

## GAUSS'S LAW FOR ELECTRON-POSITRON CREATION IN STARS

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References: Charles W. Misner, Kip S. Thorne & John Archibald Wheeler, *Gravitation*, W.H. Freeman (1973). Exercise 5.3.

Post date: 22 Aug 2020.

In some stars, the temperature is high enough that there is enough energy to create electron-positron pairs, which can collide later to annihilate themselves. We can describe the number flux of these particles by the tensor  $S$ . This is analogous to the electromagnetic four-current  $J$  that we looked at earlier. In the four-current example, total charge is conserved and the four-current divergence is zero:

$$\nabla \cdot J = 0 \quad (1)$$

In the pair creation case, the total number of particles is not conserved, so the divergence is non-zero, and we have

$$\epsilon \equiv \nabla \cdot S \quad (2)$$

The components of  $S$ , by analogy with  $J$  are:  $S^0$  is the number density of particles, and the spatial components  $S^j$  represent the flux of particles in each of the three spatial directions.

Using Gauss's theorem, we can relate the volume integral to a surface integral, so we have

$$\int_{\mathcal{V}} \nabla \cdot S \, d^4\Omega = \oint_{\partial\mathcal{V}} S \cdot d^3\Sigma \quad (3)$$

where  $\mathcal{V}$  is a spacetime volume enclosed by surface  $\partial\mathcal{V}$ , and  $d^4\Omega$  is a four-volume element and  $d^3\Sigma$  is an outward-facing 3-surface element.

To make this concrete, we can consider case (b) from our earlier treatment of the electromagnetic four-current. Here, we have a surface bounded by two constant times  $t_1$  and  $t_2$  (as viewed in one Lorentz frame) and a spatial volume that covers all space. Thus we're considering how the number of particles in all space changes between two fixed times. In this case, the surface  $\partial\mathcal{V}$  can be divided into two surfaces: a surface  $S_1$  at  $t = t_1$  and a second surface  $S_2$  at  $t = t_2$ . These two surfaces are assumed to be connected by two timelike surfaces at infinity, and we further assume that  $S$  goes to

zero fast enough at infinity that any integrals over these surfaces can be neglected. This is certainly true if the volume we're considering is that of a star.

In this case, we can write 3 in component form as

$$\oint_{\partial\mathcal{V}} \mathbf{S} \cdot d^3\boldsymbol{\Sigma} = \int_{S_2} S^0 dx dy dz - \int_{S_1} S^0 dx dy dz \quad (4)$$

$$= \int_{\mathcal{V}} \nabla \cdot \mathbf{S} d^4\Omega \quad (5)$$

$$= \int_{\mathcal{V}} \epsilon d^4\Omega \quad (6)$$

The expression 4 is the difference in particle number between times  $t_1$  and  $t_2$ , and is equal, via 2, to the volume integral of  $\epsilon$  in 6. This shows that  $\epsilon$  is the rate of particle creation per unit volume and unit time.

The argument can be extended to other spacetime volumes in the same way as we did for the four-current, so that 3 can be written as

$$\int_{\mathcal{V}} \nabla \cdot \mathbf{S} d^4\Omega = \oint_{\partial\mathcal{V}} \mathbf{S} \cdot d^3\boldsymbol{\Sigma} = \int_{\mathcal{V}} \epsilon d^4\Omega \quad (7)$$

In other words, the total number of particles created (minus those annihilated) in a spacetime volume  $\mathcal{V}$ , represented by  $\int_{\mathcal{V}} \epsilon d^4\Omega$ , is equal to the four-flux  $\mathbf{S}$  integrated over the surface  $\partial\mathcal{V}$  bounding that spacetime volume.