

NOETHER CURRENT FOR SYSTEM OF MANY SCALAR FIELDS

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Reference: Tom Lancaster and Stephen J. Blundell, *Quantum Field Theory for the Gifted Amateur*, (Oxford University Press, 2014), Problem 10.2.

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L&B derive the Noether current J^μ for a Lagrangian that depends on a single scalar field $\phi(x)$ in their Chapter 10. Problem 10.2 asks us to go through the same derivation for a Lagrangian that depends on N scalar fields ϕ_1, \dots, ϕ_N so that

$$\mathcal{L} = \mathcal{L}(\phi_1, \dots, \phi_N; \partial_\mu \phi_1, \dots, \partial_\mu \phi_N; x^\mu) \quad (1)$$

The derivation is very similar to that given in the book. The main difference is that we need to use the chain rule for derivatives in a few places. We'll run through the derivation here. We get (where repeated indices are summed, with a being summed from 1 to N):

$$\delta \mathcal{L} = \frac{\partial \mathcal{L}}{\partial \phi^a} \delta \phi^a + \frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi^a)} \delta (\partial_\mu \phi^a) \quad (2)$$

$$= \frac{\partial \mathcal{L}}{\partial \phi^a} \delta \phi^a + \Pi_a^\mu \delta (\partial_\mu \phi^a) \quad (3)$$

where

$$\Pi_a^\mu \equiv \frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi^a)} \quad (4)$$

We now use the product rule to write

$$\partial_\mu (\Pi_a^\mu \delta \phi^a) = \Pi_a^\mu \delta (\partial_\mu \phi^a) + \partial_\mu (\Pi_a^\mu) \delta \phi^a \quad (5)$$

Using this, we rewrite 3 as

$$\delta \mathcal{L} = \frac{\partial \mathcal{L}}{\partial \phi^a} \delta \phi^a + \partial_\mu (\Pi_a^\mu \delta \phi^a) - \partial_\mu (\Pi_a^\mu) \delta \phi^a \quad (6)$$

$$= \left(\frac{\partial \mathcal{L}}{\partial \phi^a} - \partial_\mu (\Pi_a^\mu) \right) \delta \phi^a + \partial_\mu (\Pi_a^\mu \delta \phi^a) \quad (7)$$

We now require the system to satisfy the Euler-Lagrange equations of motion, which state

$$\frac{\partial \mathcal{L}}{\partial \phi^a} - \partial_\mu (\Pi_a^\mu) = 0 \quad (8)$$

so that we have

$$\delta \mathcal{L} = \partial_\mu (\Pi_a^\mu \delta \phi^a) \quad (9)$$

We now consider a translation along a spacetime vector b^μ by an amount λb^μ , where λ is some parameter that can be varied to vary the size of the translation. Under this transformation the N fields will transform as

$$\phi_a(x^\mu) \rightarrow \phi_a(x^\mu + \lambda b^\mu) \quad (10)$$

We then define the quantity

$$D\phi^a \equiv \left. \frac{\partial \phi^a}{\partial \lambda} \right|_{\lambda=0} = b^\mu \partial_\mu \phi^a \quad (11)$$

so that

$$\delta \phi^a = D\phi^a \delta \lambda \quad (12)$$

The goal is that the action

$$S = \int d^4x \mathcal{L} \quad (13)$$

remain stationary under the symmetry transformation (since we're assuming that the system is invariant under this symmetry transformation). One way of ensuring this is to require $\delta \mathcal{L} = 0$, but since the change in action is

$$\delta S = \int d^4x \delta \mathcal{L} \quad (14)$$

where the integral extends over all spacetime, we can allow $\delta \mathcal{L}$ to be some total divergence, provided that we can apply the divergence theorem to convert the spacetime integral to a surface integral, which as usual vanishes as we go to infinity. That is, we can let

$$\delta \mathcal{L} = (\partial_\mu W^\mu) \delta \lambda \quad (15)$$

for some function $W^\mu(x)$. For a given W^μ , the effect this has on the action integral 13 will be a change in the total action of some constant value (which could be zero), so if $\delta S = 0$ when $\delta \mathcal{L} = 0$ it will still be zero if we add on the term $(\partial_\mu W^\mu) \delta \lambda$ since we're not varying any of the fields in doing this.

Equating this with 9 we get the condition

$$\partial_\mu (\Pi_a^\mu \delta\phi^a - (\partial_\mu W^\mu) \delta\lambda) = \partial_\mu (\Pi_a^\mu D\phi^a - \partial_\mu W^\mu) \delta\lambda \quad (16)$$

$$= 0 \quad (17)$$

which gives the Noether current

$$J_N^\mu(x) = \Pi_a^\mu D\phi^a - \partial_\mu W^\mu \quad (18)$$

which satisfies the conservation rule

$$\partial_\mu J_N^\mu(x) = 0 \quad (19)$$