

DUAL TENSORS

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References: Charles W. Misner, Kip S. Thorne & John Archibald Wheeler, *Gravitation*, W.H. Freeman (1973). Exercise 3.14

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MTW define a *dual* with reference to vectors and rank-2 and rank-3 antisymmetric tensors. A dual is indicated by an asterisk before the tensor symbol, and the definitions are

$$*J_{\alpha\beta\gamma} \equiv J^\mu \varepsilon_{\mu\alpha\beta\gamma} \quad (1)$$

$$*F_{\alpha\beta} \equiv \frac{1}{2} F^{\mu\nu} \varepsilon_{\mu\nu\alpha\beta} \quad (2)$$

$$*B_\alpha \equiv \frac{1}{3!} B^{\lambda\mu\nu} \varepsilon_{\lambda\mu\nu\alpha} \quad (3)$$

Here, J^μ are the components of a vector J , $F^{\mu\nu}$ are the components of an antisymmetric rank-2 tensor, and $B^{\lambda\mu\nu}$ are the components of a rank-3 antisymmetric tensor.

Ex 3.14 (a) We are to show that taking the dual of a dual returns the original object, up to a sign. First, we look at the vector J , and we swap upper and lower indices in 1 so we can write

$$*J^{\alpha\beta\gamma} = J_\mu \varepsilon^{\mu\alpha\beta\gamma} \quad (4)$$

Taking the dual of this, we get, using 3

$$**J_\delta = \frac{1}{3!} J_\mu \varepsilon^{\mu\alpha\beta\gamma} \varepsilon_{\alpha\beta\gamma\delta} \quad (5)$$

$$= -\frac{1}{3!} J_\mu \varepsilon^{\alpha\beta\gamma\mu} \varepsilon_{\alpha\beta\gamma\delta} \quad (6)$$

where we swapped the μ index in $\varepsilon^{\mu\alpha\beta\gamma}$ 3 times to produce $\varepsilon^{\alpha\beta\gamma\mu}$. Since this is an odd number of swaps, the sign gets flipped.

We can now use the earlier result

$$\varepsilon^{\alpha\beta\gamma\mu} \varepsilon_{\alpha\beta\gamma\delta} = -6\delta^\mu_\delta \quad (7)$$

so we have

$$**J_\delta = -\frac{1}{3!}J_\mu(-6\delta^\mu_\delta) = J_\delta \quad (8)$$

Thus we have the result

$$**J = J \quad (9)$$

Now consider F . We have

$$*F^{\alpha\beta} = \frac{1}{2}F_{\mu\nu}\varepsilon^{\mu\nu\alpha\beta} \quad (10)$$

so the double-dual is, using 2

$$**F_{\gamma\delta} = \frac{1}{2}\left(\frac{1}{2}F_{\mu\nu}\varepsilon^{\mu\nu\alpha\beta}\right)\varepsilon_{\alpha\beta\gamma\delta} \quad (11)$$

$$= \frac{1}{4}F_{\mu\nu}\varepsilon^{\mu\nu\alpha\beta}\varepsilon_{\gamma\delta\alpha\beta} \quad (12)$$

where we performed an even number of swaps to map the indices in $\varepsilon_{\alpha\beta\gamma\delta}$ to $\varepsilon_{\gamma\delta\alpha\beta}$, so there is no change of sign. We now use the earlier result

$$\varepsilon^{\mu\nu\alpha\beta}\varepsilon_{\gamma\delta\alpha\beta} = -2\delta^{\mu\nu}_{\gamma\delta} \quad (13)$$

so we have

$$**F_{\gamma\delta} = -\frac{1}{2}F_{\mu\nu}\delta^{\mu\nu}_{\gamma\delta} \quad (14)$$

We know that

$$\delta^{\mu\nu}_{\gamma\delta} = \begin{cases} +1 & \text{if } \gamma\delta \text{ is an even permutation of } \mu\nu \\ -1 & \text{if } \gamma\delta \text{ is an odd permutation of } \mu\nu \\ 0 & \text{otherwise} \end{cases} \quad (15)$$

so doing the sum over $\mu\nu$ in 14 we have

$$**F_{\gamma\delta} = -\frac{1}{2}(F_{\gamma\delta} - F_{\delta\gamma}) \quad (16)$$

$$= -\frac{1}{2}(F_{\gamma\delta} + F_{\gamma\delta}) \quad (17)$$

$$= -F_{\gamma\delta} \quad (18)$$

where in the second line we used the condition that F is antisymmetric. Therefore

$$**F = -F \quad (19)$$

Finally, we look at the double dual of B . We have

$$*B^\alpha = \frac{1}{3!} B_{\lambda\mu\nu} \varepsilon^{\lambda\mu\nu\alpha} \quad (20)$$

so we have from 1

$$**B_{\beta\gamma\delta} = \frac{1}{3!} B_{\lambda\mu\nu} \varepsilon^{\lambda\mu\nu\alpha} \varepsilon_{\alpha\beta\gamma\delta} \quad (21)$$

$$= -\frac{1}{3!} B_{\lambda\mu\nu} \varepsilon^{\lambda\mu\nu\alpha} \varepsilon_{\beta\gamma\delta\alpha} \quad (22)$$

where we used to 3 swaps to convert $\varepsilon_{\alpha\beta\gamma\delta}$ to $-\varepsilon_{\beta\gamma\delta\alpha}$. Using the earlier result

$$\varepsilon^{\lambda\mu\nu\alpha} \varepsilon_{\beta\gamma\delta\alpha} = -\delta^{\lambda\mu\nu}_{\beta\gamma\delta} \quad (23)$$

we have

$$**B_{\beta\gamma\delta} = \frac{1}{3!} B_{\lambda\mu\nu} \delta^{\lambda\mu\nu}_{\beta\gamma\delta} \quad (24)$$

We know that

$$\delta^{\lambda\mu\nu}_{\beta\gamma\delta} = \begin{cases} +1 & \text{if } \lambda\mu\nu \text{ is an even permutation of } \beta\gamma\delta \\ -1 & \text{if } \lambda\mu\nu \text{ is an odd permutation of } \beta\gamma\delta \\ 0 & \text{otherwise} \end{cases} \quad (25)$$

so doing the sum over $\lambda\mu\nu$ in 24 gives us

$$**B_{\beta\gamma\delta} = \frac{1}{3!} (B_{\beta\gamma\delta} + B_{\gamma\delta\beta} + B_{\delta\beta\gamma} - B_{\gamma\beta\delta} - B_{\beta\delta\gamma} - B_{\delta\gamma\beta}) \quad (26)$$

$$= \frac{1}{3!} \times 6B_{\beta\gamma\delta} \quad (27)$$

$$= B_{\beta\gamma\delta} \quad (28)$$

where the second line follows because B is totally antisymmetric, so an even permutation of $\beta\gamma\delta$ gives $B_{\beta\gamma\delta}$ and an odd permutation gives $-B_{\beta\gamma\delta}$. Therefore

$$**B = B \quad (29)$$

(b) We can write out the components to verify these conclusions. From 1 we have, for example

$$\begin{aligned}
*J_{012} &= J^3 \varepsilon_{3012} = -J^3 \\
*J_{123} &= J^0 \varepsilon_{0123} = J^0 \\
*J_{230} &= J^1 \varepsilon_{1230} = -J^1 \\
*J_{013} &= J^2 \varepsilon_{2013} = J^2
\end{aligned} \tag{30}$$

Permutations of the indices on the LHS can be found by swapping the indices in the ε and deducing the corresponding sign. Thus $*J$ and J contain the same information.

For F , we have from 2, for example (remember F is antisymmetric):

$$\begin{aligned}
*F_{01} &= \frac{1}{2} F^{\mu\nu} \varepsilon_{\mu\nu 01} = \frac{1}{2} (F^{23} - F^{32}) = F^{23} \\
*F_{02} &= \frac{1}{2} F^{\mu\nu} \varepsilon_{\mu\nu 02} = \frac{1}{2} (-F^{13} + F^{31}) = -F^{13} \\
*F_{23} &= \frac{1}{2} F^{\mu\nu} \varepsilon_{\mu\nu 23} = \frac{1}{2} (F^{01} - F^{10}) = F^{01}
\end{aligned} \tag{31}$$

and so on.

For B , we have from 3, for example

$$\begin{aligned}
*B_0 &= \frac{1}{6} B^{\mu\nu\lambda} \varepsilon_{\mu\nu\lambda 0} \\
&= \frac{1}{6} (-B^{123} - B^{231} - B^{312} + B^{213} + B^{132} + B^{321}) \\
&= -B^{123}
\end{aligned} \tag{32}$$

and so on.

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