system starts out in state a (because $c_a(0) = 1$), the probability that it has made a transition to state b after the perturbation is switched off is

$$(21) \quad P_{a \rightarrow b} = |c_b(\infty)|^2 = |c_b(\tau)|^2 = \sin^2\left(\frac{|\alpha|}{\hbar}\right)$$