

GREEN FUNCTION FOR KLEIN-GORDON EQUATION

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Reference: Tom Lancaster and Stephen J. Blundell, *Quantum Field Theory for the Gifted Amateur*, (Oxford University Press, 2014), Problem 17.2.

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In section 17.2, L&B derive an expression for the Feynman propagator, defined as

$$\Delta(x, y) = \langle 0 | T \phi(x) \phi^\dagger(y) | 0 \rangle \quad (1)$$

where T is the time ordering operator and ϕ is the scalar field operator. Their derivation results in equation 17.15:

$$\Delta(x, y) = \int \frac{d^3 p}{(2\pi)^3 2E_{\mathbf{p}}} \left[\theta(x^0 - y^0) e^{-ip \cdot (x-y)} + \theta(y^0 - x^0) e^{ip \cdot (x-y)} \right] \quad (2)$$

Here θ is the step function. Our task is to show that $\Delta(x, y)$ fills the role of a Green function for the Klein-Gordon equation, so that it satisfies

$$(\partial^2 + m^2) \Delta(x, y) = -i\delta(x - y) \quad (3)$$

where the derivatives on the LHS are with respect to components of x , not y .

Our task is made slightly easier since they require that we show this in a (1+1) dimensional system (one time dimension and one space dimension). That is, we have to show

$$(\partial_0^2 - \partial_1^2 + m^2) \Delta(x, y) = -i\delta(x^0 - y^0) \delta(x^1 - y^1) \quad (4)$$

To take things one step at a time, we'll split 2 into two parts:

$$\Delta_1 = \int \frac{dp}{(2\pi) 2E_{\mathbf{p}}} \theta(x^0 - y^0) e^{-ip \cdot (x-y)} \quad (5)$$

$$\Delta_2 = \int \frac{dp}{(2\pi) 2E_{\mathbf{p}}} \theta(y^0 - x^0) e^{ip \cdot (x-y)} \quad (6)$$

There is only a 2π in the denominator because we now have only one space dimension.

In what follows, we need the derivative of the step function:

$$\partial_0 \theta(x^0 - y^0) \equiv \frac{\partial}{\partial x^0} \theta(x^0 - y^0) = \delta(x^0 - y^0) \quad (7)$$

$$\frac{\partial}{\partial x^0} \theta(y^0 - x^0) = -\delta(y^0 - x^0) = -\delta(x^0 - y^0) \quad (8)$$

We also need the derivative of the delta function, which satisfies the condition for any function f :

$$f(x) \delta'(x) = -f'(x) \delta(x) \quad (9)$$

For Δ_1 , we have

$$\partial_0 \Delta_1(x, y) = \int \frac{dp}{(2\pi) 2E_{\mathbf{p}}} \left[\delta(x^0 - y^0) e^{-ip \cdot (x-y)} - iE_{\mathbf{p}} \theta(x^0 - y^0) e^{-ip \cdot (x-y)} \right] \quad (10)$$

$$\begin{aligned} \partial_0^2 \Delta_1(x, y) = \int \frac{dp}{(2\pi) 2E_{\mathbf{p}}} & \left[\delta'(x^0 - y^0) e^{-ip \cdot (x-y)} - iE_{\mathbf{p}} \delta(x^0 - y^0) e^{-ip \cdot (x-y)} \right. \\ & \left. - iE_{\mathbf{p}} \delta(x^0 - y^0) e^{-ip \cdot (x-y)} - E_{\mathbf{p}}^2 \theta(x^0 - y^0) e^{-ip \cdot (x-y)} \right] \end{aligned} \quad (11)$$

Using 9, the first term in the brackets is

$$\delta'(x^0 - y^0) e^{-ip \cdot (x-y)} = iE_{\mathbf{p}} \delta(x^0 - y^0) e^{-ip \cdot (x-y)} \quad (12)$$

Thus the first two terms in the brackets cancel and we're left with

$$\partial_0^2 \Delta_1(x, y) = - \int \frac{dp}{(2\pi) 2E_{\mathbf{p}}} \left[iE_{\mathbf{p}} \delta(x^0 - y^0) e^{-ip \cdot (x-y)} + E_{\mathbf{p}}^2 \theta(x^0 - y^0) e^{-ip \cdot (x-y)} \right] \quad (13)$$

We can do the same calculation on Δ_2 and use 8 to get

$$\partial_0^2 \Delta_2(x, y) = - \int \frac{dp}{(2\pi) 2E_{\mathbf{p}}} \left[iE_{\mathbf{p}} \delta(x^0 - y^0) e^{ip \cdot (x-y)} + E_{\mathbf{p}}^2 \theta(y^0 - x^0) e^{ip \cdot (x-y)} \right] \quad (14)$$

The spatial derivatives are

$$\partial_1^2 \Delta_1(x, y) = - \int \frac{dp}{(2\pi) 2E_{\mathbf{p}}} \theta(x^0 - y^0) (p^1)^2 e^{-ip \cdot (x-y)} \quad (15)$$

$$\partial_1^2 \Delta_2(x, y) = - \int \frac{dp}{(2\pi) 2E_{\mathbf{p}}} \theta(y^0 - x^0) (p^1)^2 e^{ip \cdot (x-y)} \quad (16)$$

Finally, the m^2 terms are

$$m^2 \Delta_1(x, y) = \int \frac{dp}{(2\pi) 2E_{\mathbf{p}}} m^2 \theta(x^0 - y^0) e^{-ip \cdot (x-y)} \quad (17)$$

$$m^2 \Delta_2(x, y) = \int \frac{dp}{(2\pi) 2E_{\mathbf{p}}} m^2 \theta(y^0 - x^0) e^{ip \cdot (x-y)} \quad (18)$$

We can now use

$$E_{\mathbf{p}}^2 = (p^1)^2 + m^2 \quad (19)$$

to combine the last term in 13 with 15 and 17 as specified on the LHS of 4:

$$\int \frac{dp}{(2\pi) 2E_{\mathbf{p}}} \left(-E_{\mathbf{p}}^2 + (p^1)^2 + m^2 \right) \theta(x^0 - y^0) e^{-ip \cdot (x-y)} = 0 \quad (20)$$

We get a similar result with the components from Δ_2 :

$$\int \frac{dp}{(2\pi) 2E_{\mathbf{p}}} \left(-E_{\mathbf{p}}^2 + (p^1)^2 + m^2 \right) \theta(y^0 - x^0) e^{ip \cdot (x-y)} = 0 \quad (21)$$

We're therefore left with

$$(\partial_0^2 - \partial_1^2 + m^2) \Delta(x, y) = - \int \frac{dp}{(2\pi) 2E_{\mathbf{p}}} \left(iE_{\mathbf{p}} \delta(x^0 - y^0) \left(e^{-ip \cdot (x-y)} + e^{ip \cdot (x-y)} \right) \right) \quad (22)$$

$$= -i \delta(x^0 - y^0) \int \frac{dp^1}{2\pi} e^{-ip^1(x^1 - y^1)} \quad (23)$$

$$= -i \delta(x^0 - y^0) \delta(x^1 - y^1) \quad (24)$$

We've converted the exponent $p \cdot (x - y)$ to just $p^1 (x^1 - y^1)$ since the delta function $\delta(x^0 - y^0)$ restricts $x^0 = y^0$. We've also used the standard definition of the delta function (in one dimension) as

$$\delta(x^0 - y^0) = \frac{1}{2\pi} \int dp^1 e^{-ip^1(x^1 - y^1)} = \frac{1}{2\pi} \int dp^1 e^{ip^1(x^1 - y^1)} \quad (25)$$

Note that it doesn't matter which sign we choose in the exponent since we're integrating over all p^1 .