

RAISING AND LOWERING TENSOR INDICES

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References: Charles W. Misner, Kip S. Thorne & John Archibald Wheeler, *Gravitation*, W.H. Freeman (1973). Sections 2.9 & 3.2; Exercise 3.3.

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A general tensor S is a function of a number of one-forms and/or vectors which outputs either a scalar or vector. In a given Lorentz frame, the components of a tensor are found by the action of the tensor on the basis one-forms and vectors.

$$S(\boldsymbol{\sigma}, \boldsymbol{\rho}, \boldsymbol{v}) = S^{\alpha\beta}{}_{\gamma} \sigma_{\alpha} \rho_{\beta} v^{\gamma} \quad (1)$$

where $\boldsymbol{\sigma}$ and $\boldsymbol{\rho}$ are one-forms and \boldsymbol{v} is a vector.

The components of a vector can be converted to the components of the corresponding one-form, and vice versa, using the metric tensor η in an operation known as raising or lowering the indices. The components of a one-form have lower indices which are raised by the operation

$$\sigma^{\alpha} = \eta^{\alpha\beta} \sigma_{\beta} \quad (2)$$

so that σ^{α} are the components of the vector $\boldsymbol{\sigma}$ corresponding to the one-form $\boldsymbol{\sigma}$. MTW use the same symbol for both the one-form and the vector as they are equivalent in the sense that they contain the same information.

The components of a vector have upper indices which can be lowered by

$$v_{\alpha} = \eta_{\alpha\beta} v^{\beta} \quad (3)$$

to convert a vector into its corresponding one-form.

The tensor S itself, as a function of its one-form and vector arguments, is independent of whether its arguments are expressed as one-forms or vectors. That is, we require

$$S^{\alpha\beta}{}_{\gamma} \sigma_{\alpha} \rho_{\beta} v^{\gamma} = S^{\alpha}{}_{\beta\gamma} \sigma_{\alpha} \rho^{\beta} v^{\gamma} \quad (4)$$

$$= S^{\alpha\beta\gamma} \sigma_{\alpha} \rho_{\beta} v_{\gamma} \quad (5)$$

$$= S^{\alpha}{}_{\beta}{}^{\gamma} \sigma_{\alpha} \rho^{\beta} v_{\gamma} \quad (6)$$

and so on for all other possible positions of the indices α , β and γ .

Ex. 3.3. Using 2 and 3 we can relate the various forms of tensor components to each other. For example, we have

$$S^{\alpha\beta}{}_{\gamma}\sigma_{\alpha\rho\beta}v^{\gamma} = S^{\alpha}{}_{\beta\gamma}\sigma_{\alpha}\rho^{\beta}v^{\gamma} \quad (7)$$

$$= S^{\alpha}{}_{\beta\gamma}\sigma_{\alpha}\eta^{\beta\delta}\rho_{\delta}v^{\gamma} \quad (8)$$

$$= \eta^{\beta\delta}S^{\alpha}{}_{\beta\gamma}\sigma_{\alpha}\rho_{\delta}v^{\gamma} \quad (9)$$

$$= \eta^{\delta\beta}S^{\alpha}{}_{\delta\gamma}\sigma_{\alpha}\rho_{\beta}v^{\gamma} \quad (10)$$

where in the last line we swapped the dummy indices $\beta \leftrightarrow \delta$, which we can do because they are summed. Comparing the LHS and RHS, we see that, since this relation must be true for all vectors and one-forms, we must have

$$S^{\alpha\beta}{}_{\gamma} = \eta^{\delta\beta}S^{\alpha}{}_{\delta\gamma} \quad (11)$$

That is, we can raise an index of a tensor components by using the metric η . A similar argument shows that we can lower an index in the same way, or we can use the fact that the inverse of the metric $\eta^{\delta\beta}$ is the same metric with lower indices:

$$\|\eta_{\alpha\beta}\| = \|\eta^{\alpha\beta}\|^{-1} \quad (12)$$

Multiplying 11 both sides by $\eta_{\beta\epsilon}$ and summing, we get

$$\eta_{\beta\epsilon}S^{\alpha\beta}{}_{\gamma} = \eta_{\beta\epsilon}\eta^{\delta\beta}S^{\alpha}{}_{\delta\gamma} \quad (13)$$

$$= \delta^{\delta}_{\epsilon}S^{\alpha}{}_{\delta\gamma} \quad (14)$$

$$= S^{\alpha}{}_{\epsilon\gamma} \quad (15)$$

Thus we've lowered the β index using the lower form of the metric $\eta_{\beta\epsilon}$.

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