

TWO-STATE PARAMAGNET - NUMERICAL SOLUTION

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We can apply the formulas for entropy, temperature and heat capacity to a real-life system by looking at a two-state paramagnet. This is a system of N magnetic dipoles which, when placed in a magnetic field B , align themselves so that their magnetic moment μ points either parallel or antiparallel to the field. The energy of a dipole that is aligned with the field is lower, and we'll call it $-\mu B$, so that the antiparallel dipole has energy $+\mu B$, and the total energy of the system is

$$U = \mu B (N_{\downarrow} - N_{\uparrow}) = \mu B (N - 2N_{\uparrow}) \quad (1)$$

where an up arrow indicates parallel alignment and a down arrow antiparallel.

The net magnetization is then

$$M = \mu (N_{\uparrow} - N_{\downarrow}) = -\frac{U}{B} \quad (2)$$

The multiplicity of states is the same as a set of N coins with N_{\uparrow} heads, so

$$\Omega = \binom{N}{N_{\uparrow}} = \frac{N!}{N_{\uparrow}!(N - N_{\uparrow})!} \quad (3)$$

For small systems, we can find the entropy directly as

$$\frac{S}{k} = \ln \Omega \quad (4)$$

For $N_{\uparrow} = 98$, we get $U/\mu B = -96$, $M/N\mu = 0.96$, $\Omega = 4950$ and $S/k = 8.507$.

For each value of N_{\uparrow} from 0 up to N , we can evaluate U and S from the formulas above and then plot S versus U (I used Maple for the plot), as in Fig. 1.

For $-100 \leq U/\mu B < 0$, the curve is a 'normal' entropy curve in that the entropy increases with increasing energy, and the curve is concave down.

From this we can get the temperature

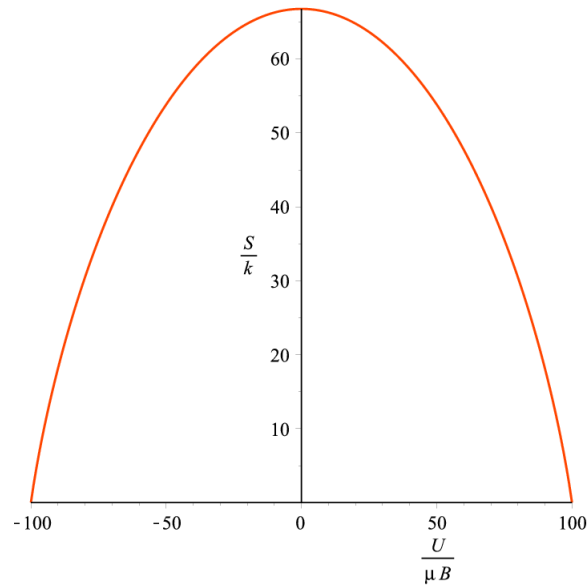


FIGURE 1. Entropy versus energy.

$$\frac{1}{T} = \frac{\partial S}{\partial U} \quad (5)$$

Thus the temperature increases with energy in the region $-100 \leq U/\mu B < 0$. At $U = 0$, however, the derivative is zero implying an infinite temperature, and for $U > 0$, the slope is negative, indicating a negative temperature. Since, in this case, negative temperatures occur at higher energies than positive temperatures, we have to interpret a negative temperature as actually being higher than a positive one, in fact, higher than an 'infinite' positive temperature.

In this case, it's probably better to use the entropy of the system as a physical measure of what's going on, since the second law implies that the system will tend to the energy with the greatest entropy, which is $U = 0$. At this energy, there are equal numbers of up and down dipoles, so the system is maximally randomized.

For a system with $N = 100$, we can approximate 5 by taking finite differences. Thus for a given value of N_{\uparrow} we can estimate the temperature as

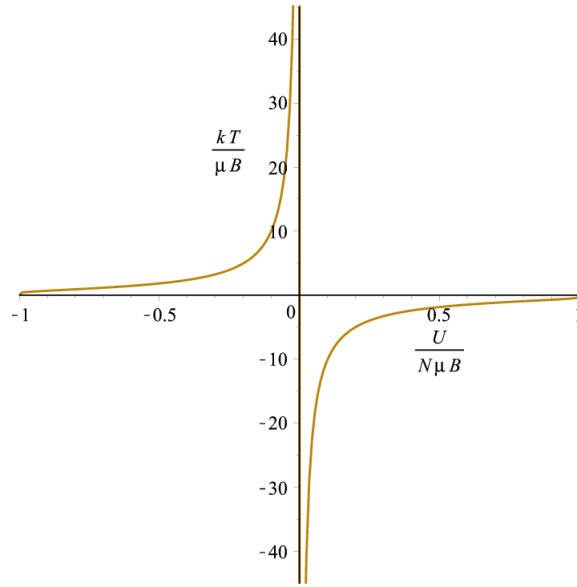


FIGURE 2. Temperature versus energy.

$$T(N_{\uparrow}) = \frac{\Delta U}{\Delta S} \quad (6)$$

$$= \frac{U(N_{\uparrow} + 1) - U(N_{\uparrow} - 1)}{S(N_{\uparrow} + 1) - S(N_{\uparrow} - 1)} \quad (7)$$

For $N_{\uparrow} = 98$, we get

$$T(98) = \frac{U(99) - U(97)}{S(99) - S(97)} = 0.541 \frac{\mu B}{k} \quad (8)$$

This allows us to plot temperature versus energy, as in Fig. 2.

We can see the flip over from $+\infty$ to $-\infty$ as the energy increases through zero.

The heat capacity can be obtained similarly as

$$C = \frac{\Delta U}{\Delta T} \quad (9)$$

$$= \frac{U(N_{\uparrow} + 1) - U(N_{\uparrow} - 1)}{T(N_{\uparrow} + 1) - T(N_{\uparrow} - 1)} \quad (10)$$

where we use 7 to calculate the temperatures. For $N_{\uparrow} = 98$ we have

$$C(98) = 0.310 N k \quad (11)$$

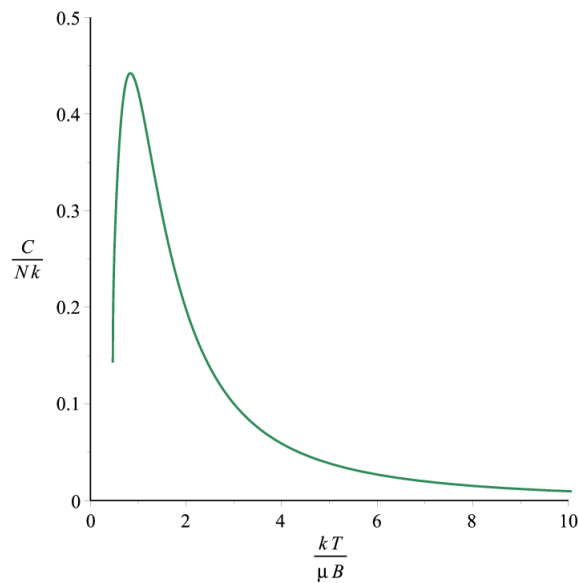


FIGURE 3. Heat capacity versus temperature.

The plot of C versus temperature is in Fig. 3.

[The plot does actually extend down to 0 at $T = 0$ if we use the analytic solution, but because we're dealing with discrete values of N_{\uparrow} , it cuts out early.] This curve is similar in shape to that of an Einstein solid at low temperatures.

Finally, we can plot the magnetization as a function of temperature, as in Fig. 4.

If we start off with $T > 0$ and lower the temperature towards zero, the magnetization gradually increases until at $T = 0$, the system is frozen into the state where all dipoles are parallel to the field. As we increase T to $+\infty$, we approach maximum randomness with equal numbers of dipoles pointing up and down so $M \rightarrow 0$. Increasing the temperature beyond $+\infty$ so it becomes negative (starting at $-\infty$) the magnetization again increases, but now the dipoles are aligned antiparallel to the field, eventually saturating as $T \rightarrow 0$ from the negative side.

PINGBACKS

Pingback: Two-state paramagnet - analytic solution

Pingback: Einstein solid - numerical solution

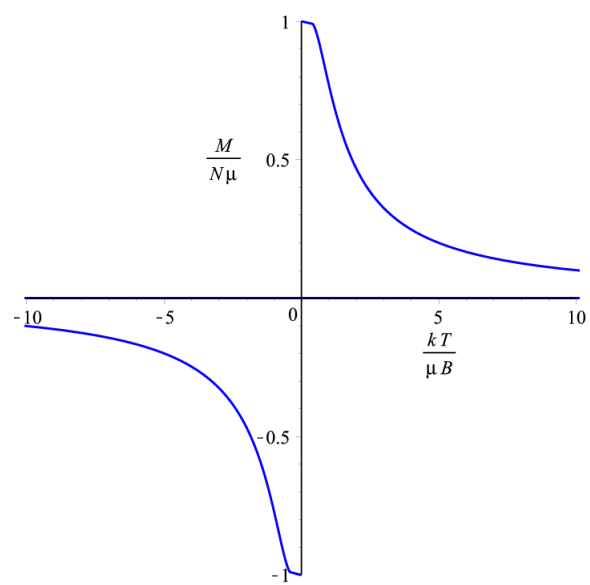


FIGURE 4. Magnetization versus temperature.