# GRAVITOELECTRIC AND GRAVITOMAGNETIC DENSITIES

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Reference: Moore, Thomas A., *A General Relativity Workbook*, University Science Books (2013) - Chapter 22, Box 22.6.

We can now have a look at the equation of motion for a particle moving in the weak field, slow velocity limit. We begin with the geodesic equation

$$\ddot{x}^m + \Gamma^m_{ij}\dot{x}^j\dot{x}^i = 0 \tag{0.1}$$

where a dot indicates a derivative with respect to proper time $\tau$. Using the chain rule we can write the first term as

$$\begin{aligned}
\frac{d^2x^m}{d\tau^2} &= \frac{d}{d\tau}\left(\frac{dx^m}{d\tau}\right) &(0.2)\\
&= \frac{dt}{d\tau}\frac{d}{dt}\left(\frac{dx^m}{d\tau}\right) &(0.3)\\
&= \frac{dt}{d\tau}\frac{d}{dt}\left(\frac{dt}{d\tau}\frac{dx^m}{dt}\right) &(0.4)\\
&= u^t\frac{dt}{d\tau}\frac{d^2x^m}{dt^2} + v^m\frac{dt}{d\tau}\frac{d}{dt}\left(\frac{dt}{d\tau}\right) &(0.5)\\
&= u^t u^t\frac{d^2x^m}{dt^2} + v^m\frac{d}{d\tau}\left(\frac{dt}{d\tau}\right) &(0.6)\\
&= u^t u^t\frac{d^2x^m}{dt^2} + v^m\frac{du^t}{d\tau} &(0.7)
\end{aligned}$$

where

$$v^m \equiv \frac{dx^m}{dt} \tag{0.8}$$

Now suppose the particle is moving slowly so that both $u^i$ (where $i$ is a spatial index) and $v^i$ are small, so we can ignore second order terms. In this approximation, $u^t \approx 1$ so we have from 0.1 with $m = t$

$$\frac{du^t}{d\tau} = -\Gamma^t{}_{ij}\dot{x}^j\dot{x}^i \tag{0.9}$$
$$= -\Gamma^t{}_{ij}u^j u^i \tag{0.10}$$
$$\approx -\Gamma^t{}_{tt}u^t u^t \tag{0.11}$$
$$= -\Gamma^t{}_{tt} \tag{0.12}$$

Therefore

$$v^m \frac{du^t}{d\tau} \approx -v^m \Gamma^t{}_{tt} \tag{0.13}$$

The Christoffel symbols in terms of the metric are

$$\Gamma^m_{ij} = \frac{1}{2} g^{ml} \left(\partial_j g_{il} + \partial_i g_{lj} - \partial_l g_{ji}\right) \tag{0.14}$$

For a stationary field, all time derivatives are zero, so

$$\Gamma^t{}_{tt} = -\frac{1}{2} g^{tl} \partial_l g_{tt} \tag{0.15}$$

In the weak field limit, we are using a metric that is perturbed from the flat space metric:

$$g_{ij} = \eta_{ij} + h_{ij} \tag{0.16}$$
$$g^{ij} = \eta^{ij} - h^{ij} \tag{0.17}$$

So to first order in $h_{ij}$:

$$\Gamma^t{}_{tt} = -\frac{1}{2}\left(\eta^{tl} - h^{tl}\right) \partial_l \left(\eta_{tt} + h_{tt}\right) \tag{0.18}$$
$$= -\frac{1}{2}\eta^{tt}\partial_t h_{tt} + \frac{1}{2} h^{tl} \partial_l \eta_{tt} \tag{0.19}$$
$$= 0 \tag{0.20}$$

since all time derivatives are zero in the stationary field, and all derivatives of the flat space metric are zero as well. Therefore in this approximation 0.7 becomes

$$\frac{d^2 x^m}{d\tau^2} = u^t u^t \frac{d^2 x^m}{dt^2} = \frac{d^2 x^m}{dt^2} \tag{0.21}$$

and the geodesic equation 0.1 becomes

$$\frac{d^2x^m}{dt^2} = -\Gamma^m_{ij}u^j u^i \tag{0.22}$$
$$\approx -\Gamma^m_{tt}u^t u^t - \Gamma^m_{it}u^t u^i - \Gamma^m_{tj}u^t u^j \tag{0.23}$$
$$= -\Gamma^m_{tt} - 2\Gamma^m_{it}u^i \tag{0.24}$$

where to get the last line we used the symmetry of the Christoffel symbols: $\Gamma^m_{it} = \Gamma^m_{ti}$.

In the weak field limit, the Christoffel symbols become

$$\Gamma^m_{ij} = \frac{1}{2}\eta^{ml}\left(\partial_j h_{il} + \partial_i h_{lj} - \partial_l h_{ji}\right) \tag{0.25}$$

In this approximation

$$\Gamma^m_{tt} = \frac{1}{2}\eta^{ml}\left(\partial_t h_{tl} + \partial_t h_{lt} - \partial_l h_{tt}\right) \tag{0.26}$$
$$= -\frac{1}{2}\eta^{ml}\partial_l h_{tt} \tag{0.27}$$

since time derivatives are zero for a stationary field.

Also

$$2\Gamma^m_{it}u^i = u^i\eta^{ml}\left(\partial_t h_{il} + \partial_i h_{lt} - \partial_l h_{it}\right) \tag{0.28}$$
$$= u^i\eta^{ml}\left(\partial_i h_{lt} - \partial_l h_{it}\right) \tag{0.29}$$

Putting it all together, the geodesic equation 0.24 becomes

$$\frac{d^2x^m}{dt^2} = \eta^{ml}\left[\frac{1}{2}\partial_l h_{tt} + u^i\left(\partial_l h_{it} - \partial_i h_{lt}\right)\right] \tag{0.30}$$

Remember that all indices (except $t$) range over only spatial components.

In the weak field limit, the perturbations to the metric are given by

$$h_{jm} = 2G\int \frac{2T_{jm} - \eta_{jm}T}{|\mathbf{r}-\mathbf{r}'|}d^3\mathbf{r}' \tag{0.31}$$

The numerator in the integrand is analogous to charge density in electromagnetism. The first term in 0.30 results from the density

$$\rho_g \equiv 2T_{tt} - \eta_{tt}T \tag{0.32}$$

which is called the *gravitoelectric energy density*, since it gives rise to an acceleration of the particle that does not depend on the particle's velocity, in a similar way to an electric field accelerating a charge independently of the charge's velocity.

The second term in 0.30 results from the density

$$\Pi_j \equiv T_{tj} - \frac{1}{2}\eta_{tj}T \tag{0.33}$$

which is called the *gravitomagnetic current density* since it contributes to an acceleration that is proportional to the particle's velocity, in a similar way to the Lorentz force arising from magnetism.

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