DIFFERENTIAL OPERATOR - EIGENVALUES AND EIGENSTATES

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Continuing with our study of differential operators, we'll look now at their eigenvalues and eigenstates. The operator we're studying is

$$K = -i\frac{d}{dx} \tag{1}$$

The eigenvalue equation is as usual:

$$K|k\rangle = k|k\rangle \tag{2}$$

where $|k\rangle$ is an eigenstate and k (outside the ket) is a (possibly complex) scalar. To find $|k\rangle$, we form the matrix element with $\langle x|$ and insert the unit operator:

$$\langle x | K | k \rangle = k \langle x | k \rangle \tag{3}$$

$$\langle x | K | k \rangle = \int \langle x | K | x' \rangle \langle x' | k \rangle dx'$$
(4)

$$= -i \int \delta' \left(x - x' \right) \psi_k \left(x' \right) dx' \tag{5}$$

$$= -i\frac{d}{dx}\psi_k(x) \tag{6}$$

In the third line we used the matrix element

$$\left\langle x \left| K \right| x' \right\rangle = -i\delta' \left(x - x' \right) \tag{7}$$

Equating the RHS on the first and last lines gives the differential equation

$$-i\frac{d}{dx}\psi_k(x) = k\psi_k(x) \tag{8}$$

which has the solution

$$\psi_k\left(x\right) = Ae^{ikx} \tag{9}$$

where A is a constant of integration. In order for $\psi_k(x)$ to be bounded as $x \to \pm \infty$, k must be real, so we'll restrict our attention to that case. The usual choice for A is $1/\sqrt{2\pi}$ so that

$$\psi_k(x) = \frac{e^{ikx}}{\sqrt{2\pi}} \tag{10}$$

This leads to the normalization condition

$$\left\langle k\left|k'\right\rangle = \int_{-\infty}^{\infty} \left\langle k\left|x\right\rangle \left\langle x\left|k'\right\rangle dx\right. \right.$$

$$(11)$$

$$=\frac{1}{2\pi}\int_{-\infty}^{\infty}e^{-i(k-k')x}dx$$
(12)

$$=\delta\left(k-k'\right) \tag{13}$$

where in the last line we used the traditional formula for the delta function. Thus the $|k\rangle$ basis is orthogonal, and normalized the same way as the $|x\rangle$ basis.

To convert between the $|k\rangle$ and $|x\rangle$ bases, we can use the unit operator in the two bases. Thus for some vector (function) $|f\rangle$ we have

$$f(k) = \langle k | f \rangle = \int \langle k | x \rangle \langle x | f \rangle dx = \int \psi_k^*(x) f(x) dx = \frac{1}{\sqrt{2\pi}} \int e^{-ikx} f(x) dx$$
(14)

Thus f(k) is the Fourier transform of f(x). We can use the same procedure to go in the reverse direction:

$$f(x) = \langle x | f \rangle = \int \langle x | k \rangle \langle k | f \rangle dk = \int \psi_k(x) f(k) dk = \frac{1}{\sqrt{2\pi}} \int e^{ikx} f(k)$$
(15)

The effect of the position operator X on a vector $|f(x)\rangle$ can be found by inserting the unit operator:

$$\langle x | X | f \rangle = \int \langle x | X | x' \rangle \langle x' | f \rangle dx'$$
(16)

$$= \int x' \langle x | x' \rangle \langle x' | f \rangle dx'$$
(17)

$$= \int x' \delta(x - x') \langle x' | f \rangle dx'$$
(18)

$$= x \left\langle x \left| f \right\rangle \right. \tag{19}$$

Thus X just multiplies any function of x by x itself. A similar argument in the $|k\rangle$ basis shows that

$$\langle k | K | f(k) \rangle = k \langle k | f(k) \rangle \tag{20}$$

We can use similar calculations to find the matrix elements of K in the $|x\rangle$ basis and of X (the position operator) in the $|k\rangle$ basis. We get

$$\left\langle k \left| X \right| k' \right\rangle = \int \int \left\langle k \left| x \right\rangle \left\langle x \left| X \right| x' \right\rangle \left\langle x' \left| k' \right\rangle dx \, dx' \right\rangle \right\rangle$$
(21)

$$= \frac{1}{2\pi} \int \int e^{-ikx} x' \langle x | x' \rangle e^{ik'x'} dx dx'$$
(22)

$$=\frac{1}{2\pi}\int\int e^{-ikx}x'\delta\left(x-x'\right)e^{ik'x'}dx\,dx'$$
(23)

$$=\frac{1}{2\pi}\int xe^{i(k'-k)x}dx$$
(24)

$$=i\frac{d}{dk}\left[\frac{1}{2\pi}\int e^{i(k'-k)x}dx\right]$$
(25)

$$=i\delta'\left(k-k'\right) \tag{26}$$

The action of X on an arbitrary vector $|g\rangle$ in the k basis can be found from this:

$$\langle k | X | g(k) \rangle = \int \langle k | X | k' \rangle \langle k' | g \rangle dk'$$
(27)

$$= i \int \delta' \left(k - k' \right) g \left(k' \right) dk' \tag{28}$$

$$=i\frac{dg\left(k\right)}{dk}\tag{29}$$

$$=i\left\langle k\left|\frac{dg\left(k\right)}{dk}\right\rangle \right. \tag{30}$$

where in the third line we've used the property of $\delta'(k-k')$ mentioned here.

By a similar calculation, we can find the matrix elements of K in the $|x\rangle$ basis:

$$\left\langle x \left| K \right| x' \right\rangle = \int \int \left\langle x \left| k \right\rangle \left\langle k \left| K \right| k' \right\rangle \left\langle k' \left| x' \right\rangle dk \, dk' \right\rangle$$
(31)

$$= \frac{1}{2\pi} \int \int e^{ikx} k' \langle k | k' \rangle e^{-ik'x'} dk dk'$$
(32)

$$=\frac{1}{2\pi}\int\int e^{ikx}k'\delta\left(k-k'\right)e^{-ik'x'}dk\,dk'$$
(33)

$$=\frac{1}{2\pi}\int xe^{i(x-x')k}dk$$
(34)

$$= -i\frac{d}{dx}\left[\frac{1}{2\pi}\int e^{i(x-x')k}dk\right]$$
(35)

$$= -i\delta'\left(x - x'\right) \tag{36}$$

Similarly, we have

$$\langle x | K | g(x) \rangle = \int \langle x | K | x' \rangle \langle x' | g \rangle dx'$$
(37)

$$= -i \int \delta' \left(x - x' \right) g\left(x' \right) dx' \tag{38}$$

$$=-i\frac{dg\left(x\right)}{dx}\tag{39}$$

$$= -i\left\langle x \left| \frac{dg\left(x\right)}{dx} \right\rangle \right. \tag{40}$$

From 30 and 40 we can work out the familiar commutator. Just for variety, we'll do this in the $|k\rangle$ basis:

$$XK|f(k)\rangle = X[k|f(k)\rangle]$$
(41)

$$=i\frac{d}{dk}[k|f(k)\rangle]$$
(42)

$$= i \left[\left| f\left(k\right) \right\rangle + k \left| \frac{df}{dk} \right\rangle \right] \tag{43}$$

$$KX|f(k)\rangle = iK\left|\frac{df}{dk}\right\rangle \tag{44}$$

$$=ik\left|\frac{df}{dk}\right\rangle \tag{45}$$

Therefore

$$[X,K]|f(k)\rangle = i|f(k)\rangle \tag{46}$$

or, looking just at the operators

$$[X,K] = iI \tag{47}$$