

## PROPAGATOR FOR FREE PARTICLE IN MOMENTUM SPACE

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Reference: Tom Lancaster and Stephen J. Blundell, *Quantum Field Theory for the Gifted Amateur*, (Oxford University Press, 2014), Problem 17.1.

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In section 17.1, L&B define the propagator for a particle in a general state  $|\Omega\rangle$  to be

$$G^+(x, y) = \theta(x^0 - y^0) \langle \Omega | \hat{\phi}(x) \hat{\phi}^\dagger(y) | \Omega \rangle \quad (1)$$

This is the probability amplitude that the system starts in state  $|\Omega\rangle$ , then at time  $y^0$  a particle is created at position  $\mathbf{y}$  and then at a later time  $x^0$ , the particle is annihilated at position  $\mathbf{x}$ . The step function  $\theta(x^0 - y^0)$  restricts  $x^0$  to be at a later time than  $y^0$ . The field operator  $\phi(x)$  (I'll leave off the hats from now on to save typing and clutter) is defined in the usual way in terms of creation and annihilation operators.

$$\phi(x) = \frac{1}{(2\pi)^{3/2}} \int \frac{d^3p}{\sqrt{2E_p}} \left( a(p) e^{-ip \cdot x} + a^\dagger(p) e^{ip \cdot x} \right) \quad (2)$$

In Exercise 17.1, L&B define the retarded field propagator for a free particle in momentum space as

$$G_0^+(\mathbf{p}, t_x, \mathbf{q}, t_y) = \theta(t_x - t_y) \langle 0 | a_{\mathbf{p}}(t_x) a_{\mathbf{q}}^\dagger(t_y) | 0 \rangle \quad (3)$$

It's not clear to me what a creation or annihilation operator is as a function of time, but I suspect what is meant is that a particle is created at time  $t_y$  with momentum  $\mathbf{q}$ . This particle then propagates to time  $t_x$  and is then subjected to annihilation if its momentum is then  $\mathbf{p}$ . The time evolution operator is the usual  $e^{-iHt}$ , where  $H$  is the hamiltonian, so we can write 3 as

$$G_0^+(\mathbf{p}, t_x, \mathbf{q}, t_y) = \theta(t_x - t_y) \langle 0 | a_{\mathbf{p}} e^{-iH(t_x - t_y)} a_{\mathbf{q}}^\dagger | 0 \rangle \quad (4)$$

That is, a particle is just created by  $a_{\mathbf{q}}^\dagger$ , then evolves for a time  $t_x - t_y$  and is then annihilated by  $a_{\mathbf{p}}$ .

We can now use

$$a_{\mathbf{q}}^\dagger |0\rangle = |\mathbf{q}\rangle \quad (5)$$

$$\langle 0| a_{\mathbf{p}} = \langle \mathbf{p}| \quad (6)$$

$$\langle \mathbf{p}|\mathbf{q}\rangle = \delta(\mathbf{p} - \mathbf{q}) \quad (7)$$

Also, since the states  $|\mathbf{p}\rangle$  and  $|\mathbf{q}\rangle$  are states with a fixed momentum, they are energy eigenstates so operating on them with  $H$  results in replacing  $H$  with  $E_{\mathbf{p}}$  or  $E_{\mathbf{q}}$ . Therefore we have

$$\langle 0| a_{\mathbf{p}} e^{-iH(t_x - t_y)} a_{\mathbf{q}}^\dagger |0\rangle = \langle 0| a_{\mathbf{p}} e^{-iE_{\mathbf{q}}t_x + iE_{\mathbf{q}}t_y} a_{\mathbf{q}}^\dagger |0\rangle \quad (8)$$

$$= e^{-iE_{\mathbf{q}}(t_x - t_y)} \langle \mathbf{p}|\mathbf{q}\rangle \quad (9)$$

$$= e^{-iE_{\mathbf{q}}(t_x - t_y)} \delta(\mathbf{p} - \mathbf{q}) \quad (10)$$

The delta function forces  $\mathbf{p} = \mathbf{q}$ , so this expression is equivalent to

$$G_0^+(\mathbf{p}, t_x, \mathbf{q}, t_y) = \theta(t_x - t_y) e^{-i(E_{\mathbf{p}}t_x - E_{\mathbf{q}}t_y)} \delta(\mathbf{p} - \mathbf{q}) \quad (11)$$

I'm not sure how we could justify splitting the exponential operator  $e^{-iH(t_x - t_y)}$  so that the  $t_x$  part operates on the  $\langle \mathbf{p}|$  to its left and the  $t_y$  part operates on the  $|\mathbf{q}\rangle$  to its right, so I'm not entirely sure that this is the correct way of solving this problem. Comments welcome.