

HARMONIC OSCILLATOR: MOMENTUM SPACE FUNCTIONS AND HERMITE POLYNOMIAL RECURSION RELATIONS FROM RAISING AND LOWERING OPERATORS

Link to: physicspages home page.

To leave a comment or report an error, please use the auxiliary blog.

Shankar, R. (1994), *Principles of Quantum Mechanics*, Plenum Press.
Section 7.5, Exercises 7.5.1 - 7.5.3.

Earlier, we found the position space energy eigenfunctions of the harmonic oscillator to be

$$(0.1) \quad \psi_n(y) = \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} \frac{1}{\sqrt{2^n n!}} H_n(y) e^{-y^2/2}$$

$$(0.2) \quad \psi_n(x) = \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} \frac{1}{\sqrt{2^n n!}} H_n\left(\sqrt{\frac{m\omega}{\hbar}}x\right) e^{-m\omega x^2/2\hbar}$$

where y in the first equation is shorthand for

$$(0.3) \quad y = \sqrt{\frac{m\omega}{\hbar}}x$$

It turns out that an alternative method for deriving these functions uses the lowering operator a . Shankar gives the derivation of $\psi_n(x)$ in his section 7.5, but we can use the same technique to derive the momentum space functions. We start with the ground state and use

$$(0.4) \quad a|0\rangle = 0$$

In terms of X and P , we have

$$(0.5) \quad a = \sqrt{\frac{m\omega}{2\hbar}}X + i\frac{1}{\sqrt{2m\omega\hbar}}P$$

To find the momentum space functions, we need to express X and P in terms of p :

$$(0.6) \quad X = i\hbar \frac{d}{dp}$$

$$(0.7) \quad P = p$$

We thus have

$$(0.8) \quad \left[i\hbar\sqrt{\frac{m\omega}{2\hbar}} \frac{d}{dp} + i\frac{1}{\sqrt{2m\omega\hbar}} p \right] \psi_0(p) = 0$$

If we define the auxiliary variable

$$(0.9) \quad z \equiv \frac{p}{\sqrt{\hbar m\omega}}$$

we get

$$(0.10) \quad \left(\frac{d}{dz} + z \right) \psi_0(z) = 0$$

This has the solution

$$(0.11) \quad \psi_0(z) = Ae^{-z^2/2}$$

for some normalization constant A . Thus in terms of p we have

$$(0.12) \quad \psi_0(p) = Ae^{-p^2/2\hbar m\omega}$$

Normalizing in the usual way, making use of the Gaussian integral, we have

$$(0.13) \quad \int_{-\infty}^{\infty} \psi_0^2(p) dp = A^2 \int_{-\infty}^{\infty} e^{-p^2/\hbar m\omega} dp = 1$$

$$(0.14) \quad A = \frac{1}{(\pi\hbar m\omega)^{1/4}}$$

This agrees with the earlier result which was obtained by solving a second-order differential equation.

We can also use a and a^\dagger to verify a couple of recursion relations for Hermite polynomials. Reverting back to position space we have

$$(0.15) \quad X = x$$

$$(0.16) \quad P = -i\hbar \frac{d}{dx}$$

so 0.5 becomes

$$(0.17) \quad a = \sqrt{\frac{m\omega}{2\hbar}} x + \frac{\hbar}{\sqrt{2m\omega\hbar}} \frac{d}{dx}$$

Also from 0.5 we have, since X and P are both hermitian operators

$$(0.18) \quad a^\dagger = \sqrt{\frac{m\omega}{2\hbar}}X - i\frac{1}{\sqrt{2m\omega\hbar}}P$$

$$(0.19) \quad = \sqrt{\frac{m\omega}{2\hbar}}x - \frac{\hbar}{\sqrt{2m\omega\hbar}}\frac{d}{dx}$$

Defining

$$(0.20) \quad y \equiv \sqrt{\frac{m\omega}{\hbar}}x$$

we have

$$(0.21) \quad a = \frac{1}{\sqrt{2}}\left(y + \frac{d}{dy}\right)$$

$$(0.22) \quad a^\dagger = \frac{1}{\sqrt{2}}\left(y - \frac{d}{dy}\right)$$

We also recall the normalization conditions on the raising and lowering operators:

$$(0.23) \quad a|n\rangle = \sqrt{n}|n-1\rangle$$

$$(0.24) \quad a^\dagger|n\rangle = \sqrt{n+1}|n+1\rangle$$

Applying 0.23 to 0.1 we have, after cancelling common factors from each side:

$$(0.25) \quad \frac{1}{\sqrt{2}}\frac{1}{\sqrt{2^n n!}}\left(y + \frac{d}{dy}\right)\left[H_n(y)e^{-y^2/2}\right] = \frac{\sqrt{n}}{\sqrt{2^{n-1}(n-1)!}}H_{n-1}(y)e^{-y^2/2}$$

$$(0.26) \quad \frac{1}{2\sqrt{n}}\frac{1}{\sqrt{2^{n-1}(n-1)!}}e^{-y^2/2}\left[yH_n(y) - yH_n(y) + \frac{dH_n}{dy}\right] = \frac{\sqrt{n}}{\sqrt{2^{n-1}(n-1)!}}H_{n-1}(y)e^{-y^2/2}$$

$$(0.27) \quad yH_n(y) - yH_n(y) + \frac{dH_n}{dy} = 2nH_{n-1}(y)$$

$$(0.28) \quad H'_n(y) = 2nH_{n-1}(y)$$

Another recursion relation for Hermite polynomials can be found as follows. We start with 0.22 to get

$$(0.29) \quad a + a^\dagger = \sqrt{2}y$$

We now apply 0.23 and 0.24 to 0.1. We can cancel common factors, including $e^{-y^2/2}$, from both sides to get

$$(0.30)$$

$$(a + a^\dagger) \psi_n = \sqrt{2}y \psi_n$$

$$(0.31)$$

$$\frac{\sqrt{2}y}{\sqrt{2^n n!}} H_n(y) = \frac{\sqrt{n}}{\sqrt{2^{n-1} (n-1)!}} H_{n-1}(y) + \frac{\sqrt{n+1}}{\sqrt{2^{n+1} (n+1)!}} H_{n+1}(y)$$

$$(0.32)$$

$$\frac{y}{\sqrt{2^{n-1} n (n-1)!}} H_n(y) = \frac{\sqrt{n}}{\sqrt{2^{n-1} (n-1)!}} H_{n-1}(y) + \frac{1}{2\sqrt{2^{n-1} n (n-1)!}} H_{n+1}(y)$$

$$(0.33) \quad yH_n(y) = nH_{n-1}(y) + \frac{1}{2}H_{n+1}(y)$$

$$(0.34) \quad H_{n+1}(y) = 2yH_n(y) - 2nH_{n-1}(y)$$