

CHARGE CONSERVATION

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References: Charles W. Misner, Kip S. Thorne & John Archibald Wheeler, *Gravitation*, W.H. Freeman (1973). Exercise 3.16

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From the expression of Maxwell's equations given earlier

$$F^{\alpha\beta}{}_{,\beta} = 4\pi J^\alpha \quad (1)$$

we may derive the equation of conservation of charge. We take the divergence of 1 with respect to the index α to get

$$F^{\alpha\beta}{}_{,\beta,\alpha} = 4\pi J^\alpha{}_{,\alpha} \quad (2)$$

We can write the LHS as

$$F^{\alpha\beta}{}_{,\beta,\alpha} = F^{\beta\alpha}{}_{,\alpha,\beta} \quad (3)$$

$$= F^{\beta\alpha}{}_{,\beta,\alpha} \quad (4)$$

$$= -F^{\alpha\beta}{}_{,\beta,\alpha} \quad (5)$$

In the first line, we swapped the dummy labels $\alpha \leftrightarrow \beta$. In the second line, we reversed the order of derivatives, since the order doesn't matter. In the third line, we used the antisymmetry of $F^{\alpha\beta} = -F^{\beta\alpha}$. Thus

$$F^{\alpha\beta}{}_{,\beta,\alpha} = -F^{\alpha\beta}{}_{,\beta,\alpha} \quad (6)$$

so we must have

$$F^{\alpha\beta}{}_{,\beta,\alpha} = 4\pi J^\alpha{}_{,\alpha} = 0 \quad (7)$$

Since the four-current J is composed of the 3-current vector \mathbf{J} and the charge density $J^0 = \rho$, the RHS of this last equation can be written as

$$J^\alpha{}_{,\alpha} = \nabla \cdot \mathbf{J} + \frac{\partial J^0}{\partial t} = 0 \quad (8)$$

Integrating over all space and using Gauss's theorem to convert the volume integral of a divergence to a surface integral, we have

$$\int_V \nabla \cdot \mathbf{J} d^3x = -\frac{\partial}{\partial t} \int_V J^0 d^3x \quad (9)$$

$$\int_A \mathbf{J} \cdot \hat{\mathbf{n}} d\mathbf{a} = -\frac{\partial Q}{\partial t} \quad (10)$$

This states that the integral of the current over a surface A bounding the volume V is the negative of the rate of change of charge Q within V . If we take the surface far enough out so that no current crosses it, then the surface integral is zero, and the total charge remains unchanged.

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