

INTEGRAL FORM OF THE SCHRÖDINGER EQUATION

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Reference: Griffiths, David J. (2005), Introduction to Quantum Mechanics, 2nd Edition; Pearson Education - Problem 11.8.

As a prelude to the Born approximation in quantum scattering, we need to look at the integral form of the time-independent Schrödinger equation. The equation in its original differential equation form is

$$(0.1) \quad -\frac{\hbar^2}{2m}\nabla^2\psi + V\psi = E\psi$$

which can be written as

$$(0.2) \quad (\nabla^2 + k^2)\psi = Q$$

$$(0.3) \quad k \equiv \frac{\sqrt{2mE}}{\hbar}$$

$$(0.4) \quad Q \equiv \frac{2m}{\hbar^2}V\psi$$

To convert this to an integral equation, we need to define a *Green's function* $G(\mathbf{r})$ which satisfies the differential equation

$$(0.5) \quad (\nabla^2 + k^2)G(\mathbf{r}) = \delta^3(\mathbf{r})$$

Using this function we can write ψ as an integral equation

$$(0.6) \quad \psi(\mathbf{r}) = \int G(\mathbf{r} - \mathbf{r}_0)Q(\mathbf{r}_0)d^3\mathbf{r}_0$$

We can show this works by plugging in G from 0.5:

$$(0.7) \quad (\nabla^2 + k^2)\psi(\mathbf{r}) = \int (\nabla^2 + k^2)G(\mathbf{r} - \mathbf{r}_0)Q(\mathbf{r}_0)d^3\mathbf{r}_0$$

$$(0.8) \quad = \int \delta(\mathbf{r} - \mathbf{r}_0)Q(\mathbf{r}_0)d^3\mathbf{r}_0$$

$$(0.9) \quad = Q(\mathbf{r})$$

which gives us back 0.2.

This isn't a solution of the Schrödinger equation, of course, because Q contains ψ , so we'd need to actually know ψ in advance in order to work out the integral with the Green's function. Rather, it's just a different way of writing the Schrödinger equation which proves useful in scattering theory.

Because 0.5 doesn't depend on the potential V , we can work out the Green's function which is valid for every potential. The process is rather involved, but Griffiths goes through the details in section 11.4.1, so I won't reproduce them here, apart from noting that the solution uses what is, to me, one of the most beautiful theorems in mathematics: Cauchy's theorem on contour integration. Maybe I'll return to it later.

Anyway, the Green's function turns out to be

$$(0.10) \quad G(\mathbf{r}) = -\frac{e^{ikr}}{4\pi r}$$

We can verify this is in fact a solution by plugging it back into 0.5. We need the Laplacian of G which we can get by calculating the divergence of the gradient. Taking the gradient first, we use the product rule for gradients:

$$(0.11) \quad \nabla(fg) = f\nabla g + g\nabla f$$

We get

$$(0.12) \quad \nabla G = -\frac{1}{4\pi} \left(\frac{1}{r} \nabla e^{ikr} + e^{ikr} \nabla \frac{1}{r} \right)$$

To calculate the divergence of the gradient, we use the identity for the divergence of the product of a scalar and a vector:

$$(0.13) \quad \nabla \cdot (f\mathbf{A}) = \mathbf{A} \cdot \nabla f + f\nabla \cdot \mathbf{A}$$

We therefore have

$$\nabla \cdot (\nabla G) = -\frac{1}{4\pi} \left[\left(\nabla \frac{1}{r} \right) \cdot \left(\nabla e^{ikr} \right) + \frac{1}{r} \nabla^2 e^{ikr} + \left(\nabla e^{ikr} \right) \cdot \left(\nabla \frac{1}{r} \right) + e^{ikr} \nabla^2 \frac{1}{r} \right]$$

The last term turns out to be a delta function:

$$(0.15) \quad \nabla^2 \frac{1}{r} = \nabla \cdot \left(\nabla \frac{1}{r} \right) = -\nabla \cdot \left(\frac{\hat{\mathbf{r}}}{r^2} \right) = -4\pi \delta^3(\mathbf{r})$$

To work out the second term, we use the formula for the Laplacian in spherical coordinates, for a function that depends only on r :

$$(0.16) \quad \nabla f(r) = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial f}{\partial r} \right)$$

We get

$$(0.17) \quad \nabla^2 e^{ikr} = \frac{1}{r^2} \frac{\partial}{\partial r} \left(ikr^2 e^{ikr} \right)$$

$$(0.18) \quad = \frac{2ik}{r} e^{ikr} - k^2 e^{ikr}$$

Putting this back into 0.14 we get

$$(0.19) \quad \nabla^2 G = -\frac{1}{4\pi} \left[2 \left(\nabla \frac{1}{r} \right) \cdot \left(\nabla e^{ikr} \right) + \frac{2ik}{r^2} e^{ikr} - \frac{k^2}{r} e^{ikr} - 4\pi \delta^3(\mathbf{r}) e^{ikr} \right]$$

$$(0.20) \quad = -\frac{1}{4\pi} \left[-\frac{2ik}{r^2} e^{ikr} + \frac{2ik}{r^2} e^{ikr} - \frac{k^2}{r} e^{ikr} - 4\pi \delta^3(\mathbf{r}) e^{ikr} \right]$$

$$(0.21) \quad = -\frac{1}{4\pi} \left[-\frac{k^2}{r} e^{ikr} - 4\pi \delta^3(\mathbf{r}) \right]$$

$$(0.22) \quad = \frac{k^2}{4\pi r} e^{ikr} + \delta^3(\mathbf{r})$$

$$(0.23) \quad = -k^2 G(\mathbf{r}) + \delta^3(\mathbf{r})$$

$$(0.24)$$

$$(\nabla^2 + k^2) G(\mathbf{r}) = \delta^3(\mathbf{r})$$

where we dropped the e^{ikr} from the last term in the third line since the delta function is zero except when $\mathbf{r} = 0$.

Using this Green's function, the integral form of the Schrödinger equation is

$$(0.25) \quad \psi(\mathbf{r}) = \psi_0(\mathbf{r}) - \frac{m}{2\pi\hbar^2} \int \frac{e^{ik|\mathbf{r}-\mathbf{r}_0|}}{|\mathbf{r}-\mathbf{r}_0|} V(\mathbf{r}_0) \psi(\mathbf{r}_0) d^3\mathbf{r}_0$$

where ψ_0 is a solution of the free particle Schrödinger equation

$$(0.26) \quad (\nabla^2 + k^2) \psi_0(\mathbf{r}) = 0$$

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