

ENERGY-MOMENTUM TENSOR: COMMUTATOR OF FIELD WITH CONSERVED CHARGE

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Reference: Tom Lancaster and Stephen J. Blundell, *Quantum Field Theory for the Gifted Amateur*, (Oxford University Press, 2014), Problem 10.1.

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We've seen before the derivation of the energy-momentum tensor from Noether's theorem. L&B derive the tensor using a slightly different argument from those we've seen before, but I won't copy it all out again here, as the result is the same. In L&B's notation, we have

$$T^{\mu\nu} = \Pi^\mu \partial^\nu \phi - g^{\mu\nu} \mathcal{L} \quad (1)$$

where \mathcal{L} is the Lagrangian density, ϕ is the field and the canonical momentum density is

$$\Pi^\mu \equiv \frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi)} \quad (2)$$

The conserved charges from Noether's theorem are

$$P^\alpha = \int d^3x T^{0\alpha} \quad (3)$$

In their problem 10.1, L&B ask us to show

$$[\phi(x), P^\alpha] = i\partial^\alpha \phi(x) \quad (4)$$

It's not clear to me how we can do this using only material that has been presented in L&B's book up to this point. As far as I can tell, the presentation in Chapter 10 refers to classical field theory, in which commutators don't have any meaning. We *have* seen a derivation of 4 before (here) but this relies on Fourier expansions of ϕ and Π in terms of creation and annihilation operators, which haven't yet been introduced in L&B.

However, if we assume the usual quantum commutation relation between the field and the conjugate momentum $\pi(x)$

$$\pi(x) = \frac{\partial \mathcal{L}}{\partial \dot{\phi}} = \Pi^0(x) \quad (5)$$

which is

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$$[\phi(x), \pi(y)] = i\delta^{(3)}(\mathbf{x} - \mathbf{y}) \quad (6)$$

then we can write 3 as (changing the variable of integration from x to y for convenience later):

$$P^\alpha = \int d^3y (\pi(y) \partial^\alpha \phi(y) - g^{\mu\nu} \mathcal{L}(y)) \quad (7)$$

and thus get (assuming that the field commutes with the Lagrangian density \mathcal{L}):

$$[\phi(x), P^\alpha] = \int d^3y [\phi(x), \pi(y)] \partial^\alpha \phi(y) \quad (8)$$

$$= \int d^3y i\delta^{(3)}(\mathbf{x} - \mathbf{y}) \partial^\alpha \phi(y) \quad (9)$$

$$= \partial^\alpha \phi(x) \quad (10)$$

We can take $\phi(x)$ inside the integral since its spacetime variable x is not integrated over.

PINGBACKS

Pingback: [Energy-momentum tensor for Klein-Gordon Lagrangian](#)

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