

## TIME REVERSAL IN THE DIRAC EQUATION

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Post date: 6 Jan 2019.

References: Amitabha Lahiri & P. B. Pal, *A First Book of Quantum Field Theory*, Second Edition (Alpha Science International, 2004) - Chapter 10, Exercise 10.16.

We've had a look at time reversal earlier, where we saw that it is an antilinear operator. A system is time reversal invariant if it remains the same when we keep the spatial coordinates the same and reverse the direction of time, so that  $t \rightarrow -t$ . The coordinates in the time reversed system are thus

$$x_T \equiv (-t, \mathbf{x}) = -\tilde{x} \quad (1)$$

where  $\tilde{x}$  is defined by this equation.

The original Dirac Lagrangian is

$$\mathcal{L} = \bar{\psi} (i\cancel{\partial} - m) \psi \quad (2)$$

in which we can separate the time and space coordinates to give

$$\mathcal{L} = \psi^\dagger(x) (i\partial_t + i\gamma_0 \boldsymbol{\gamma} \cdot \nabla_x - m\gamma_0) \psi(x) \quad (3)$$

In section 10.4, L&P derive the time-reversed Dirac Lagrangian.

$$\mathcal{T} \mathcal{L}(x) \mathcal{T}^{-1} = \psi_T^\dagger(x) (-i\partial_t - i\gamma_0^* \boldsymbol{\gamma}^* \cdot \nabla_x - m\gamma_0^*) \psi_T(x) \quad (4)$$

$$= \psi_T^\dagger(x) \left( -i\partial_t + i\gamma_0^\top \boldsymbol{\gamma}^\top \cdot \nabla_x - m\gamma_0^\top \right) \psi_T(x) \quad (5)$$

$$\psi^\dagger(-\tilde{x}) \mathbb{T}^\dagger \left( i\partial_{-t} + i\gamma_0^\top \boldsymbol{\gamma}^\top \cdot \nabla_x - m\gamma_0^\top \right) \mathbb{T} \psi(-\tilde{x}) \quad (6)$$

In the first line, we use the antilinearity of time reversal, so that all values are replaced by their complex conjugates. We also note that for a square matrix such as  $\gamma_\mu$ , taking the hermitian conjugate followed by the transpose (which is the meaning L&P give to the superscript \* when applied to a matrix) is equivalent to just taking the complex conjugate, since the hermitian conjugate is complex conjugate followed by the transpose, so the two transpose operations cancel out. Thus  $\gamma_0^* = \gamma_0^{\dagger\top}$  is also just the complex conjugate of  $\gamma_0$ .

To get the second line, we use the hermiticity property of  $\gamma^\mu$  (L&P equation 4.18) and anticommutators:

$$\gamma^{\mu\dagger} = \gamma^0 \gamma^\mu \gamma^0 \quad (7)$$

Thus

$$\gamma_0^* = \gamma_0^{\dagger T} = \gamma_0^T \quad (8)$$

$$\gamma_i^* = \gamma_i^{\dagger T} = (\gamma_0 \gamma_i \gamma_0)^T = -(\gamma_0 \gamma_0 \gamma_i)^T = -\gamma_i^T \quad (9)$$

In the last line, the  $4 \times 4$  matrix  $T$  is used to generate the time reversed Dirac field according to

$$\psi_T(x) = T\psi(-\tilde{x}) \quad (10)$$

Time reversal invariance of the Lagrangian requires that  $\mathcal{L}$  be the same as the original Lagrangian with  $x$  replaced by  $-\tilde{x}$ , so that

$$\mathcal{T}\mathcal{L}(x)\mathcal{T}^{-1} = \mathcal{L}(-\tilde{x}) \quad (11)$$

This gives three conditions on the matrix  $T$ :

$$T^\dagger T = 1 \quad (12)$$

$$T^\dagger \gamma_0^T \gamma_i^T T = \gamma_0 \gamma_i \quad (13)$$

$$T^\dagger \gamma_0^T T = \gamma_0 \quad (14)$$

Remember that  $T$  as a superscript indicates a matrix transpose, and is not the same as the  $4 \times 4$  matrix  $T$ .

Using the relation

$$C^{-1} \gamma_\mu C = -\gamma_\mu^T \quad (15)$$

we can get rid of the matrix transposes in these conditions. From 14 we have (since  $T$  is antiunitary, we have  $T^{-1} = T^\dagger$ ):

$$T^\dagger \gamma_0^T T = T^{-1} \gamma_0^T T \quad (16)$$

$$= -T^{-1} C^{-1} \gamma_0 C T \quad (17)$$

$$= \gamma_0 \quad (18)$$

or

$$CT\gamma_0 = -\gamma_0 CT \quad (19)$$

which gives the anticommutator

$$\{CT, \gamma_0\} = 0 \quad (20)$$

Now, using 13 we have

$$T\gamma_0\gamma_iT^{-1} = \gamma_0^T\gamma_i^T \quad (21)$$

$$= C^{-1}\gamma_0CC^{-1}\gamma_iC \quad (22)$$

$$= C^{-1}\gamma_0\gamma_iC \quad (23)$$

This gives

$$CT\gamma_0\gamma_i = \gamma_0\gamma_iCT \quad (24)$$

$$-\gamma_0CT\gamma_i = \gamma_0\gamma_iCT \quad (25)$$

We now multiply on the left by  $\gamma_0$  to get

$$-CT\gamma_i = \gamma_iCT \quad (26)$$

$$\{CT, \gamma_i\} = 0 \quad (27)$$

Thus CT anticommutes with all the gamma matrices:

$$\{CT, \gamma_\mu\} = 0 \quad (28)$$

As the matrix  $\gamma_5$  also anticommutes with all the gamma matrices, this makes CT a multiple of  $\gamma_5$ , which is written as

$$T = \eta_T C^{-1} \gamma_5 \quad (29)$$

where  $\eta_T$  is the time reversal phase factor. Referring back to 10, we thus have for a Dirac field:

$$\psi_T(x) = \eta_T C^{-1} \gamma_5 \psi(-\tilde{x}) \quad (30)$$

We can now see what happens if we have two different representations of the  $\gamma_\mu$  related by the unitary transformation

$$\tilde{\gamma}_\mu = U\gamma_\mu U^\dagger \quad (31)$$

The time reversed Dirac Lagrangian using the new gamma matrices is obtained from 5, so we have

$$\mathcal{T}\mathcal{L}(x)\mathcal{T}^{-1} = \tilde{\psi}_T^\dagger(x) \left( -i\partial_t + i\tilde{\gamma}_0^\top \tilde{\gamma}^\top \cdot \nabla_x - m\tilde{\gamma}_0^\top \right) \tilde{\psi}_T(x) \quad (32)$$

$$= \tilde{\psi}_T^\dagger(x) \left( -i\partial_t + i \left( U\gamma_0 U^\dagger \right)^\top \left( U\gamma U^\dagger \right)^\top \cdot \nabla_x - m \left( U\gamma_0 U^\dagger \right)^\top \right) \tilde{\psi}_T(x) \quad (33)$$

$$= \tilde{\psi}_T^\dagger(x) \left( -i\partial_t + iU^{\dagger\top} \gamma_0^\top U^\top U^{\dagger\top} \gamma^\top U^\top \cdot \nabla_x - mU^{\dagger\top} \gamma_0^\top U^\top \right) \tilde{\psi}_T(x) \quad (34)$$

$$= \tilde{\psi}_T^\dagger(x) \left( -U^{\dagger\top} U^\top i\partial_t + iU^{\dagger\top} \gamma_0^\top \gamma^\top U^\top \cdot \nabla_x - mU^{\dagger\top} \gamma_0^\top U^\top \right) \tilde{\psi}_T(x) \quad (35)$$

where we've inserted the identity matrix in the form  $U^{\dagger\top} U^\top$  in the first term in the last line.

This matches the original time-reversed Lagrangian 5 if

$$\psi_T(x) = U^\top \tilde{\psi}_T(x) \quad (36)$$

or, multiplying both sides by  $U^{\dagger\top} = U^*$

$$\tilde{\psi}_T(x) = U^* \psi_T(x) \quad (37)$$

#### PINGBACKS

Pingback: Time reversal in fermion interactions