

## TRANSLATION OPERATOR FROM PASSIVE TRANSFORMATIONS

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Shankar, R. (1994), *Principles of Quantum Mechanics*, Plenum Press. Chapter 11.

We've seen that the translation operator  $T(\epsilon)$  in quantum mechanics can be derived by considering the translation to be an active transformation, that is, a transformation where the state vectors, rather than the operators, get transformed according to

$$(0.1) \quad T(\epsilon) |\psi\rangle = |\psi_\epsilon\rangle$$

Using this approach, we found that

$$(0.2) \quad T(\epsilon) = I - \frac{i\epsilon}{\hbar} P$$

so that the momentum  $P$  is the generator of the transformation.

We can also derive  $T$  using a passive transformation, where the state vectors remain the same but the operators are transformed according to

$$(0.3) \quad T^\dagger(\epsilon) X T(\epsilon) = X + \epsilon I$$

$$(0.4) \quad T^\dagger(\epsilon) P T(\epsilon) = P$$

This is equivalent to an active transformation since

$$(0.5) \quad \langle \psi | T^\dagger(\epsilon) X T(\epsilon) | \psi \rangle = \langle T(\epsilon) \psi | X | T(\epsilon) \psi \rangle$$

$$(0.6) \quad = \langle \psi_\epsilon | X | \psi_\epsilon \rangle$$

$$(0.7) \quad = x + \epsilon$$

As before we start by taking

$$(0.8) \quad T(\epsilon) = I - \frac{i\epsilon}{\hbar} G$$

where  $G$  is some Hermitian operator, so that  $G^\dagger = G$ . Plugging this into 0.3 we get, keeping only terms up to order  $\epsilon$ :

$$(0.9) \quad T^\dagger(\varepsilon)XT(\varepsilon) = \left(I + \frac{i\varepsilon}{\hbar}G\right)X\left(I - \frac{i\varepsilon}{\hbar}G\right)$$

$$(0.10) \quad = X + \frac{i\varepsilon}{\hbar}I(GX - XG)$$

$$(0.11) \quad = X - \frac{i\varepsilon}{\hbar}[X, G]$$

$$(0.12) \quad = X + \varepsilon I$$

Therefore

$$(0.13) \quad -\frac{i\varepsilon}{\hbar}[X, G] = \varepsilon I$$

$$(0.14) \quad [X, G] = i\hbar I$$

Since  $[X, P] = i\hbar$  we see that

$$(0.15) \quad G = P + f(X)$$

The extra  $f(X)$  is there because any function of  $X$  alone commutes with  $X$ , so

$$(0.16) \quad [X, G] = [X, P] + [X, f(X)] = i\hbar I + 0$$

We can eliminate  $f(X)$  by considering 0.4.

$$(0.17) \quad T^\dagger(\varepsilon)PT(\varepsilon) = \left(I + \frac{i\varepsilon}{\hbar}G\right)P\left(I - \frac{i\varepsilon}{\hbar}G\right)$$

$$(0.18) \quad = P + \frac{i\varepsilon}{\hbar}I(GP - PG)$$

$$(0.19) \quad = P - \frac{i\varepsilon}{\hbar}[P, G]$$

$$(0.20) \quad = P$$

Thus we must have  $[P, G] = 0$ , which means that  $G$  must be a function of  $P$  alone. This means that the most general form for  $f(X)$  is  $f(X) = \text{constant}$ , but there's nothing to be gained by adding some non-zero constant to  $G$ , so we can take  $f(X) = 0$ . Thus we end up with the same form 0.2 that we got from the active transformation.

Translational invariance is the condition that the Hamiltonian is unaltered by a translation. In the passive representation this is stated by the condition

$$(0.21) \quad T^\dagger(\varepsilon)HT(\varepsilon) = H$$

Since translation is unitary, we can apply a theorem that is valid for any operator  $\Omega$  which can be expanded in powers of  $X$  and  $P$ . For any unitary operator  $U$ , we have

$$(0.22) \quad U^\dagger\Omega(X,P)U = \Omega(U^\dagger XU, U^\dagger PU)$$

This follows because for a unitary operator  $U^\dagger U = UU^\dagger = I$  so we can insert the product  $UU^\dagger$  anywhere we like. In particular, we can insert it between each pair of factors in every term of the power series expansion of  $\Omega$ , for example

$$(0.23) \quad U^\dagger X^2 P^2 U = U^\dagger X X P P U$$

$$(0.24) \quad = U^\dagger X U U^\dagger X U U^\dagger P U U^\dagger P U$$

$$(0.25) \quad = (U^\dagger X U)^2 (U^\dagger P U)^2$$

For 0.21 this means that

$$(0.26) \quad T^\dagger(\varepsilon)H(X,P)T(\varepsilon) = H(X + \varepsilon I, P) = H(X,P)$$

As before, this leads to the condition

$$(0.27) \quad [P, H] = 0$$

which means that  $P$  is conserved, according to Ehrenfest's theorem.

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