

VECTOR OPERATORS - TRANSFORMATION UNDER ROTATION

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A vector operator \mathbf{V} is defined as an operator whose components transform under rotation according to

$$U^\dagger [R] V_i U [R] = \sum_j R_{ij} V_j \quad (1)$$

where R is the rotation matrix in either 2 or 3 dimensions. We've seen that, for an infinitesimal rotation about an arbitrary axis $\delta\boldsymbol{\theta}$, a vector transforms like

$$\mathbf{V} \rightarrow \mathbf{V} + \delta\boldsymbol{\theta} \times \mathbf{V} \quad (2)$$

This can be written more compactly using the totally antisymmetric Levi-Civita tensor ε_{ijk} , where $\varepsilon_{123} = +1$ and the tensor is antisymmetric when any two indices are swapped. The component i of a cross product is

$$(\delta\boldsymbol{\theta} \times \mathbf{V})_i = \sum_{j,k} \varepsilon_{ijk} (\delta\theta)_j V_k \quad (3)$$

We get

$$\sum_j R_{ij} V_j = V_i + \sum_{j,k} \varepsilon_{ijk} (\delta\theta)_j V_k \quad (4)$$

The operator $U [R]$ is given by

$$U [R(\delta\boldsymbol{\theta})] = I - \frac{i}{\hbar} \delta\boldsymbol{\theta} \cdot \mathbf{L} \quad (5)$$

where \mathbf{L} is the angular momentum. Plugging this into 1, we have, to first order in $\delta\boldsymbol{\theta}$ (remembering that the components of \mathbf{L} do not commute with each other and, in general also do not commute with the components of \mathbf{V}):

$$\left(I + \frac{i}{\hbar} \delta \boldsymbol{\theta} \cdot \mathbf{L}\right) V_i \left(I - \frac{i}{\hbar} \delta \boldsymbol{\theta} \cdot \mathbf{L}\right) = V_i + \frac{i}{\hbar} \sum_j (\delta \theta_j L_j) V_i - \frac{i}{\hbar} V_i \sum_j (\delta \theta_j L_j) \quad (6)$$

$$= V_i + \frac{i}{\hbar} \sum_j \delta \theta_j [L_j, V_i] \quad (7)$$

Setting this equal to the RHS of 4 we have, equating coefficients of $\delta \theta_j$:

$$\frac{i}{\hbar} [L_j, V_i] = \sum_k \varepsilon_{ijk} V_k \quad (8)$$

$$[V_i, L_j] = i\hbar \sum_k \varepsilon_{ijk} V_k \quad (9)$$

With $\mathbf{V} = \mathbf{L}$, we regain the commutation relations for the components of angular momentum

$$[L_x, L_y] = i\hbar L_z \quad (10)$$

$$[L_y, L_z] = i\hbar L_x \quad (11)$$

$$[L_z, L_x] = i\hbar L_y \quad (12)$$

By the way, it is possible to write these commutation relations in the compact form

$$\mathbf{L} \times \mathbf{L} = i\hbar \mathbf{L} \quad (13)$$

This looks wrong if you're used to the standard definition of the cross product for vectors whose components are ordinary numbers, since for such a vector \mathbf{a} , we always have $\mathbf{a} \times \mathbf{a} = 0$. However, if the components of the vector are *operators* that don't commute, then the result is not zero, as we can see:

$$(\mathbf{L} \times \mathbf{L})_i = \sum_{j,k} \varepsilon_{ijk} L_j L_k \quad (14)$$

If $i = x$, for example, then the sum on the RHS gives

$$(\mathbf{L} \times \mathbf{L})_x = \sum_{j,k} \varepsilon_{xjk} L_j L_k \quad (15)$$

$$= L_y L_z - L_z L_y \quad (16)$$

$$= [L_y, L_z] \quad (17)$$

From 13, this gives

$$[L_y, L_z] = i\hbar L_x \quad (18)$$