COMPLEX SCALAR FIELD: INTERNAL TRANSFORMATION

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Reference: Tom Lancaster and Stephen J. Blundell, *Quantum Field Theory for the Gifted Amateur*, (Oxford University Press, 2014), Problem 12.7. Post date: 7 May 2019.

The complex scalar field considered by L&B in their Chapter 12 is invariant under the internal transformation

$$\psi \to \psi' = e^{i\alpha}\psi \tag{1}$$

They point out that this transformation can also be written using the unitary operator

$$U(\alpha) = e^{iQ_{Nc}\alpha} \tag{2}$$

where Q_{Nc} is the conserved number charge operator, which is defined in terms of the Noether charge Q_N by normal ordering:

$$Q_{Nc} = -:Q_N: \tag{3}$$

To verify that this unitary operator works correctly, we consider $U^{\dagger}\psi U$ for an infinitesimal α . We have, keeping only terms up to first order in α

$$U^{\dagger}(\alpha)\psi U(\alpha) = (1 - iQ_{Nc}\alpha)\psi(1 + iQ_{Nc}\alpha)$$
(4)

$$=\psi - i\alpha \left[Q_{Nc},\psi\right] \tag{5}$$

We can now use the formula

$$[Q_N,\psi] = -iD\psi \tag{6}$$

where

$$D\psi = \frac{\partial\psi'}{\partial\alpha}\bigg|_{\alpha=0} = i\psi \tag{7}$$

From 3 we have

$$[Q_{Nc},\psi] = -[Q_N,\psi] = -\psi \tag{8}$$

so we have

$$U^{\dagger}(\alpha)\psi U(\alpha) = \psi(1+i\alpha)$$
(9)

Taking α to be finite, we can apply this formula n times and let $n \to \infty$ to get

$$U^{\dagger}(\alpha)\psi U(\alpha) = \lim_{n \to \infty} \psi \left(1 + i\frac{\alpha}{n}\right)^n \tag{10}$$

$$=e^{i\alpha}\psi\tag{11}$$

Thus the unitary transformation is equivalent to the transformation 1.

In case you're worried that the normal ordering done to get from Q_N to Q_{Nc} messes up the relation 8, we can see that this isn't a problem by considering Q_{Nc} and ψ in their mode expansions. From L&B's equations 12.5 and 12.15, we have

$$\psi\left(x\right) = \int \frac{d^3p}{\left(2\pi\right)^{3/2} \sqrt{2E_{\mathbf{p}}}} \left(a_{\mathbf{p}}e^{-ip\cdot x} + b_{\mathbf{p}}^{\dagger}e^{ip\cdot x}\right)$$
(12)

$$Q_{Nc} = \int d^3q \left(n_{\mathbf{q}}^{(a)} - n_{\mathbf{q}}^{(b)} \right) \tag{13}$$

$$= \int d^3q \left(a_{\mathbf{q}}^{\dagger} a_{\mathbf{q}} - b_{\mathbf{q}}^{\dagger} b_{\mathbf{q}} \right) \tag{14}$$

Taking the commutator explicitly and using

$$\begin{bmatrix} a_{\mathbf{p}}, a_{\mathbf{q}}^{\dagger} \end{bmatrix} = \begin{bmatrix} b_{\mathbf{p}}, b_{\mathbf{q}}^{\dagger} \end{bmatrix} = \delta\left(\mathbf{p} - \mathbf{q}\right)$$
(15)

with all other commutators being zero, we have

$$[Q_{Nc},\psi] = \int \frac{d^3p \ d^3q}{(2\pi)^{3/2} \sqrt{2E_{\mathbf{p}}}} \left\{ \left[a_{\mathbf{q}}^{\dagger}, a_{\mathbf{p}} \right] a_{\mathbf{q}} e^{-ip \cdot x} - b_{\mathbf{q}}^{\dagger} \left[b_{\mathbf{q}}, b_{\mathbf{p}}^{\dagger} \right] e^{ip \cdot x} \right\}$$
(16)
$$= \int \frac{d^3p \ d^3q}{(2\pi)^{3/2} \sqrt{2E_{\mathbf{p}}}} \left\{ - \left[a_{\mathbf{p}}, a_{\mathbf{q}}^{\dagger} \right] a_{\mathbf{q}} e^{-ip \cdot x} - b_{\mathbf{q}}^{\dagger} \left[b_{\mathbf{q}}, b_{\mathbf{p}}^{\dagger} \right] e^{ip \cdot x} \right\}$$
(17)

$$= \int \frac{d^3 p \, d^3 q}{\left(2\pi\right)^{3/2} \sqrt{2E_{\mathbf{p}}}} \left\{ -\delta\left(\mathbf{p}-\mathbf{q}\right) a_{\mathbf{q}} e^{-ip\cdot x} - b_{\mathbf{q}}^{\dagger} \delta\left(\mathbf{p}-\mathbf{q}\right) e^{ip\cdot x} \right\}$$
(18)

$$= -\int \frac{d^3p}{\left(2\pi\right)^{3/2}\sqrt{2E_{\mathbf{p}}}} \left(a_{\mathbf{p}}e^{-ip\cdot x} + b_{\mathbf{p}}^{\dagger}e^{ip\cdot x}\right)$$
(19)

$$= -\psi \tag{20}$$