

SPONTANEOUS EMISSION RATES FOR THE HYDROGEN ATOM

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References: Griffiths, David J. (2005), Introduction to Quantum Mechanics, 2nd Edition; Pearson Education - Problem 9.10, 9.11.

A collection of atoms an excited state will decay over time into the ground state by spontaneous emission of radiation. If the rate of decay is A (this is the rate at which a single atom decays, defined as the derivative of the transition probability with respect to time), then over a time dt a fraction $A dt$ of the number N_b of atoms still in the excited state will decay (on average). That is

$$(1) \quad dN_b = -AN_b dt$$

which is the ODE for simple exponential decay, so

$$(2) \quad N_b(t) = N_b(0) e^{-At}$$

The time constant is

$$(3) \quad \tau = \frac{1}{A}$$

and is the time required for the original population to be reduced to $1/e \approx 0.368$ of its size at $t = 0$. An often-quoted, related quantity is the *half-life* $t_{1/2}$ which is the time required for the population to be reduced to half its original size. That is

$$(4) \quad e^{-At_{1/2}} = 0.5$$
$$(5) \quad t_{1/2} = -\frac{\ln 0.5}{A}$$
$$(6) \quad = \frac{\ln 2}{A}$$
$$(7) \quad \approx 0.693 \tau$$

As an example, we can calculate the lifetimes of the four $n = 2$ states of the hydrogen atom. The transition rate is

$$(8) \quad A = \frac{\omega_0^3 |\mathbf{p}|^2}{3\epsilon_0 \pi \hbar c^3}$$

The main calculation is the evaluation of the matrix elements

$$(9) \quad \mathbf{p} = q \langle \psi_b | \mathbf{r} | \psi_a \rangle$$

where q is the electron charge and ψ_i is one of the hydrogen wave functions, given by (with $a = 4\pi\epsilon_0\hbar^2/mq^2$ as the Bohr radius)

$$(10) \quad \psi_{nlm} = \sqrt{\left(\frac{2}{na}\right)^3 \frac{(n-l-1)!}{2n[(n+l)!]^3}} e^{-r/na} \left(\frac{2r}{na}\right)^l L_{n-l-1}^{2l+1} \left(\frac{2r}{na}\right) Y_l^m(\theta, \phi)$$

where L is an associated Laguerre polynomial and Y is a spherical harmonic. Doing these integrals by hand gets quite tedious and error-prone, so it's easier to use Maple to do them. The Maple code for evaluating the Laguerre polynomial is

$$(11) \quad \text{simplify}\left(\text{LaguerreL}\left(n-l-1, 2*l+1, \frac{2*r}{n*a}\right), \text{LaguerreL}'\right)$$

The spherical harmonic can be evaluated using

$$(12) \quad \text{simplify}\left(\text{convert}(\text{SphericalY}(l, m, t, p), \text{LegendreP}), \text{LegendreP}'\right)$$

The t stands for θ and the p stands for ϕ .

With these definitions, we can get expressions for ψ_{200} , ψ_{21-1} , ψ_{210} , ψ_{211} and ψ_{100} and use the rectangular to spherical conversion formulas

$$(13) \quad x = r \sin \theta \cos \phi$$

$$(14) \quad y = r \sin \theta \sin \phi$$

$$(15) \quad z = r \cos \theta$$

With $\psi_a = \psi_{100}$ and ψ_b set to each of the four $n = 2$ wave functions in turn, we have a total of 12 integrals to work out to evaluate 9 in all cases. We can do all these in Maple, but here's the explicit formula for one of the integrals so you can see what they look like.

(16)

$$\langle \psi_{210} | z | \psi_{100} \rangle = \int_0^\infty dr \int_0^\pi d\theta \int_0^{2\pi} d\phi r^2 \sin \theta [r \cos \theta \psi_{210}^* \psi_{100}]$$

(17)

$$= \int_0^\infty dr \int_0^\pi d\theta \int_0^{2\pi} d\phi r^2 \sin \theta (r \cos \theta) \left[\frac{\sqrt{2}}{8a^{5/2}\sqrt{\pi}} r e^{-r/2a} \cos \theta \right] \left[\frac{e^{-r/a}}{a^{3/2}\sqrt{\pi}} \right]$$

(18)

$$= \frac{128\sqrt{2}}{243} a$$

The other non-zero matrix elements are

$$(19) \quad \langle \psi_{211} | x | \psi_{100} \rangle = \frac{128}{243} ai$$

$$(20) \quad \langle \psi_{211} | y | \psi_{100} \rangle = \frac{128}{243} a$$

$$(21) \quad \langle \psi_{21-1} | x | \psi_{100} \rangle = \frac{128}{243} ai$$

$$(22) \quad \langle \psi_{21-1} | y | \psi_{100} \rangle = -\frac{128}{243} a$$

From these we can evaluate \mathbf{p} for each state.

$$(23) \quad \mathbf{p}_{200} = 0$$

$$(24) \quad \mathbf{p}_{210} = \frac{128\sqrt{2}}{243} qa\hat{\mathbf{z}}$$

$$(25) \quad \mathbf{p}_{211} = \frac{128}{243} qai\hat{\mathbf{x}} + \frac{128}{243} qa\hat{\mathbf{y}}$$

$$(26) \quad \mathbf{p}_{21-1} = \frac{128}{243} qai\hat{\mathbf{x}} - \frac{128}{243} qa\hat{\mathbf{y}}$$

From this we get $|\mathbf{p}|^2 = |\mathbf{p}_x|^2 + |\mathbf{p}_y|^2 + |\mathbf{p}_z|^2$.

$$(27) \quad |\mathbf{p}_{200}|^2 = 0$$

$$(28) \quad |\mathbf{p}_{210}|^2 = 2 \left(\frac{128}{243} aq \right)^2$$

$$(29) \quad |\mathbf{p}_{211}|^2 = 2 \left(\frac{128}{243} aq \right)^2$$

$$(30) \quad |\mathbf{p}_{21-1}|^2 = 2 \left(\frac{128}{243} aq \right)^2$$

To evaluate the transition rate 8 we need ω_0 which is the energy of the emitted photon, which we can get from the difference in energy between the $n = 2$ and $n = 1$ levels in the Bohr formula

$$(31) \quad E_n = -\frac{1}{n^2} \frac{mq^4}{2\hbar^2(4\pi\epsilon_0)^2}$$

We have

$$(32) \quad \omega_0 = \frac{E_2 - E_1}{\hbar}$$

$$(33) \quad = \frac{3mq^4}{8\hbar^3(4\pi\epsilon_0)^2}$$

$$(34) \quad = 1.55 \times 10^{16} \text{ s}^{-1}$$

Plugging these results into 8 we get for all states except ψ_{200} :

$$(35) \quad \tau = \frac{1}{A} = 1.595 \times 10^{-9} \text{ s}$$

For ψ_{200} , the lifetime is $1/A = \infty$ so this state is stable.

PINGBACKS

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