

QUANTUM ELECTRODYNAMICS: NO FIRST-ORDER PROCESSES ARE POSSIBLE

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Post date: 12 Oct 2018.

References: Amitabha Lahiri & P. B. Pal, *A First Book of Quantum Field Theory*, Second Edition (Alpha Science International, 2004) - Chapter 9, Exercise 9.3.

The Lagrangian for a system consisting of the electromagnetic field and fermion fields is given by

$$\mathcal{L} = \bar{\psi} (i\cancel{\partial} - m) \psi - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} - eQ \bar{\psi} \gamma^\mu \psi A_\mu \quad (1)$$

The first term is the Lagrangian for the Dirac field and the second term is that for a free electromagnetic field. The last term involves both the Dirac spinors ψ and the electromagnetic field A_μ , so represents an interaction between these two types of field. For the special case of the electron field, the charge is $-e$ so $Q = -1$ and the interaction Lagrangian is

$$\mathcal{L}_{\text{int}} = e \bar{\psi} \gamma^\mu \psi A_\mu \quad (2)$$

The corresponding interaction Hamiltonian is obtained from the usual formula

$$\mathcal{H}_I = \sum \pi \dot{\phi} - \mathcal{L}_{\text{int}} \quad (3)$$

$$= -e \bar{\psi} \gamma^\mu \psi A_\mu \quad (4)$$

Imposing the usual requirement of normal ordering we have

$$:\mathcal{H}_I: = -e : \bar{\psi} \gamma^\mu \psi A_\mu : \quad (5)$$

To use this hamiltonian to calculate S-matrix elements, we look at first order terms first, and then higher order terms if necessary. For anything beyond first order, we need Wick's theorem to convert a time-ordered product into a sum over normal ordered products.

For first order, however, the S-matrix element is

$$S_{fi}^{(1)} = \int d^4x \langle f | \mathcal{H}_I | i \rangle \quad (6)$$

where i and f represent the initial and final states.

The various fields in the interaction Hamiltonian 5 can be written as the sum of a creation and annihilation parts. For the photon field we have

$$A^\mu(x) = \int \frac{d^3k}{\sqrt{(2\pi)^3 2\omega_{\mathbf{k}}}} \sum_{r=0}^3 \left[\epsilon_r^\mu(k) a_r(k) e^{-ik \cdot x} + \epsilon_r^{*\mu}(k) a_r^\dagger(k) e^{ik \cdot x} \right] \quad (7)$$

which we can write as

$$A^\mu(x) = A_+^\mu(x) + A_-^\mu(x) \quad (8)$$

where A_+^μ contains the annihilation operators and A_-^μ the creation operators:

$$A_+^\mu(x) = \int \frac{d^3k}{\sqrt{(2\pi)^3 2\omega_{\mathbf{k}}}} \sum_{r=0}^3 \epsilon_r^\mu(k) a_r(k) e^{-ik \cdot x} \quad (9)$$

$$A_-^\mu(x) = \int \frac{d^3k}{\sqrt{(2\pi)^3 2\omega_{\mathbf{k}}}} \sum_{r=0}^3 \epsilon_r^{*\mu}(k) a_r^\dagger(k) e^{ik \cdot x} \quad (10)$$

For the Dirac field, the Fourier decompositions are

$$\psi(x) = \int \frac{d^3p}{\sqrt{2(2\pi)^3 E_p}} \sum_{s=1,2} \left(f_s(\mathbf{p}) u_s(\mathbf{p}) e^{-ip \cdot x} + \hat{f}_s^\dagger(\mathbf{p}) v_s(\mathbf{p}) e^{ip \cdot x} \right) \quad (11)$$

$$\bar{\psi}(x) = \int \frac{d^3p}{\sqrt{2(2\pi)^3 E_p}} \sum_{s=1,2} \left(f_s^\dagger(\mathbf{p}) \bar{u}_s(\mathbf{p}) e^{ip \cdot x} + \hat{f}_s(\mathbf{p}) \bar{v}_s(\mathbf{p}) e^{-ip \cdot x} \right) \quad (12)$$

We can split each of these into two parts to get

$$\psi_+(x) = \int \frac{d^3p}{\sqrt{2(2\pi)^3 E_p}} \sum_{s=1,2} f_s(\mathbf{p}) u_s(\mathbf{p}) e^{-ip \cdot x} \quad (13)$$

$$\psi_-(x) = \int \frac{d^3p}{\sqrt{2(2\pi)^3 E_p}} \sum_{s=1,2} \hat{f}_s^\dagger(\mathbf{p}) v_s(\mathbf{p}) e^{ip \cdot x} \quad (14)$$

$$\bar{\psi}_+(x) = \int \frac{d^3p}{\sqrt{2(2\pi)^3 E_p}} \sum_{s=1,2} \hat{f}_s(\mathbf{p}) \bar{v}_s(\mathbf{p}) e^{-ip \cdot x} \quad (15)$$

$$\bar{\psi}_-(x) = \int \frac{d^3p}{\sqrt{2(2\pi)^3 E_p}} \sum_{s=1,2} f_s^\dagger(\mathbf{p}) \bar{u}_s(\mathbf{p}) e^{ip \cdot x} \quad (16)$$

Here \hat{f} is a positron operator and f is an electron operator. Thus we see that ψ_+ annihilates an electron, ψ_- creates a positron, $\bar{\psi}_+$ annihilates a positron and $\bar{\psi}_-$ creates an electron.

Given all this, what interactions are possible in first order? We see from 5 that, because each of the fields can be written as the sum of two terms, when we expand out 5, we will get a sum of $2 \times 2 \times 2 = 8$ terms. Some terms are clearly impossible from the outset. For example, the term involving $\bar{\psi}_+ \psi_+ A_+$ annihilates an electron, positron and photon and creates nothing in return. Thus it corresponds to a process in which all three particles come together at the same spacetime point and mutually annihilate each other, leaving nothing. Such a process clearly cannot conserve energy or momentum, so is impossible.

Other processes might appear to be possible, however. We could, for example, envision a process in which an electron and positron are annihilated and a single photon created as a result. Such a process would be given by the term $\bar{\psi}_+ \psi_+ A_-$.

However, it turns out that *none* of the processes at first order are physically possible. We can see why by considering conservation of momentum. Suppose the two fermions have 4-momenta p and p' , and the photon has 4-momentum k . Then, for any of the first order processes we must have

$$p' = \pm p \pm k \quad (17)$$

for some choice of the signs. For example, for the process $e^- + e^+ \rightarrow \gamma$ mentioned above, we must have

$$k = p + p' \quad (18)$$

Returning to 17, we can square both sides to get

$$p'^2 = p^2 + k^2 \pm 2p \cdot k \quad (19)$$

For a photon $k^2 = 0$ and, if both fermions have the same mass (as for electrons and positrons), $p'^2 = p^2 = m^2$ so we must have

$$p \cdot k = 0 \quad (20)$$

Remembering that $p \cdot k$ is the 4-dimensional scalar product, we have

$$p \cdot k = p_0 k_0 - \mathbf{p} \cdot \mathbf{k} \quad (21)$$

$$= E\omega - \mathbf{p}\mathbf{k} \cos \theta \quad (22)$$

where $E = p_0$ is the fermion energy, $\omega = k_0$ is the photon energy and $\mathbf{p} \equiv |\mathbf{p}|$ and $\mathbf{k} \equiv |\mathbf{k}|$. Therefore

$$\cos \theta = \frac{E\omega}{\mathbf{p}\mathbf{k}} \quad (23)$$

However, for the photon $\omega = \mathbf{k}$ and for the fermion $E = \sqrt{\mathbf{p}^2 + m^2}$ so we have

$$\cos \theta = \frac{\sqrt{\mathbf{p}^2 + m^2}}{\mathbf{p}} \quad (24)$$

$$= \sqrt{1 + \frac{m^2}{\mathbf{p}^2}} > 1 \quad (25)$$

Thus we cannot satisfy this condition for any real angle θ , so for any first-order process, momentum cannot be conserved. Thus the lowest order at which an interaction can take place in QED is at second order.

PINGBACKS

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