

## EVENT HORIZON: TIME AND SPACE SWAP ROUND

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Reference: Moore, Thomas A., *A General Relativity Workbook*, University Science Books (2013) - Chapter 14; Box 14.3.

When we first looked at the  $t$  component of the Schwarzschild metric we noted that for an object at rest, the  $t$  component is related to the object's proper time by

$$(1) \quad \Delta\tau = \sqrt{1 - \frac{2GM}{r}} \Delta t$$

If the object is at the event horizon, that is,  $r = 2GM$ , then  $\Delta\tau = 0$  and since  $ds^2 = -d\tau^2$ , this means that  $ds^2 = 0$  for the object, no matter what  $\Delta t$  is. A zero space-time interval can exist only for photons (or other massless particles), so what happens if a massive particle approaches  $r = 2GM$ ? The only way we can reconcile the Schwarzschild metric with an object at the event horizon is if the object is not at rest (remember the assumption that the object *was* at rest led to the zero space-time interval, so that assumption must be wrong). In other words, as an object approaches the event horizon, it is compelled to keep moving; there is no way it can stop itself from continuing past the event horizon.

The explanation of this phenomenon goes like this. In any metric, a time-like interval is always represented by a value of  $ds^2 < 0$ . In the Schwarzschild metric:

$$(2) \quad ds^2 = - \left(1 - \frac{2GM}{r}\right) dt^2 + \left(1 - \frac{2GM}{r}\right)^{-1} dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2$$

this means that the two events that define the interval must always be separated by a non-zero interval in the coordinate  $t$ . It is possible for all the other intervals (the space intervals  $dr$ ,  $d\theta$  and  $d\phi$ , for  $r > 2GM$ ) to be zero, that is, it's possible for the two events to occur at the same place, but they must always be separated by a non-zero time interval.

If we carry the Schwarzschild metric through the event horizon so that  $r < 2GM$ , then the signs of the  $dt^2$  and  $dr^2$  components flip, so that  $-\left(1 - \frac{2GM}{r}\right) dt^2 > 0$  and  $\left(1 - \frac{2GM}{r}\right)^{-1} dr^2 < 0$ . Thus a timelike interval now means that  $dr$

must be non-zero, rather than  $dt$ . In other words, the physical meanings of  $r$  and  $t$  have swapped;  $r$  now plays the role of a time component and  $t$  of a space component.

Inside the event horizon, we can write the metric as

$$(3) \quad ds^2 = - \left( \frac{2GM}{r} - 1 \right)^{-1} dr^2 + \left( \frac{2GM}{r} - 1 \right) dt^2 + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2$$

The proper time for an object at rest is now

$$(4) \quad d\tau^2 = -ds^2 = \left( \frac{2GM}{r} - 1 \right)^{-1} dr^2 - \left( \frac{2GM}{r} - 1 \right) dt^2 - r^2 d\theta^2 - r^2 \sin^2 \theta d\phi^2$$

Since  $\frac{2GM}{r} - 1 > 0$  inside the event horizon, the maximum proper time interval occurs when  $dt = d\theta = d\phi = 0$ , in other words, for a purely radial path. Since a geodesic is the path of longest proper time, the geodesic is a purely radial path through space-time. For a path from  $r = 0$  to  $r = 2GM$ , this time interval is given by the same formula we evaluated in the last post, with  $R = 2GM$

$$(5) \quad \Delta\tau = \frac{\pi R^{3/2}}{\sqrt{8GM}} = \pi GM$$

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