

## SPHERICAL BESSEL FUNCTIONS: BEHAVIOUR FOR SMALL ARGUMENTS

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Shankar, R. (1994), *Principles of Quantum Mechanics*, Plenum Press. Chapter 12, Exercise 12.6.7.

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The general solution for a free particle in spherical coordinates involves the radial function, which turns out to be

$$(1) \quad \frac{R_{l+1}}{\rho^{l+1}} = \left( -\frac{1}{\rho} \frac{d}{d\rho} \right)^{l+1} \frac{R_0}{\rho^0}$$

where  $l$  is the total angular momentum quantum number and

$$(2) \quad k^2 \equiv \frac{2\mu E}{\hbar^2}$$

$$(3) \quad \rho \equiv kr$$

We can rewrite this as

$$(4) \quad R_l = (-\rho)^l \left( \frac{1}{\rho} \frac{d}{d\rho} \right)^l R_0$$

We saw earlier that the solutions for  $l = 0$  are, with  $U_l = \rho R_l$

$$(5) \quad U_0^A(\rho) = \sin \rho$$

$$(6) \quad U_0^B(\rho) = -\cos \rho$$

Thus the two solutions for  $l = 0$  are

$$(7) \quad R_0^A = \frac{\sin \rho}{\rho}$$

$$(8) \quad R_0^B = -\frac{\cos \rho}{\rho}$$

From these starting points, we can generate all the solutions for higher values of  $l$  using 4. These functions are

$$(9) \quad j_l(\rho) = (-\rho)^l \left( \frac{1}{\rho} \frac{d}{d\rho} \right)^l \frac{\sin \rho}{\rho}$$

$$(10) \quad n_l(\rho) = -(-\rho)^l \left( \frac{1}{\rho} \frac{d}{d\rho} \right)^l \frac{\cos \rho}{\rho}$$

and are known as spherical Bessel functions  $j_l$  and spherical Neumann functions  $n_l$ .

The asymptotic behaviour is given by

$$(11) \quad j_l \xrightarrow{\rho \rightarrow \infty} \frac{1}{\rho} \sin \left( \rho - \frac{l\pi}{2} \right)$$

$$(12) \quad n_l \xrightarrow{\rho \rightarrow \infty} -\frac{1}{\rho} \cos \left( \rho - \frac{l\pi}{2} \right)$$

For  $\rho \rightarrow 0$ , we have

$$(13) \quad j_l \xrightarrow{\rho \rightarrow 0} \frac{\rho^l}{(2l+1)!!}$$

$$(14) \quad n_l \xrightarrow{\rho \rightarrow 0} -\frac{(2l-1)!!}{\rho^{l+1}}$$

We can verify the latter equation for  $j_l$  for a couple of cases with small  $l$ . From 9, we can generate a couple of  $j_l$ s:

$$(15) \quad j_0 = \frac{\sin \rho}{\rho}$$

$$(16) \quad j_1 = -\rho \frac{1}{\rho} \frac{d}{dr} \left( \frac{\sin \rho}{\rho} \right)$$

$$(17) \quad = -\rho \frac{1}{\rho} \left( \frac{\cos \rho}{\rho} - \frac{\sin \rho}{\rho^2} \right)$$

$$(18) \quad = \frac{\sin \rho}{\rho^2} - \frac{\cos \rho}{\rho}$$

$$(19) \quad j_2 = (-\rho)^2 \frac{1}{\rho} \frac{d}{d\rho} \left[ \frac{1}{\rho} \frac{d}{dr} \left( \frac{\sin \rho}{\rho} \right) \right]$$

$$(20) \quad (-\rho)^2 \frac{1}{\rho} \frac{d}{d\rho} \left[ \frac{1}{\rho} \left( \frac{\cos \rho}{\rho} - \frac{\sin \rho}{\rho^2} \right) \right]$$

$$(21) \quad = \left( \frac{3}{\rho^3} - \frac{1}{\rho} \right) \sin \rho - \frac{3 \cos \rho}{\rho^2}$$

We can get the limits for  $\rho \rightarrow 0$  by expanding the sine and cosine. That is, we use the limiting forms

$$(22) \quad \sin \rho \rightarrow \rho - \frac{\rho^3}{3!} + \dots$$

$$(23) \quad \cos \rho \rightarrow 1 - \frac{1}{2} \rho^2 + \dots$$

We need to retain enough terms for  $j_l$  so that we get all the terms up to the first power of  $\rho$  that doesn't cancel out when we do the algebra. We get

$$(24) \quad j_0 \rightarrow 1 = \frac{\rho^0}{1!!}$$

$$(25) \quad j_1 \rightarrow \frac{1}{\rho} - \frac{\rho}{6} - \frac{1}{\rho} \left(1 - \frac{1}{2}\rho^2\right)$$

$$(26) \quad = \frac{\rho}{3} = \frac{\rho^1}{3!!}$$

$$(27) \quad j_2 \rightarrow \left(\frac{3}{\rho^3} - \frac{1}{\rho}\right) \left(\rho - \frac{\rho^3}{6} + \frac{\rho^5}{120}\right) - \frac{3}{\rho^2} \left(1 - \frac{1}{2}\rho^2 + \frac{1}{24}\rho^4\right)$$

$$(28) \quad \rightarrow \left(\frac{1}{6} + \frac{1}{40} - \frac{1}{8}\right) \rho^2$$

$$(29) \quad = \frac{20 + 3 - 15}{120} \rho^2$$

$$(30) \quad = \frac{\rho^2}{15} = \frac{\rho^2}{5!!}$$

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