

ENSEMBLE OF ELECTRONS IN MAGNETIC FIELD

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Shankar, R. (1994), *Principles of Quantum Mechanics*, Plenum Press.

Exercise 14.4.6.

[If some equations are too small to read easily, use your browser's magnifying option (Ctrl + on Chrome, probably something similar on other browsers).]

As an example of a density matrix, we can apply it to an ensemble of spin $\frac{1}{2}$ particles. The density matrix is defined as

$$(0.1) \quad \rho \equiv \sum_i p_i |i\rangle \langle i|$$

where p_i is the probability of a single system being state $|i\rangle$. For a spin $\frac{1}{2}$ particle, there are only 2 states, so the density matrix can be written as a 2×2 matrix, once we define a basis for the states (for example, the basis of S_z states where $S_z = \pm \frac{\hbar}{2}$). Since any 2×2 matrix can be written as a linear combination of the Pauli matrices and the identity matrix, the density matrix can be written as

$$(0.2) \quad \rho = a_0 I + \mathbf{A} \cdot \boldsymbol{\sigma}$$

Since the trace of each of the Pauli matrices is zero, and $\text{Tr}I = 2$, we have

$$(0.3) \quad \text{Tr}\rho = 2a_0$$

However, we know that $\text{Tr}\rho = 1$, so we must have $a_0 = \frac{1}{2}$, so we can write

$$(0.4) \quad \rho = \frac{1}{2} (I + \mathbf{a} \cdot \boldsymbol{\sigma})$$

for some vector \mathbf{a} whose elements are complex numbers.

To find the average value $\langle \bar{\Omega} \rangle$ of an observable Ω in an ensemble, we can use the density matrix in the form

$$(0.5) \quad \langle \bar{\Omega} \rangle = \text{Tr}(\Omega \rho)$$

To find $\langle \vec{\sigma} \rangle$, we can work out each component separately. For σ_x we have, using the properties of the σ_i :

$$(0.6) \quad \langle \vec{\sigma}_x \rangle = \text{Tr}(\sigma_x \rho)$$

$$(0.7) \quad = \frac{1}{2} \text{Tr}(\sigma_x I + a_x \sigma_x^2 + a_y \sigma_x \sigma_y + a_z \sigma_x \sigma_z)$$

$$(0.8) \quad = \frac{1}{2} \text{Tr}(\sigma_x + a_x I + i a_y \sigma_z - i a_z \sigma_y)$$

$$(0.9) \quad = \frac{1}{2} (0 + 2a_x + 0 + 0)$$

$$(0.10) \quad = a_x$$

We can do similar calculations to get the other two components, with the result

$$(0.11) \quad \langle \vec{\sigma} \rangle = \mathbf{a}$$

Finally, suppose we have an ensemble of electrons in a constant magnetic field $\mathbf{B} = B\hat{\mathbf{z}}$, and that this ensemble is in thermal equilibrium at temperature T . A central result of statistical mechanics (which we haven't covered yet) is that particles in thermal equilibrium obey the Boltzmann distribution, which states that the probability of finding a particle with energy E in the ensemble is

$$(0.12) \quad p_E \propto e^{-E/kT}$$

where k is the Boltzmann constant. In this case, the energy is that of a magnetic moment μ in a constant magnetic field \mathbf{B} , which is

$$(0.13) \quad H = -\mu \cdot \mathbf{B} = -\gamma \mathbf{S} \cdot \mathbf{B} = -\gamma S_z B$$

There are only two states ($S_z = \pm \frac{\hbar}{2}$), so the probabilities are

$$(0.14) \quad p_{\uparrow} = \frac{1}{P} e^{\gamma B \hbar / 2kT}$$

$$(0.15) \quad p_{\downarrow} = \frac{1}{P} e^{-\gamma B \hbar / 2kT}$$

$$(0.16) \quad P = e^{\gamma B \hbar / 2kT} + e^{-\gamma B \hbar / 2kT}$$

The density matrix is therefore

$$(0.17) \quad \rho = \frac{1}{P} \left(e^{\gamma B \hbar / 2kT} |\uparrow\rangle \langle \uparrow| + e^{-\gamma B \hbar / 2kT} |\downarrow\rangle \langle \downarrow| \right)$$

In the S_z basis, this is

$$(0.18) \quad \rho = \frac{1}{P} \begin{bmatrix} e^{\gamma B \hbar / 2kT} & 0 \\ 0 & e^{-\gamma B \hbar / 2kT} \end{bmatrix}$$

We can work out the average magnetic moment for the ensemble as

$$(0.19) \quad \langle \bar{\mu} \rangle = \text{Tr}(\mu \rho)$$

$$(0.20) \quad = \frac{\hat{\mathbf{z}}}{P} \left[\frac{\gamma \hbar}{2} e^{\gamma B \hbar / 2kT} - \frac{\gamma \hbar}{2} e^{-\gamma B \hbar / 2kT} \right]$$

$$(0.21) \quad = \frac{e^{\gamma B \hbar / 2kT} - e^{-\gamma B \hbar / 2kT}}{e^{\gamma B \hbar / 2kT} + e^{-\gamma B \hbar / 2kT}} \frac{\gamma \hbar}{2} \hat{\mathbf{z}}$$

$$(0.22) \quad = \frac{\gamma \hbar}{2} \tanh \frac{\gamma B \hbar}{2kT} \hat{\mathbf{z}}$$

$$(0.23) \quad = -\frac{e \hbar}{2mc} \tanh \left(-\frac{e B \hbar}{2mckT} \right) \hat{\mathbf{z}}$$

$$(0.24) \quad = \frac{e \hbar}{2mc} \tanh \frac{e B \hbar}{2mckT} \hat{\mathbf{z}}$$