

COMMUTATOR OF REAL SCALAR FIELD WITH MOMENTUM

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Reference: Tom Lancaster and Stephen J. Blundell, *Quantum Field Theory for the Gifted Amateur*, (Oxford University Press, 2014), Problem 11.2.

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The real scalar field $\phi(x)$ can be written as a Fourier expansion

$$\phi(x) = \frac{1}{(2\pi)^{3/2}} \int \frac{d^3p}{\sqrt{2E_p}} \left(a(p) e^{-ip \cdot x} + a^\dagger(p) e^{ip \cdot x} \right) \quad (1)$$

where the energy is

$$E_{\mathbf{p}} = +\sqrt{\mathbf{p}^2 + m^2} \quad (2)$$

The conjugate momentum density is defined as

$$\Pi^0(x) = \frac{\partial \phi(x)}{\partial t} \quad (3)$$

$$= \frac{i}{(2\pi)^{3/2}} \int \frac{d^3p E_p}{\sqrt{2E_p}} \left(-a(p) e^{-ip \cdot x} + a^\dagger(p) e^{ip \cdot x} \right) \quad (4)$$

$$= \frac{i}{\sqrt{2}(2\pi)^{3/2}} d^3p \sqrt{E_p} \left(a^\dagger(p) e^{ip \cdot x} - a(p) e^{-ip \cdot x} \right) \quad (5)$$

The commutators of the creation and annihilation operators are then

$$\left[a(p), a^\dagger(q) \right] = \delta^{(3)}(\mathbf{p} - \mathbf{q}) \quad (6)$$

$$\left[a(p), a(q) \right] = \left[a^\dagger(p), a^\dagger(q) \right] = 0 \quad (7)$$

From these we can work out the following commutator at equal times, that is, when $x^0 = y^0$:

$$[\phi(x), \Pi^0(y)] = \frac{i}{2(2\pi)^3} \int \frac{d^3p d^3q}{\sqrt{E_p}} \sqrt{E_q} [a(p) e^{-i\mathbf{p}\cdot\mathbf{x}} + a^\dagger(p) e^{i\mathbf{p}\cdot\mathbf{x}}, a^\dagger(q) e^{i\mathbf{q}\cdot\mathbf{y}} - a(q) e^{-i\mathbf{q}\cdot\mathbf{y}}] \quad (8)$$

$$= \frac{i}{2(2\pi)^3} \int \frac{d^3p d^3q}{\sqrt{E_p}} \sqrt{E_q} \left\{ \delta^{(3)}(\mathbf{p} - \mathbf{q}) e^{-i(\mathbf{p}\cdot\mathbf{x} - \mathbf{q}\cdot\mathbf{y})} + \delta^{(3)}(\mathbf{p} - \mathbf{q}) e^{i(\mathbf{p}\cdot\mathbf{x} - \mathbf{q}\cdot\mathbf{y})} \right\} \quad (9)$$

$$= \frac{i}{2(2\pi)^3} \int d^3p \left(e^{-i\mathbf{p}\cdot(\mathbf{x}-\mathbf{y})} + e^{i\mathbf{p}\cdot(\mathbf{x}-\mathbf{y})} \right) \quad (10)$$

Because we're integrating over all values of \mathbf{p} , we can swap $\mathbf{p} \rightarrow -\mathbf{p}$ in the second term. This introduces a sign change in d^3p and another sign change from the reversal of the integration limits, so the two signs cancel each other, giving

$$[\phi(x), \Pi^0(y)] = \frac{2i}{2(2\pi)^3} \int d^3p e^{-i\mathbf{p}\cdot(\mathbf{x}-\mathbf{y})} \quad (11)$$

$$= i\delta^{(3)}(\mathbf{x} - \mathbf{y}) \quad (12)$$

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