

LORENTZ BOOST IN AN ARBITRARY DIRECTION

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References: Charles W. Misner, Kip S. Thorne & John Archibald Wheeler, *Gravitation*, W.H. Freeman (1973). Exercise 2.7.

Post date: 17 Jul 2020.

Usually we see Lorentz transformations for a boost parallel to one of the coordinate axes. MTW give the general form for a boost in an arbitrary direction. In matrix form, from an unprimed system to a primed system, we have

$$\Lambda = \begin{bmatrix} \gamma & -\beta\gamma n^1 & -\beta\gamma n^2 & -\beta\gamma n^3 \\ -\beta\gamma n^1 & 1 + (\gamma - 1)(n^1)^2 & (\gamma - 1)n^1 n^2 & (\gamma - 1)n^1 n^3 \\ -\beta\gamma n^2 & (\gamma - 1)n^1 n^2 & 1 + (\gamma - 1)(n^2)^2 & (\gamma - 1)n^2 n^3 \\ -\beta\gamma n^3 & (\gamma - 1)n^1 n^3 & (\gamma - 1)n^2 n^3 & 1 + (\gamma - 1)(n^3)^2 \end{bmatrix} \quad (1)$$

Here

$$\gamma = \frac{1}{\sqrt{1 - \beta^2}} \quad (2)$$

with β being the relative speed as usual, and

$$\mathbf{n} = [n^1, n^2, n^3] \quad (3)$$

is a unit vector, so that

$$(n^1)^2 + (n^2)^2 + (n^3)^2 = 1 \quad (4)$$

Ex 2.7(a). We are first asked to verify that Λ satisfies the relation

$$\Lambda^T \eta \Lambda = \eta \quad (5)$$

where η is the flat space metric

$$\eta = \begin{bmatrix} -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} \quad (6)$$

I'm not aware of any fast way of doing this, so it seems we need to resort to brute force. I used Maple to do the matrix multiplications and simplification using 2 and 4. The result is that 5 is, in fact, true.

I'll give a few of the matrix elements of 5 so you can see how they simplify. Defining

$$L \equiv \Lambda^T \eta \Lambda \quad (7)$$

we have

$$L[0,0] = \beta^2 \gamma^2 (n^1)^2 + \beta^2 \gamma^2 (n^2)^2 + \beta^2 \gamma^2 (n^3)^2 + -\gamma^2 \quad (8)$$

$$= \gamma^2 \left(\beta^2 \left((n^1)^2 + (n^2)^2 + (n^3)^2 \right) - 1 \right) \quad (9)$$

$$= \gamma^2 (\beta^2 - 1) \quad (10)$$

$$= \gamma^2 \left(-\frac{1}{\gamma^2} \right) \quad (11)$$

$$= -1 \quad (12)$$

Another example:

$$L[1,2] = -\beta^2 \gamma^2 n^1 n^2 + \left(1 + (\gamma - 1) (n^1)^2 \right) (\gamma - 1) n^1 n^2 + \quad (13)$$

$$(\gamma - 1) n^1 n^2 \left(1 + (\gamma - 1) (n^2)^2 \right) + (\gamma - 1)^2 n^1 (n^3)^2 n^2 \quad (14)$$

$$= n^1 n^2 \left[-\beta^2 \gamma^2 + 2(\gamma - 1) + (\gamma - 1)^2 \left((n^1)^2 + (n^2)^2 + (n^3)^2 \right) \right] \quad (15)$$

$$= n^1 n^2 \left[-\beta^2 \gamma^2 + 2(\gamma - 1) + (\gamma - 1)^2 \right] \quad (16)$$

$$= n^1 n^2 \left[\gamma^2 (1 - \beta^2) + 2\gamma - 2\gamma - 2 + 1 \right] \quad (17)$$

$$= n^1 n^2 [1 - 2 + 1] \quad (18)$$

$$= 0 \quad (19)$$

The other elements of L can be worked out similarly, if you have the patience. I'd imagine that you would only need to work out $L[0,0]$, one of the other diagonal elements and one of the off-diagonal elements, since the others would be obtained by permuting indices.

Ex 2.7(b&c) . The matrix 1 transforms from the unprimed to the primed system, so we have

$$x^{\alpha'} = \Lambda^{\alpha'}_{\beta} x^{\beta} \quad (20)$$

To work out the relative velocity of the two frames, we would like to find $dx^{i'}/dx^{0'}$. Using the entries from 1, we have (repeated Latin indexes are summed over 1 to 3):

$$dx^{i'} = (\gamma - 1)n^i n^j dx^j + dx^i - \beta \gamma n^i dx^0 \quad (21)$$

$$dx^{0'} = \gamma dx^0 - \beta \gamma n^j dx^j \quad (22)$$

We can now multiply 21 on both sides by n^i and sum over i , using 4:

$$n^i dx^{i'} = (\gamma - 1)n^i n^i n^j dx^j + n^i dx^i - \beta \gamma n^i n^i dx^0 \quad (23)$$

$$= (\gamma - 1)n^j dx^j + n^i dx^i - \beta \gamma dx^0 \quad (24)$$

We also multiply 22 on both sides by β :

$$\beta dx^{0'} = \beta \gamma dx^0 - \beta^2 \gamma n^j dx^j \quad (25)$$

Adding this to 24 we have

$$n^i dx^{i'} + \beta dx^{0'} = (\gamma - 1)n^j dx^j + n^i dx^i - \beta^2 \gamma n^j dx^j \quad (26)$$

$$= \gamma (1 - \beta^2) n^i dx^i \quad (27)$$

$$= \frac{1}{\gamma} n^i dx^i \quad (28)$$

where we relabelled the dummy index j to i , since it's summed.

Now if we want the relative velocity of the two frames, we want to see how the location of a single point in one frame is viewed in the other. That is, we can take $dx^i = 0$ in the unprimed frame to see how fast a point in the unprimed frame is seen to move in the primed frame. Doing this gives us

$$n^i \frac{dx^{i'}}{dx^{0'}} = -\beta \quad (29)$$

Defining the vector \mathbf{v}' to be

$$\mathbf{v}' \equiv \left[\frac{dx^{1'}}{dx^{0'}}, \frac{dx^{2'}}{dx^{0'}}, \frac{dx^{3'}}{dx^{0'}} \right] \quad (30)$$

we have

$$\mathbf{n} \cdot \mathbf{v}' = -\beta \quad (31)$$

Since \mathbf{n} is a unit vector, we can take

$$\mathbf{v}' = -\beta [n^1, n^2, n^3] = -\beta \mathbf{n} \quad (32)$$

Strictly speaking, \mathbf{v}' could also have a component perpendicular to \mathbf{n} , since this would disappear in the dot product. However, movement perpendicular to the relative velocity is not affected by a Lorentz transformation.

To find the velocity of the primed frame as seen in the unprimed frame, we just replace β by $-\beta$ in the above argument and everything works out the same except for the flipped sign, so we have

$$\mathbf{v} = \beta \mathbf{n} \quad (33)$$

Ex 2.7(d). To verify that 1 reduces to the usual Lorentz transformation if the velocity is in the z direction, we set $n^1 = n^2 = 0$ and $n^3 = 1$ which gives

$$\|\Lambda^{\nu'}_{\mu}\| = \begin{bmatrix} \gamma & 0 & 0 & -\beta\gamma \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ -\beta\gamma & 0 & 0 & \gamma \end{bmatrix} \quad (34)$$

This is the transformation from the unprimed to the primed system. To get the inverse transformation, we replace β by $-\beta$:

$$\|\Lambda^{\mu}_{\nu'}\| = \begin{bmatrix} \gamma & 0 & 0 & \beta\gamma \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ \beta\gamma & 0 & 0 & \gamma \end{bmatrix} \quad (35)$$

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